

# The chemistry and thermodynamics of Pop III star formation

# Where do the first stars form

- Form in dark matter minihalos with mass

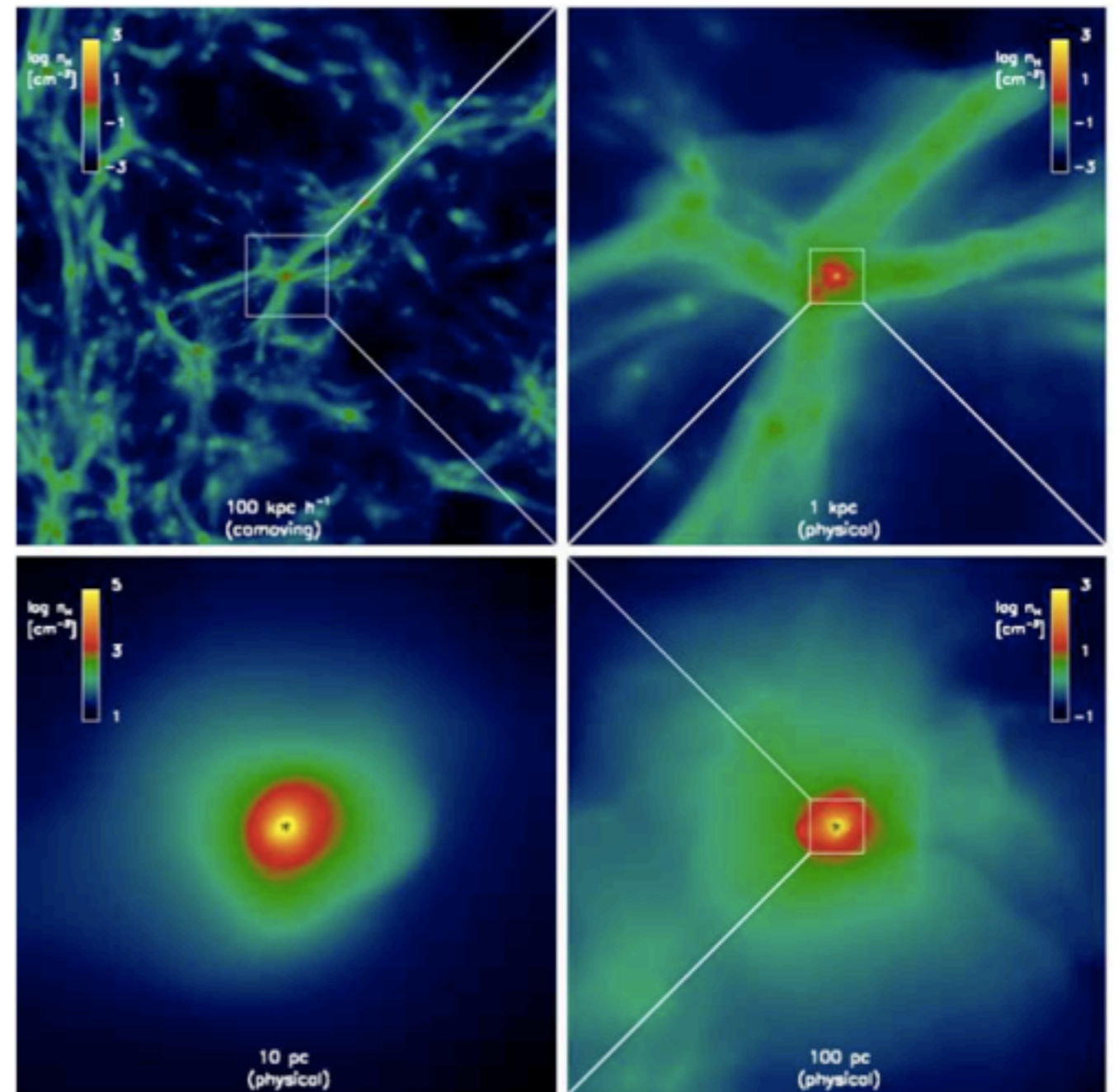
$$M_{\text{halo}} \gtrsim 5 \times 10^5 M_{\odot}$$

- Redshift  $z = 16 - 20$

- $T_{\text{vir}} \sim 1000 \text{ K}$

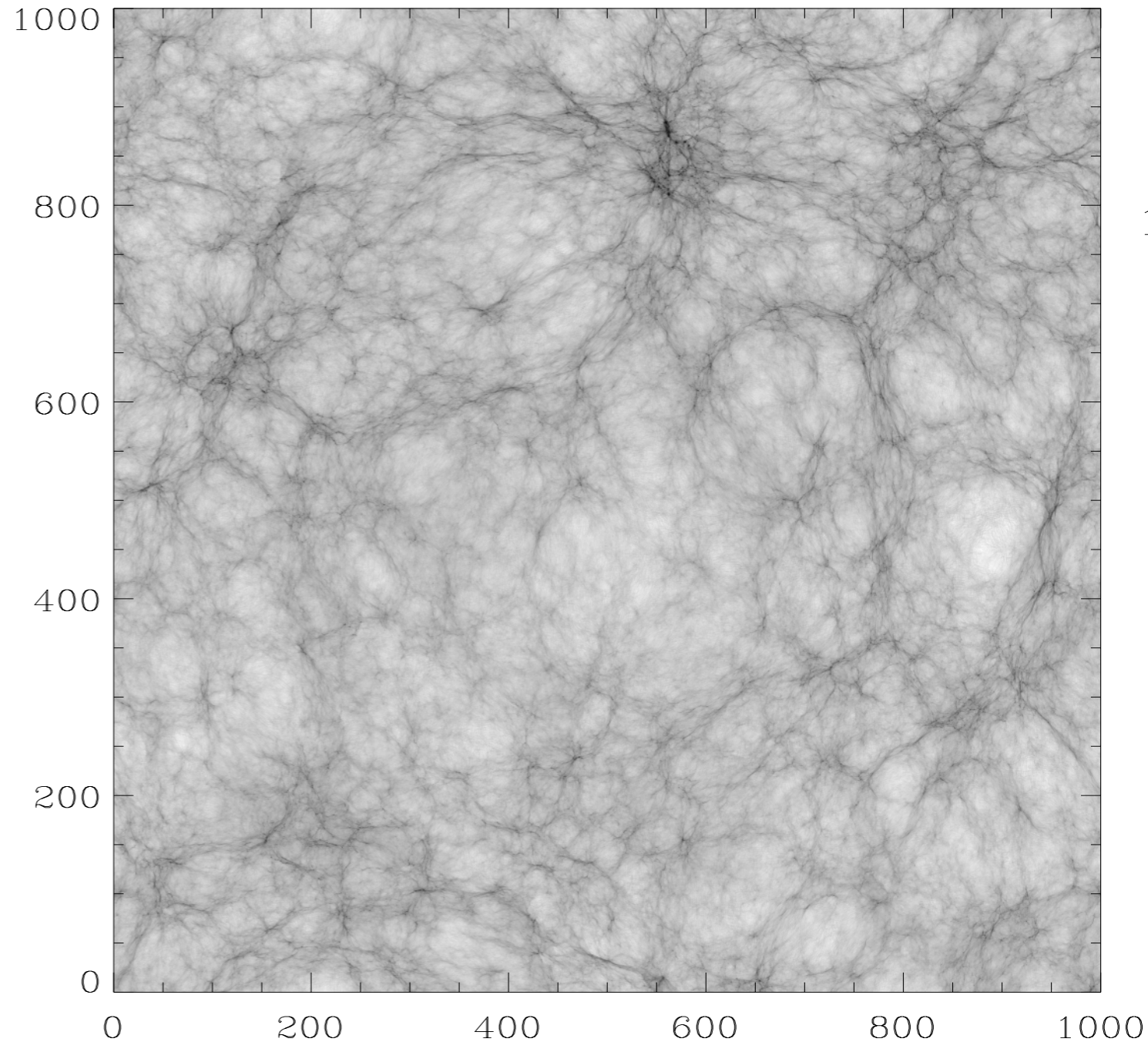
- Gas density is around

$$1 \text{ cm}^{-3}$$



# Current state-of-the-art?

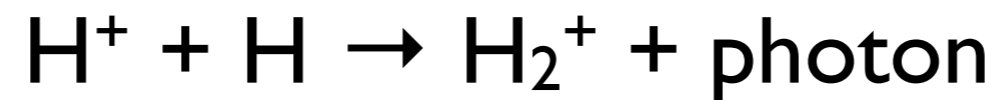
- $2048^3$  DM particles
- 1 Mpc (co-moving)
- Particle mass  $\sim 9 M_{\odot}$



Sasaki et al. (2013, in prep.)

# The H<sub>2</sub> chemistry

- The ion-neutral reactions dominate in primordial gas



- As we pointed out in lecture 2, these reactions depend on the ionisation state of the gas.
- Electrons and protons act as catalysts in these reactions.

# Why do they form in these minihalos?

- Collapse occurs in minihalos with  $t_{\text{cool}} < t_{\text{H}}$
- Below 7000 K,  $\text{H}_2$  is the only useful coolant
- Collapse only occurs in minihalos that can form enough  $\text{H}_2$  to permit them to cool
- Amount of  $\text{H}_2$  produced is a strong function of temperature.
- Once  $T_{\text{vir}}$  reaches  $\sim 1000$  K, enough  $\text{H}_2$  can form to satisfy the above cooling time requirement.

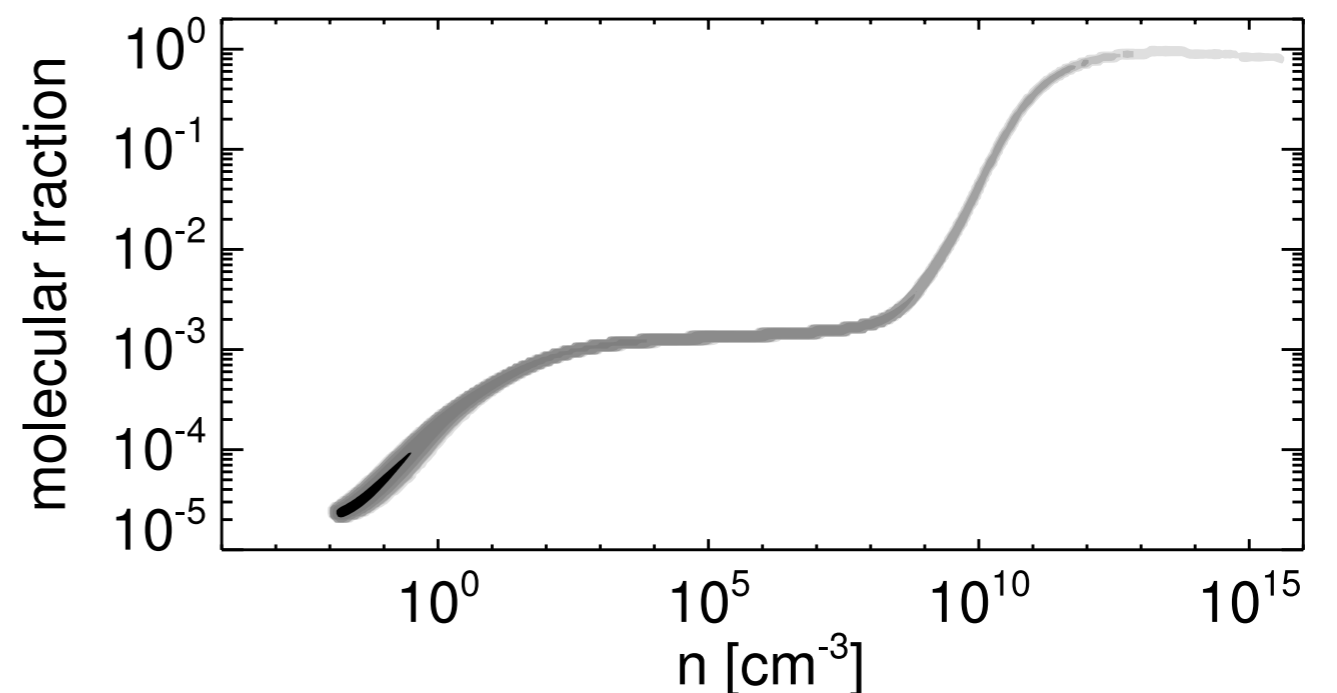
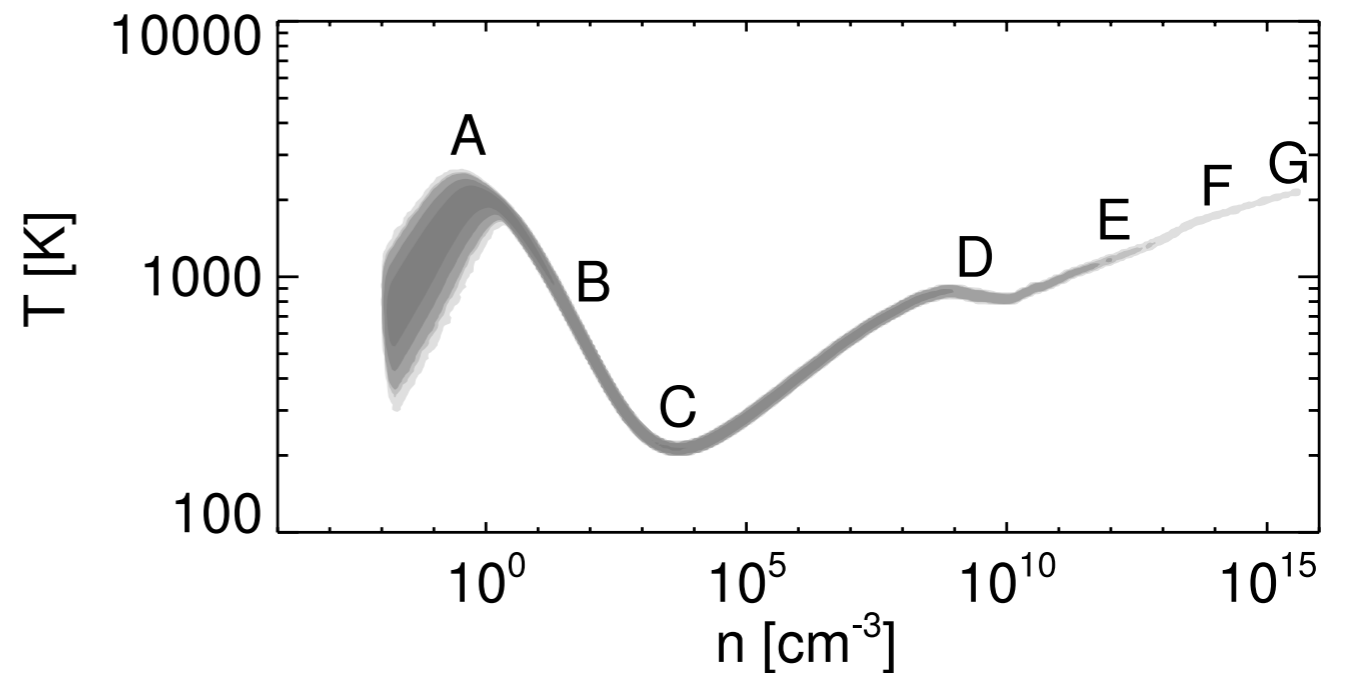
# What sets the H<sub>2</sub> fraction?

- Rate of ion-neutral reactions depend on the ionisation fraction in the gas.
- Ionisation fraction is reduced by recombination.
- H<sub>2</sub> formation is a race against the recombination!
- Limits the amount of H<sub>2</sub> that a halo can form by this process to around  $n_{\text{H}_2} \sim 10^{-3} n$
- Sets a limit to the cooling available in the halos as they collapse.

# So what happens next?

- The evolution of the temperature during the collapse is not well understood.
- Characterised by several key stages.
- Different H<sub>2</sub> physics (formation/heating/cooling) controls each stage.

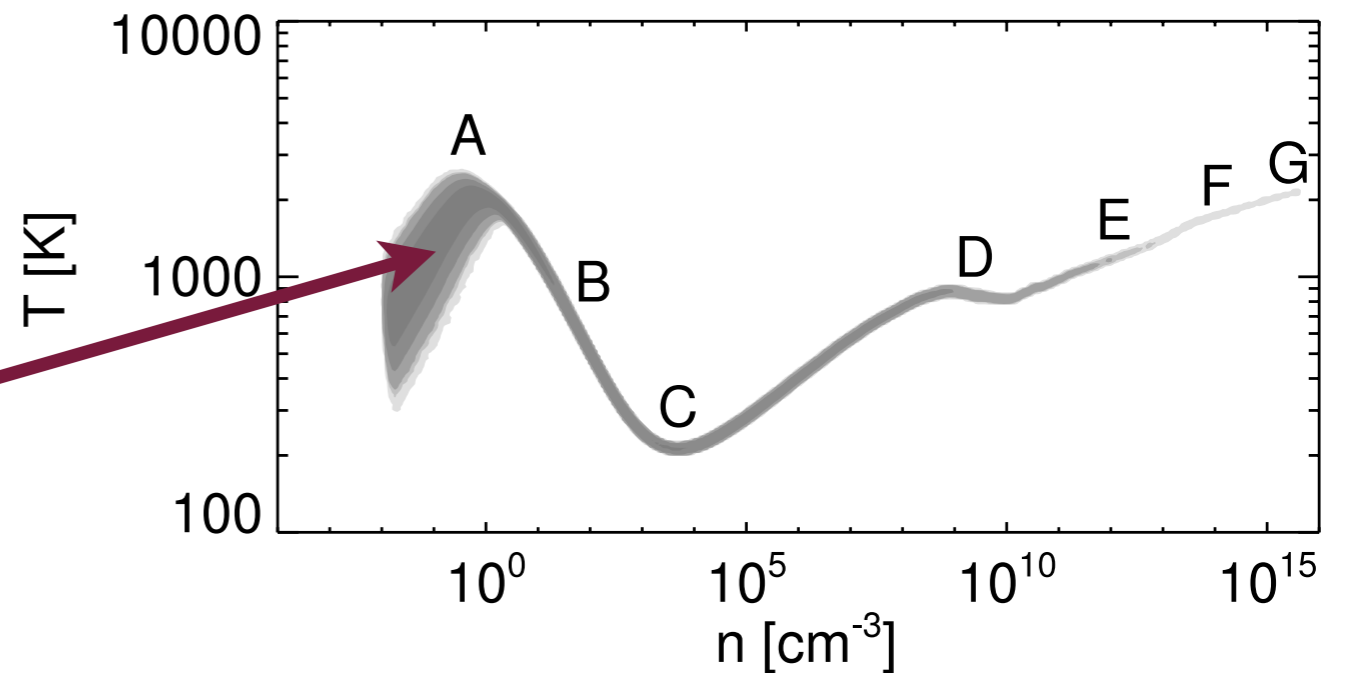
Yoshida et al. (2006)



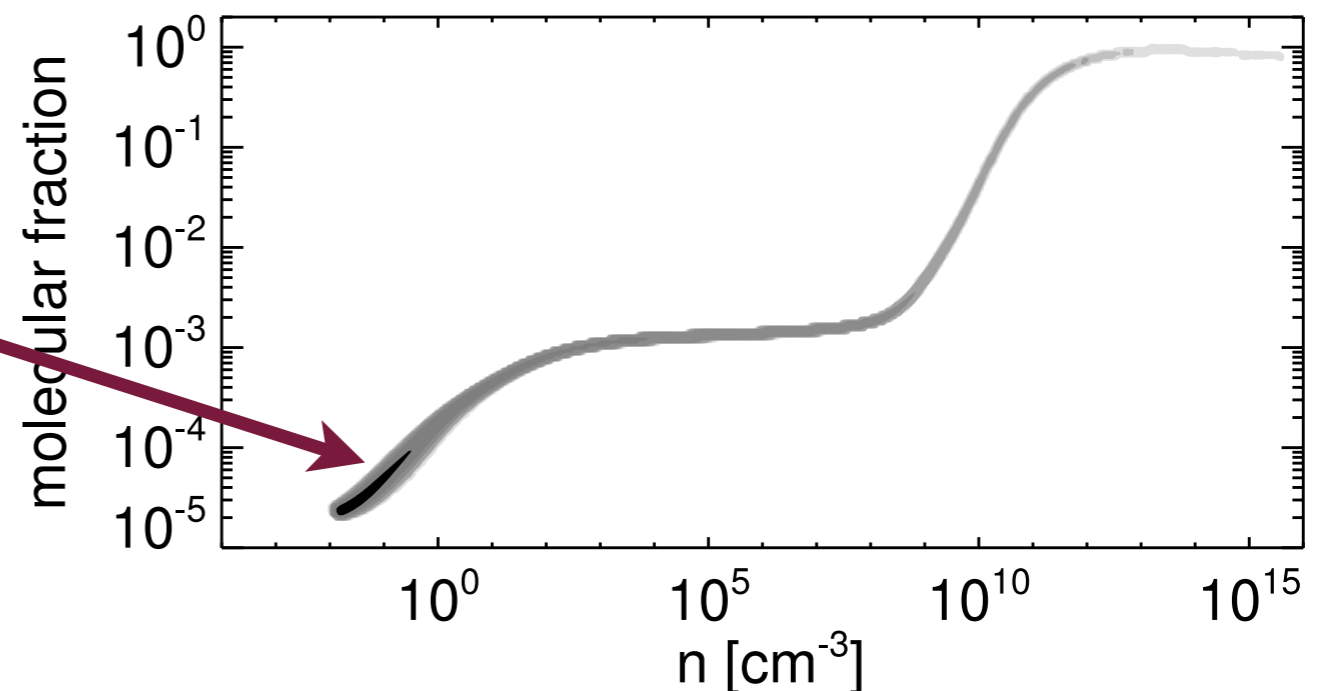
# What's happening at each of these stages?

Yoshida et al. (2006)

Gas falls into minihalo and shock heats to virial temperature



$\text{H}_2$  formation is enhanced by the temperature increase

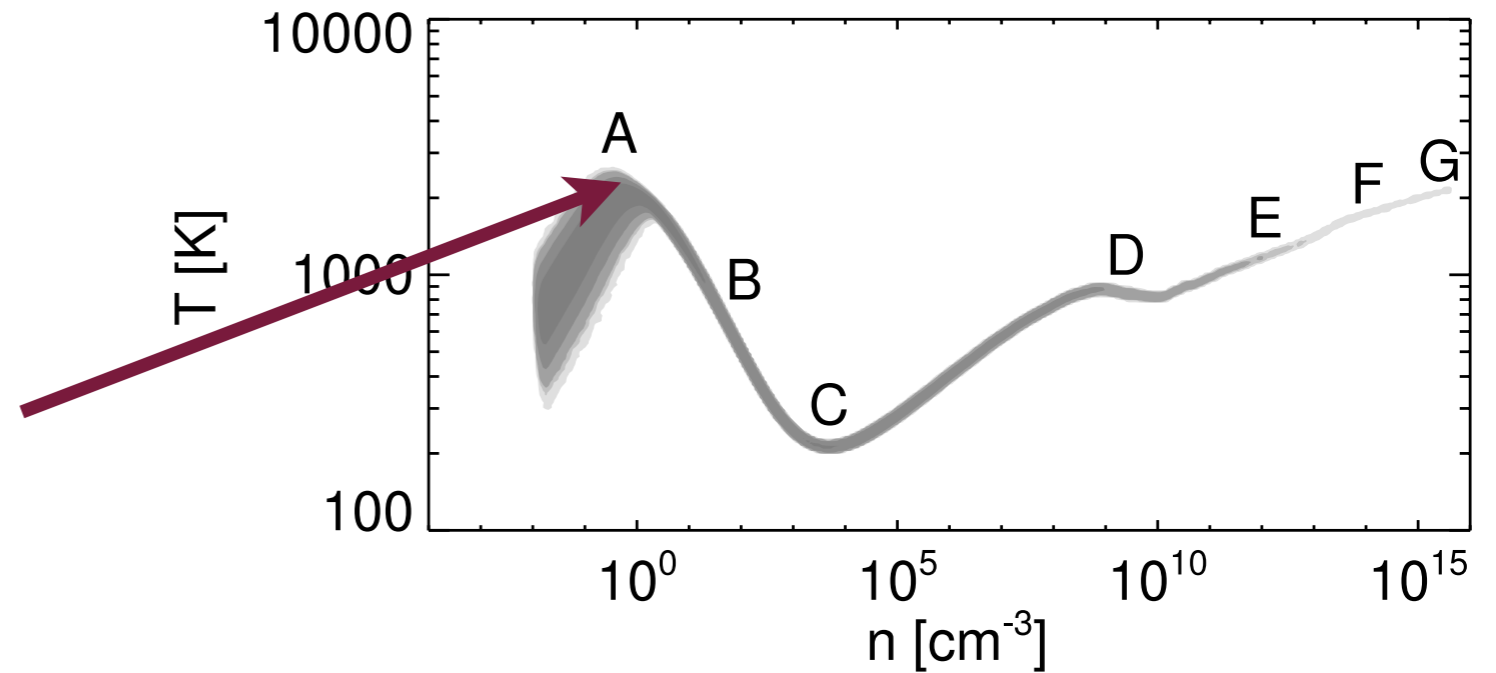




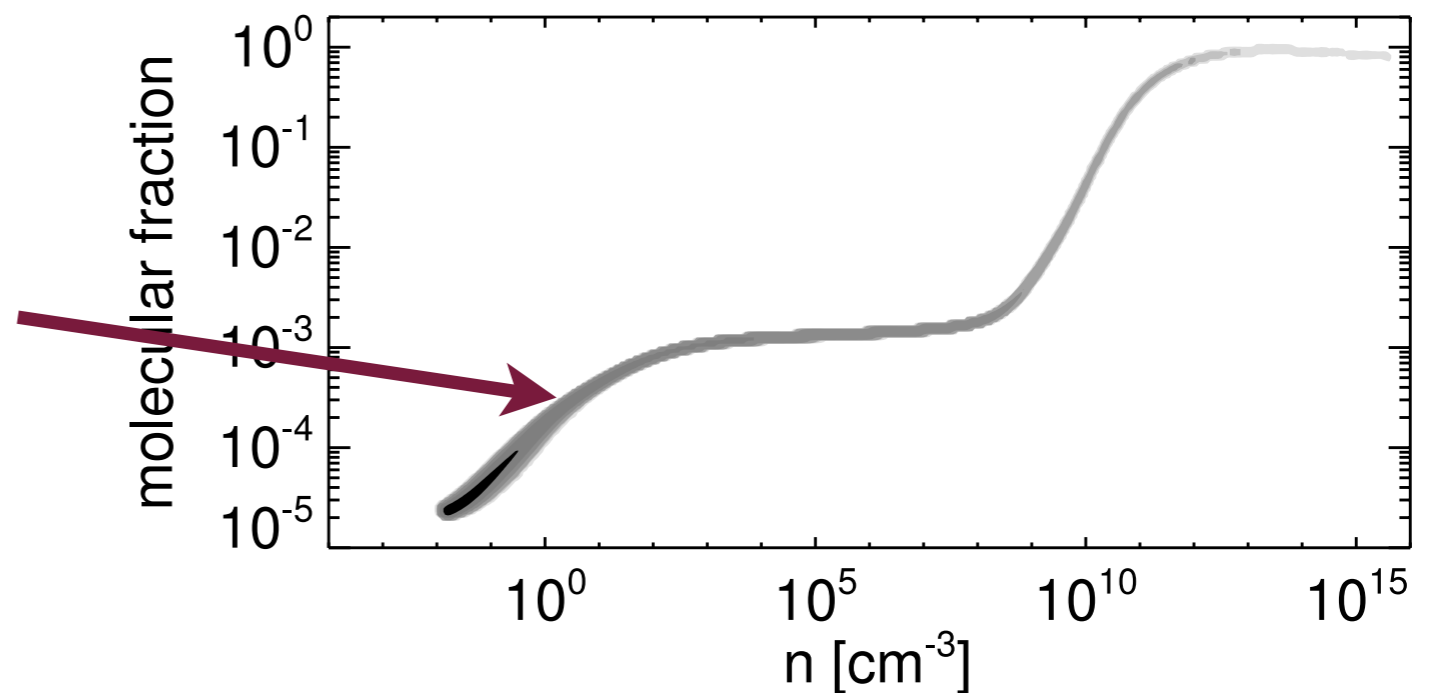
# What's happening at each of these stages?

Yoshida et al. (2006)

H<sub>2</sub> cooling counteracts the heating, and gas starts to cool down



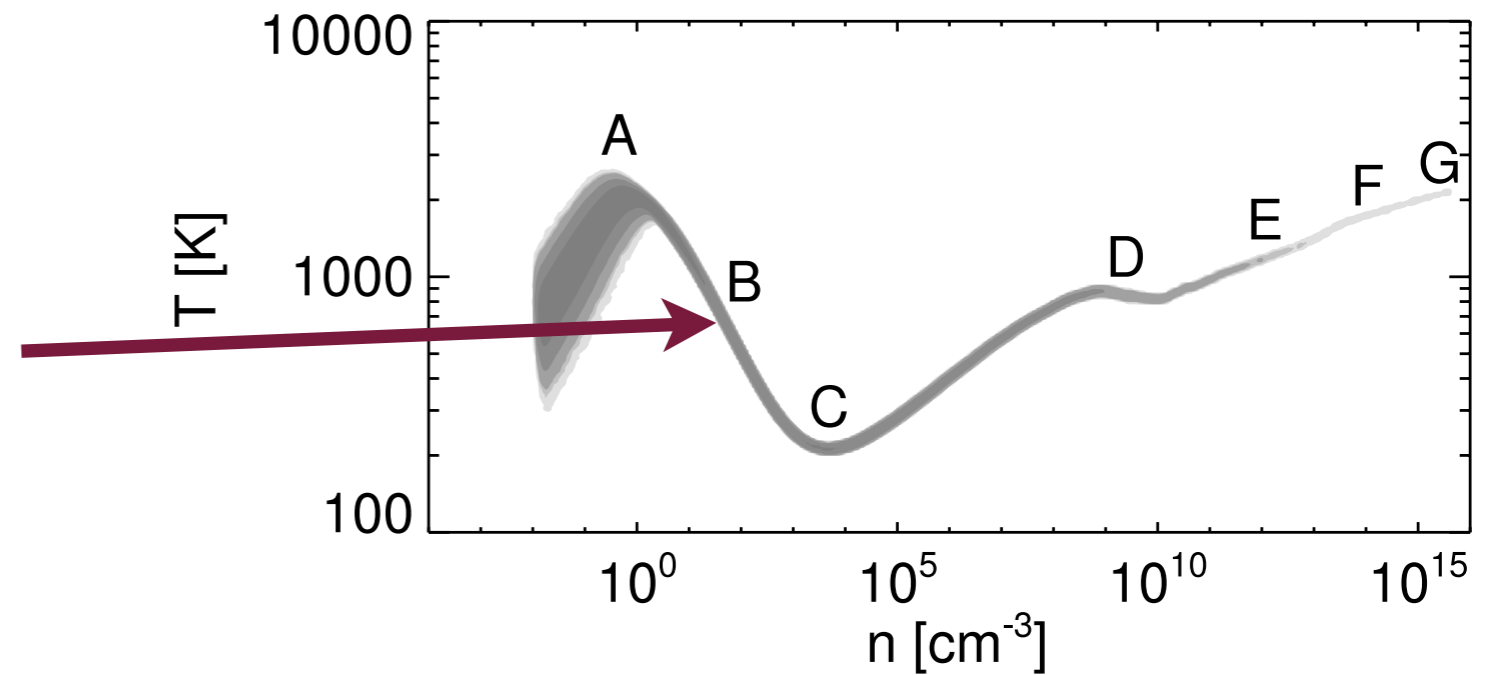
H<sub>2</sub> formation continues



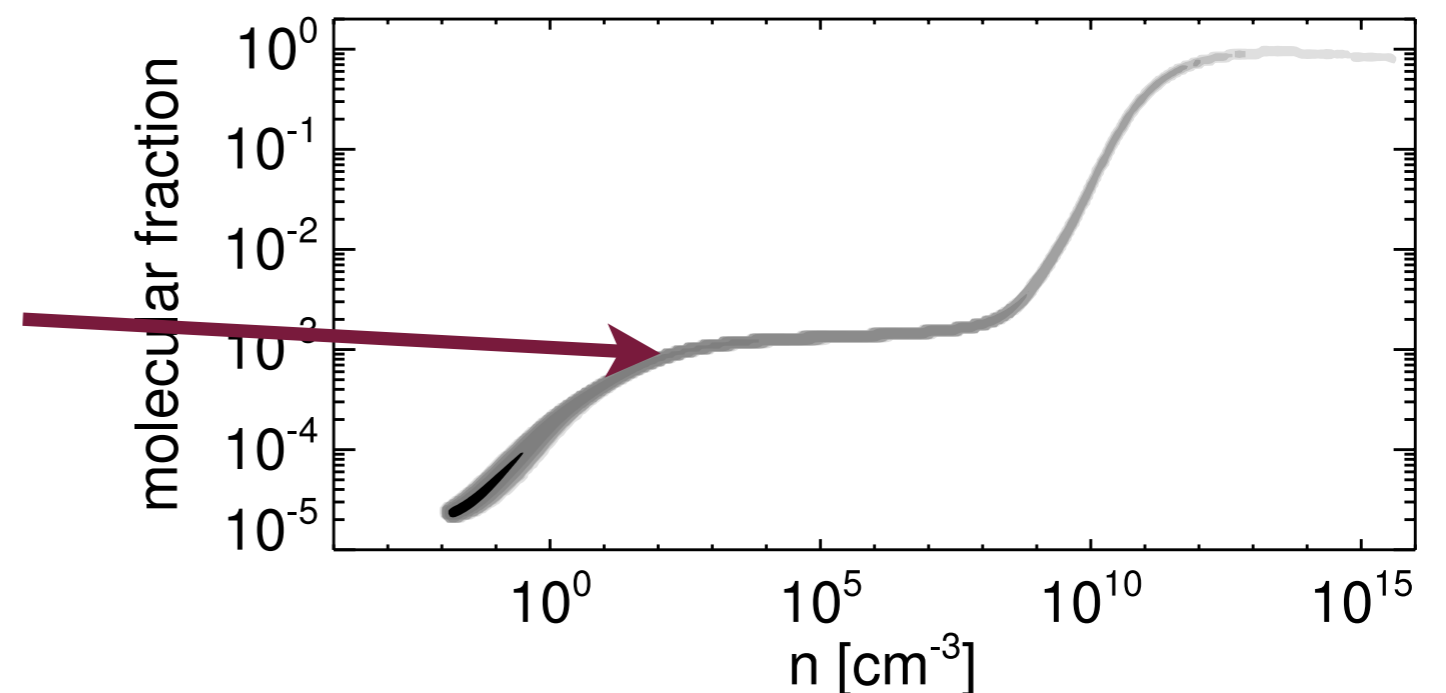
# What's happening at each of these stages?

Yoshida et al. (2006)

Cooling rate is faster than the heating rate (to do gravitational contraction). Density increases and  $T$  drops.



$\text{H}_2$  formation rate drops to around zero -- ionisation of gas is too low.



# What's happening at each of these stages?

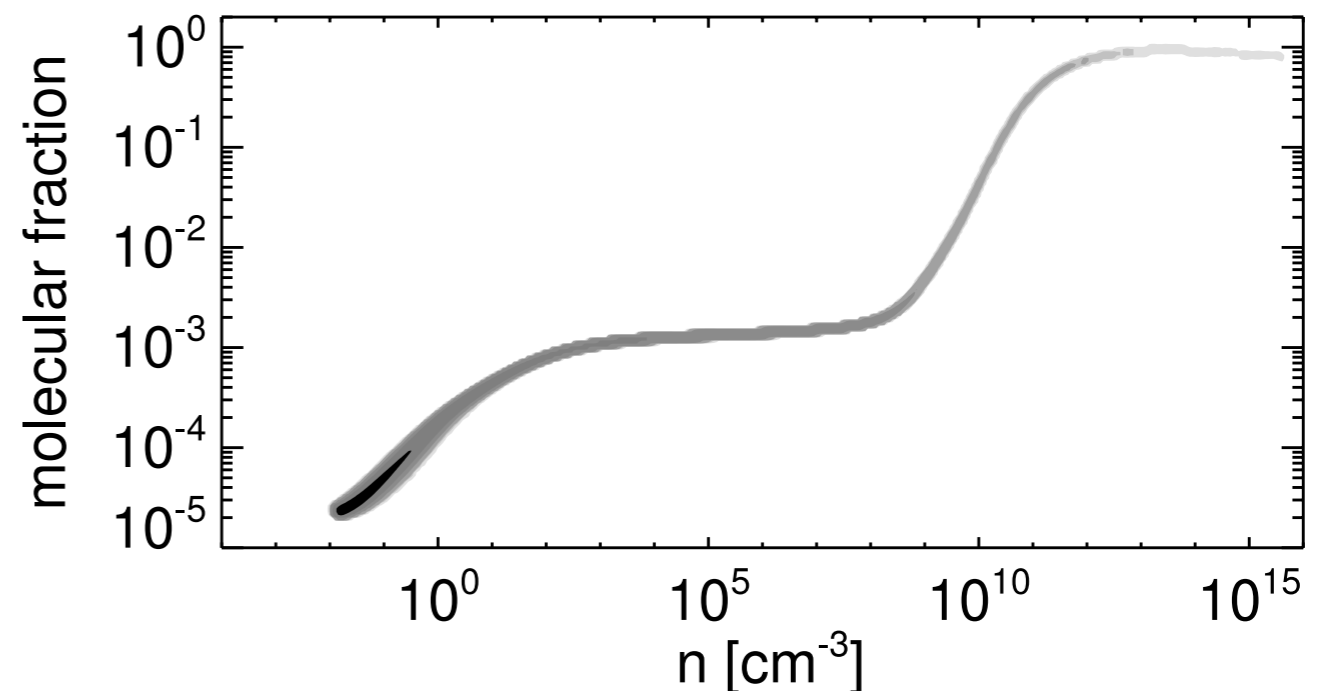
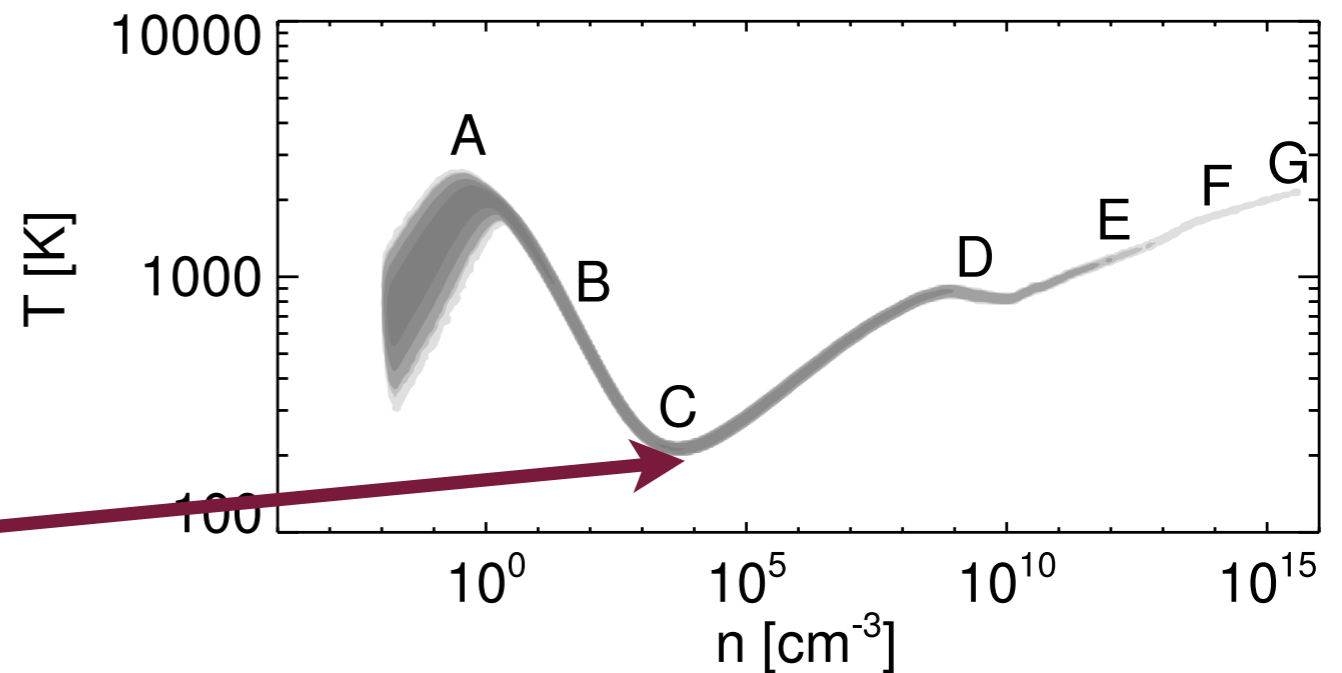
Yoshida et al. (2006)

H<sub>2</sub> is ineffective below around 200 K!

J = 2 rotational state sits at 512K

cooling rate drops as  $e^{-T/512K}$

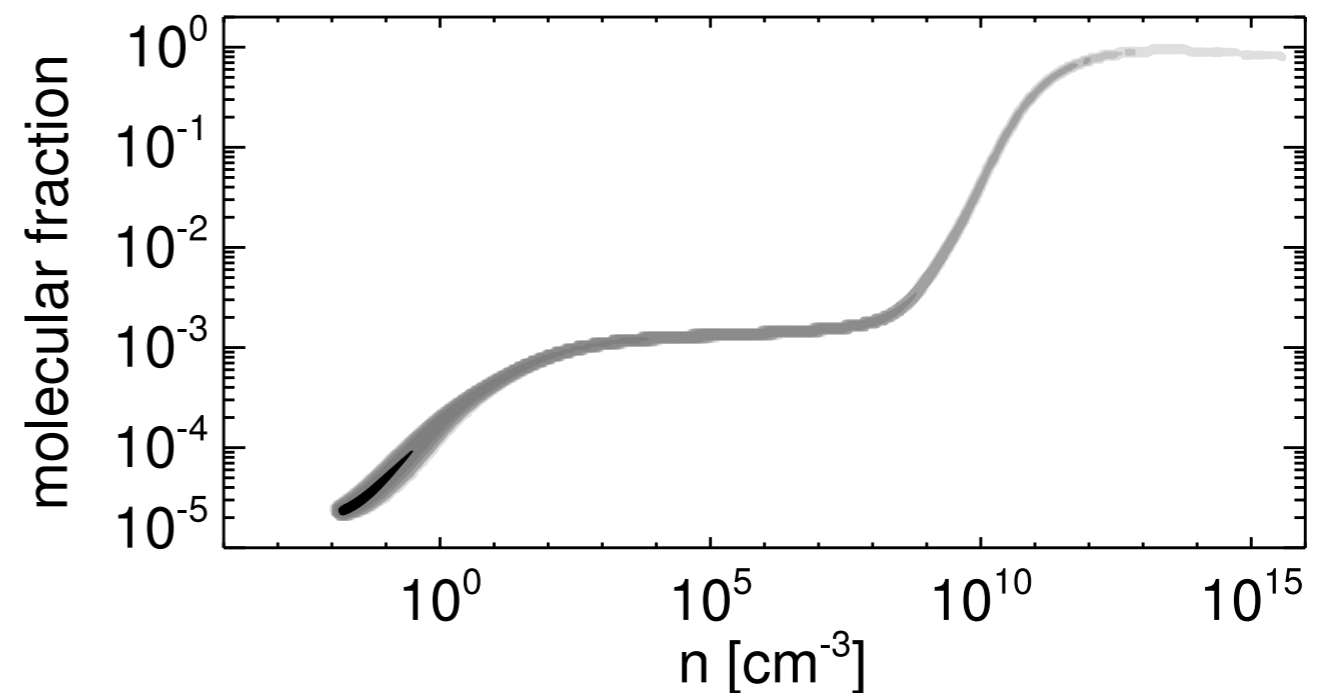
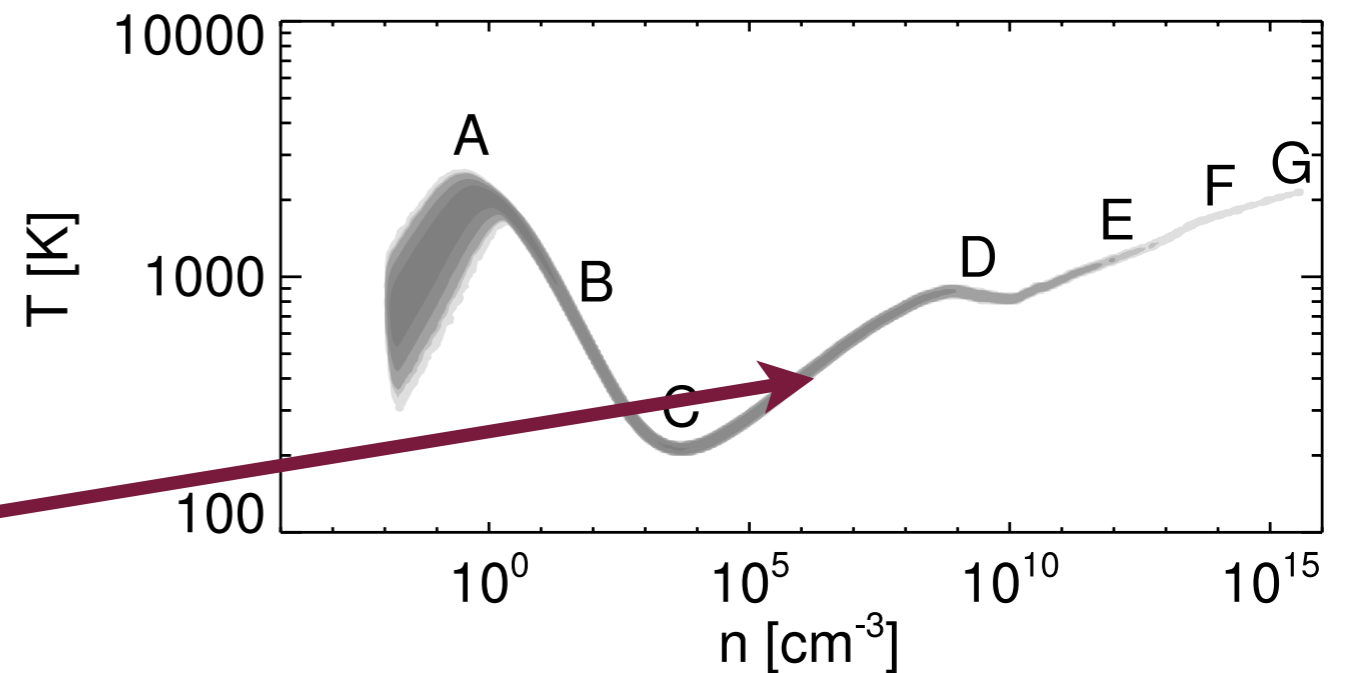
Limits minimum T to around few 100 K



# What's happening at each of these stages?

Yoshida et al. (2006)

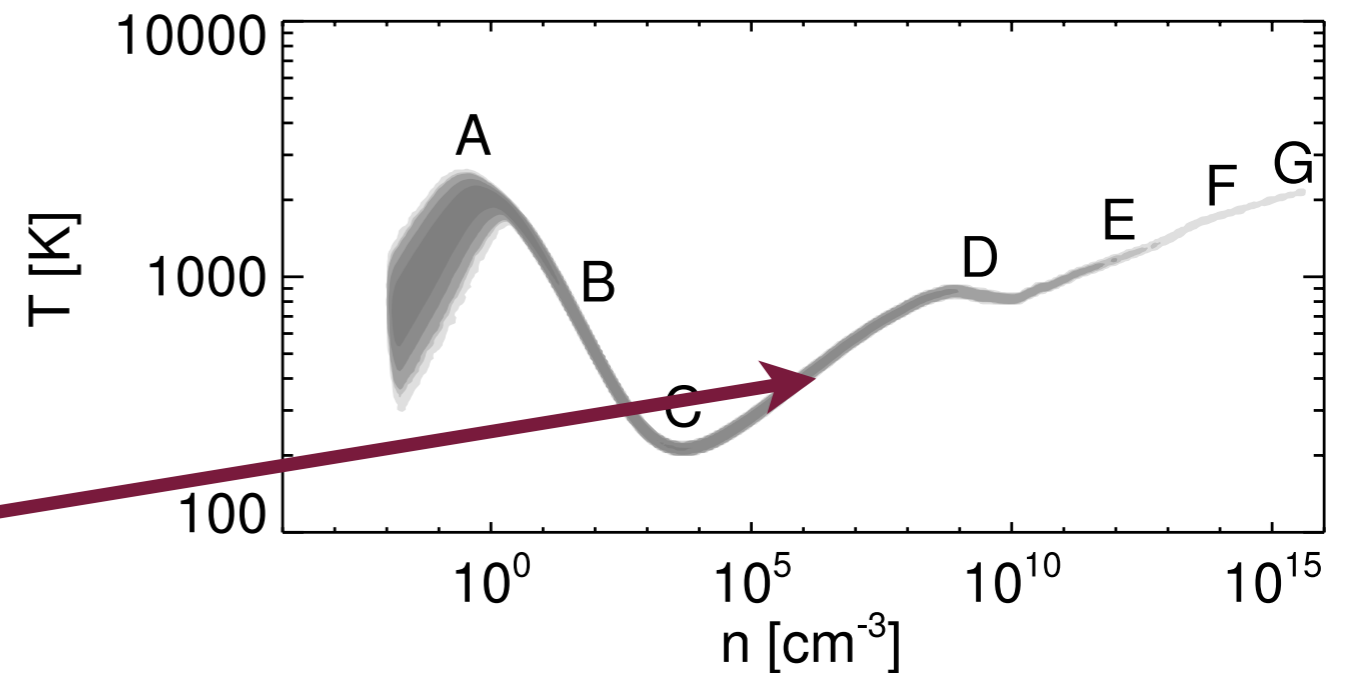
H<sub>2</sub> reaches its critical density, and levels go into LTE  
Cooling goes as  $n$ , rather than  $n^2$



# What's happening at each of these stages?

Yoshida et al. (2006)

H<sub>2</sub> reaches its critical density, and levels go into LTE  
Cooling goes as  $n$ , rather than  $n^2$



Compressional heating from gravitational collapse

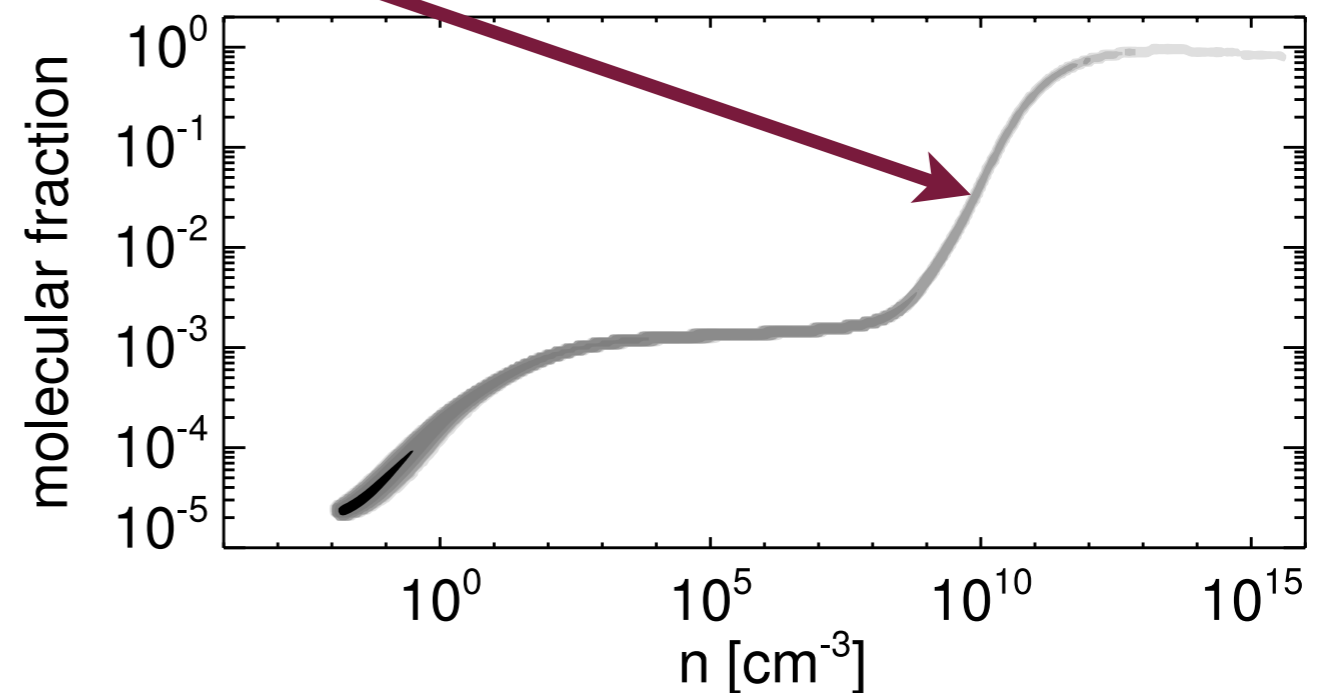
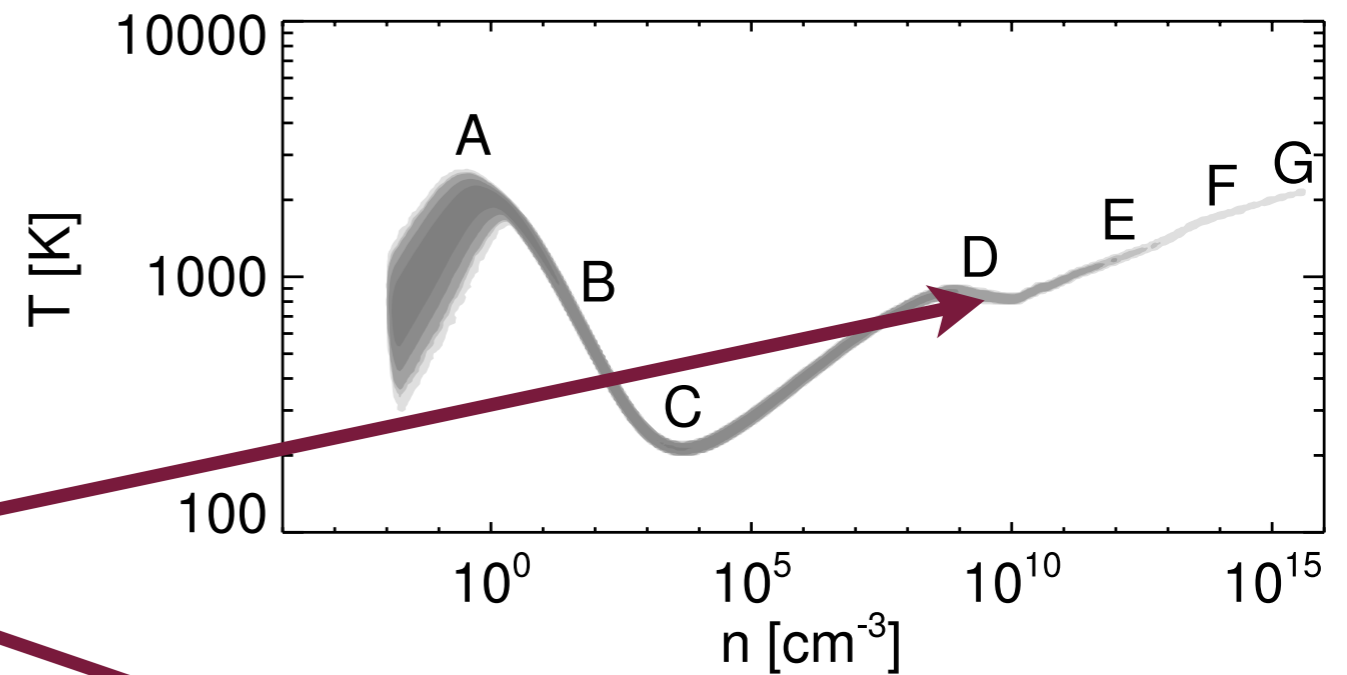
$$\frac{de}{dt} = \frac{kT}{\mu m_p} \left[ \frac{32G}{3\pi} \right]^{1/2} \rho^{3/2} - \Lambda_V + \Gamma_V$$

Gas has to heat up as it collapses

# What's happening at each of these stages?

Yoshida et al. (2006)

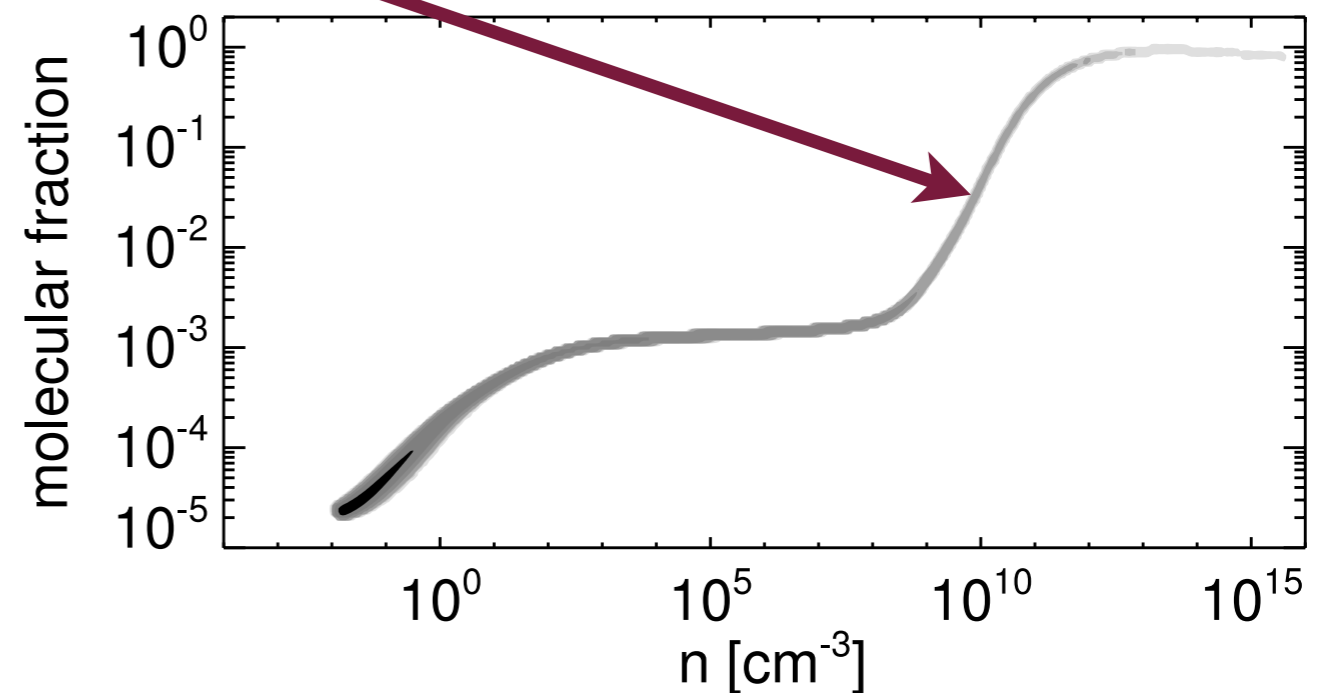
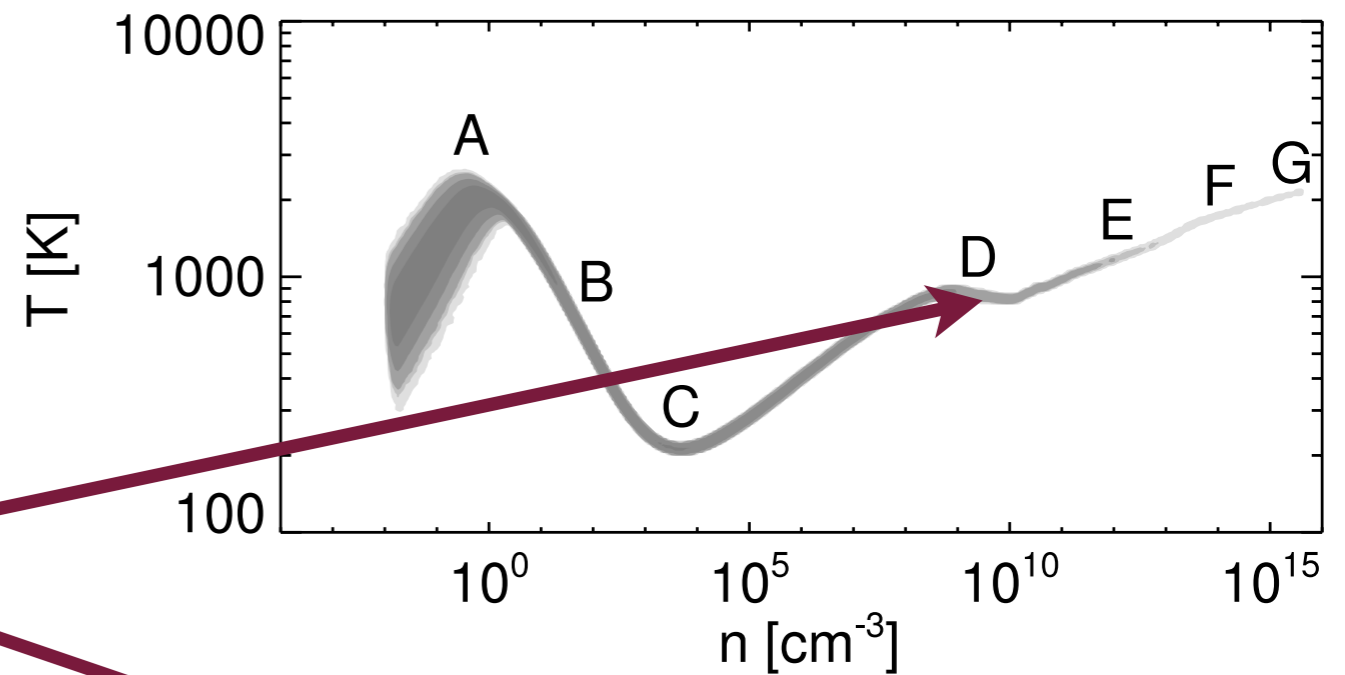
Rapid rise in the H<sub>2</sub> fraction due to 3-body reactions. Results in cooling once H<sub>2</sub> fraction is high



# H<sub>2</sub> becomes optically thick!

Yoshida et al. (2006)

Rapid rise in the H<sub>2</sub> fraction  
due to 3-body reactions.  
Results in cooling once H<sub>2</sub>  
fraction is high



# The 3-body formation regime

- Two path ways:



- Consider the first. Using the rate from Glover (2008):

$$\frac{dn_{\text{H}_2}}{dt} = k_{3b} n_{\text{H}}^3 \quad k_{3b} = 7.7 \times 10^{-31} T^{-0.464} \text{ cm}^6 \text{ s}^{-1}$$

$$\Gamma_{3b} = 4.4 \text{ eV} \frac{dn_{\text{H}_2}}{dt} \quad t_{\text{H}_2} \sim \frac{n_{\text{H}_2}}{k_{3b} n_{\text{H}}^3}$$

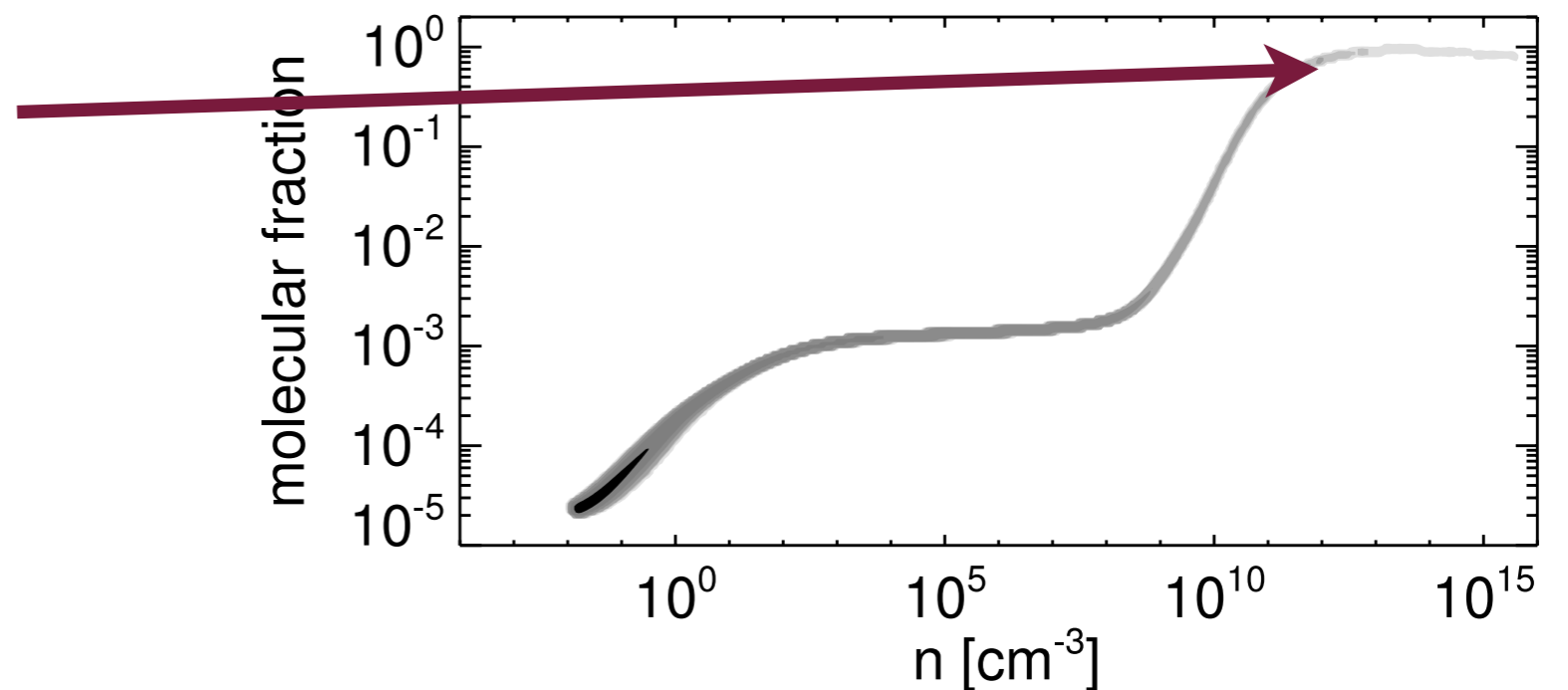
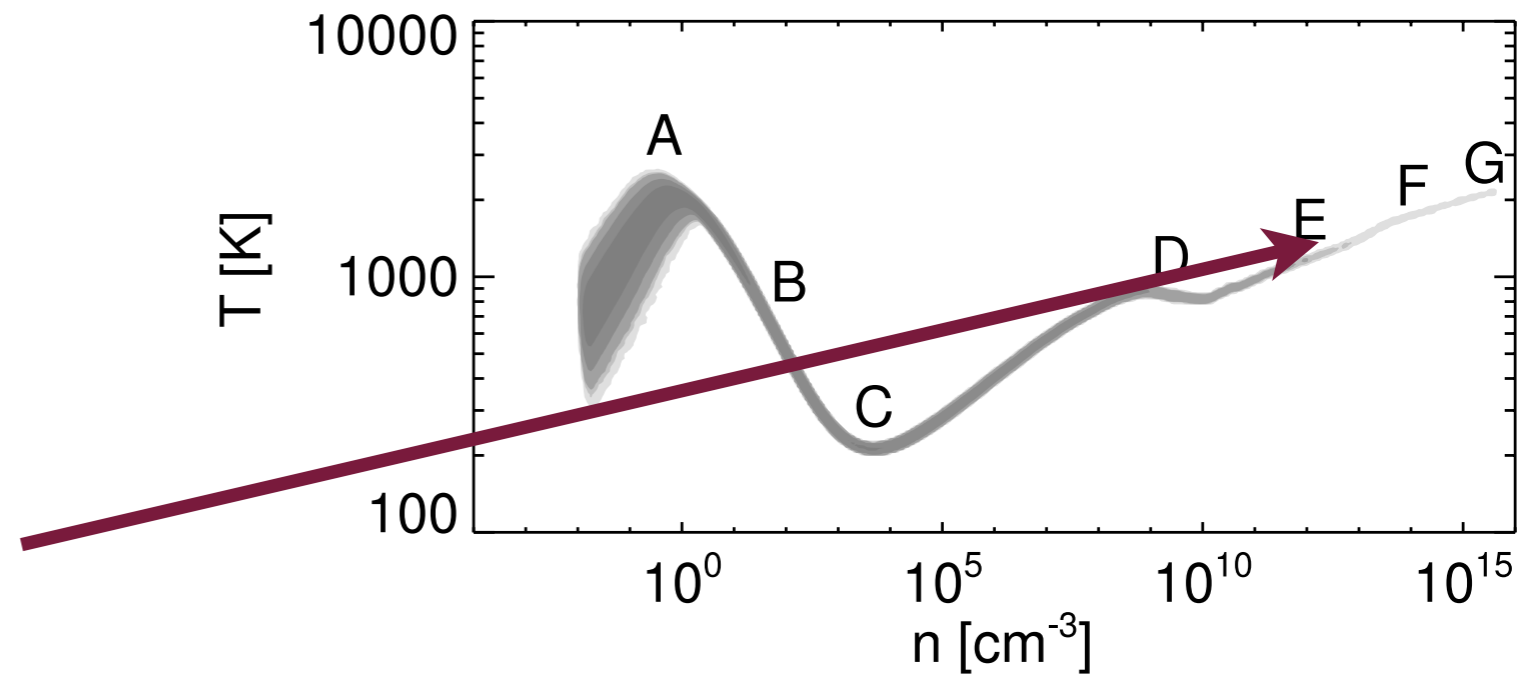


# H<sub>2</sub> becomes optically thick!

Yoshida et al. (2006)

H<sub>2</sub> cooling becomes optically thick.

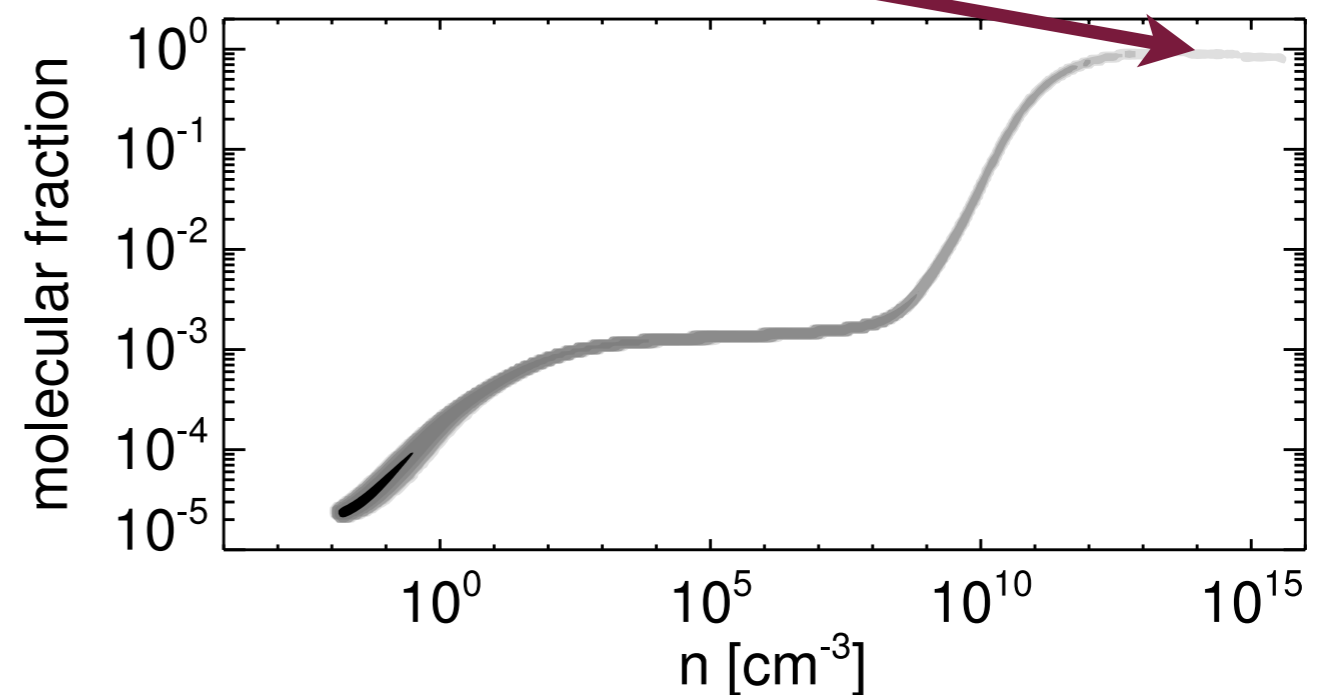
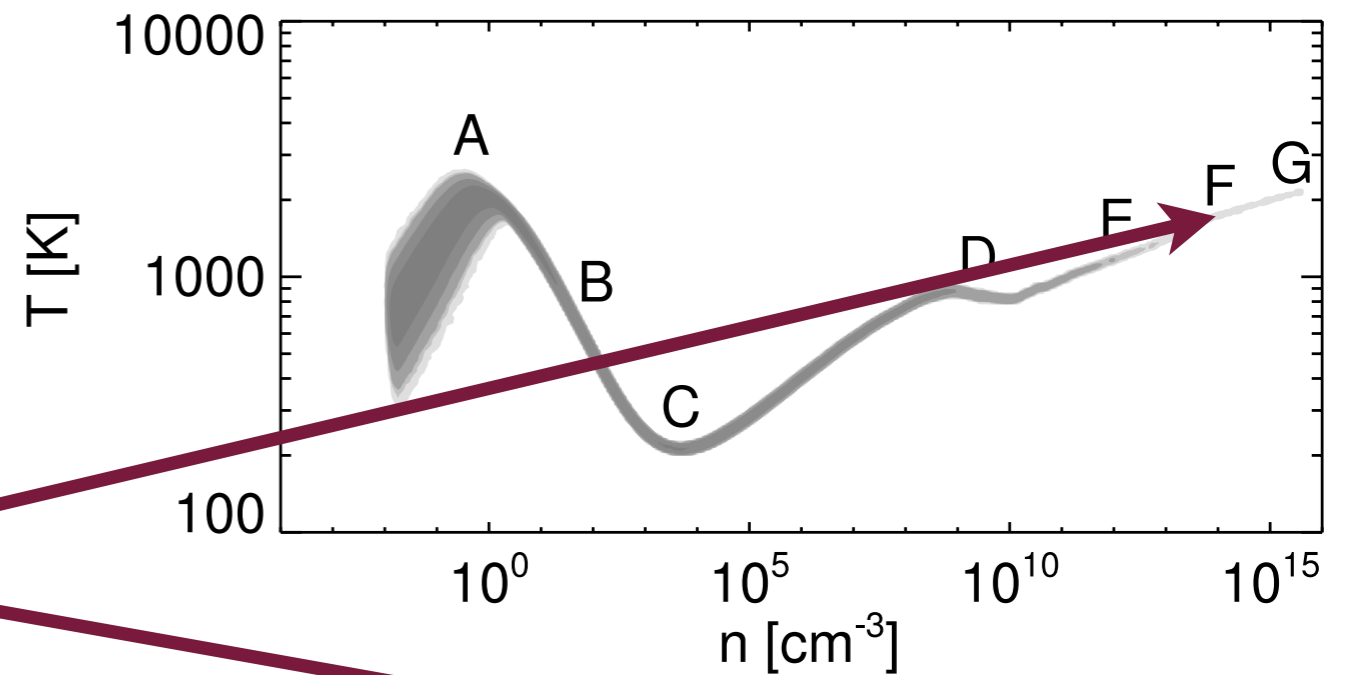
Reason:  $n_{\text{H}_2}$  is very high!



# H<sub>2</sub> becomes optically thick!

Yoshida et al. (2006)

Collision induced  
emission (CIE)  
cooling takes over

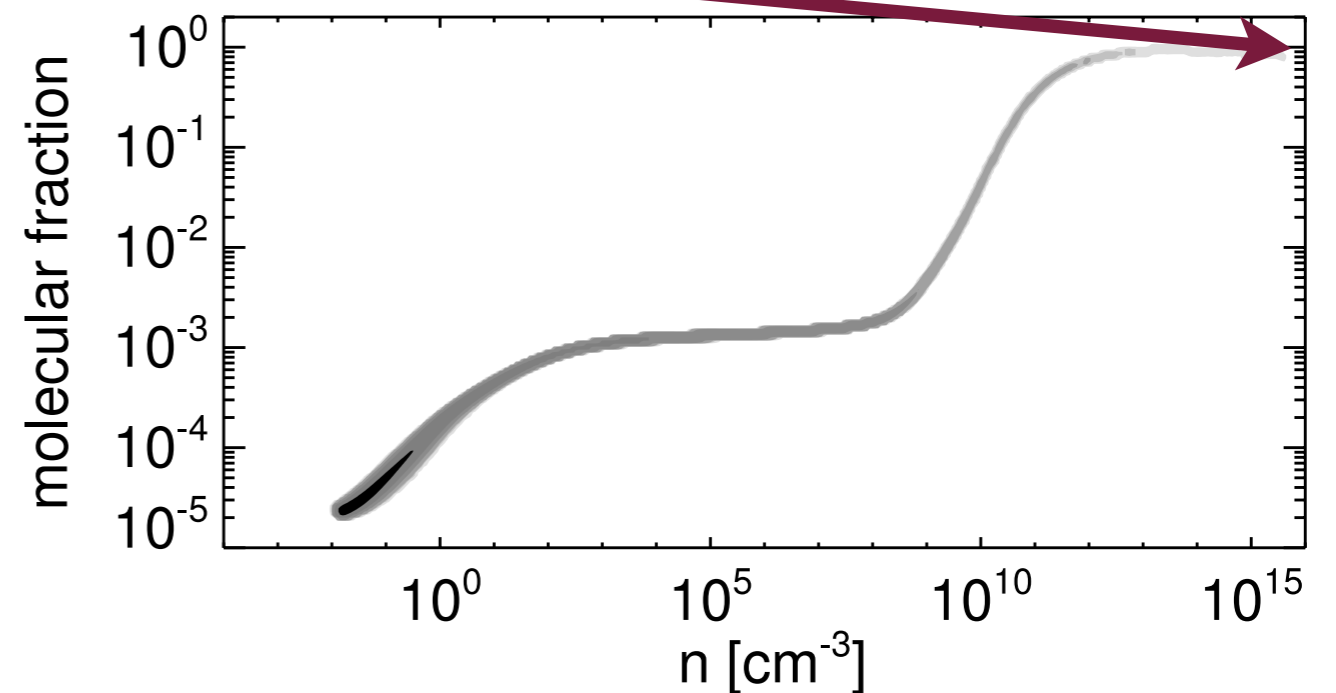
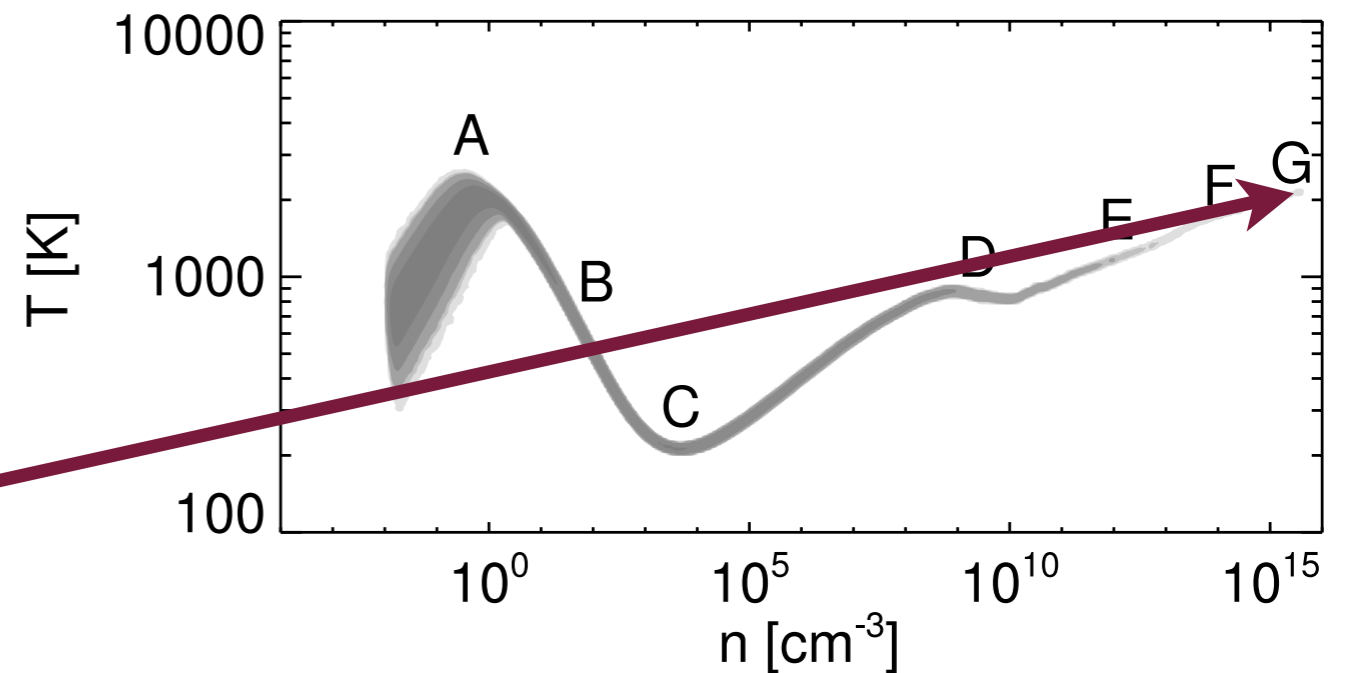


# H<sub>2</sub> becomes optically thick!

Yoshida et al. (2006)

CIE cooling becomes optically thick

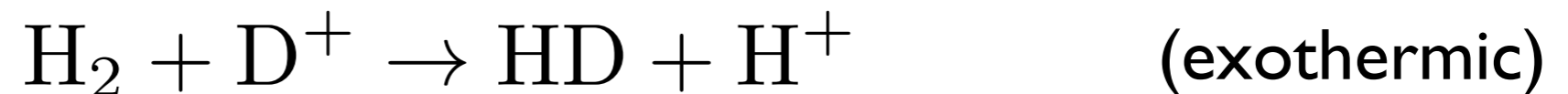
Gas cools via collisional dissociation (recover 4.4eV per molecule)



# Different path to Pop III?

## HD cooling:

- Formed and destroyed via the reactions,



- Below 462K, reaction 1 is favoured over reaction 2.
- Results in chemical fractionation

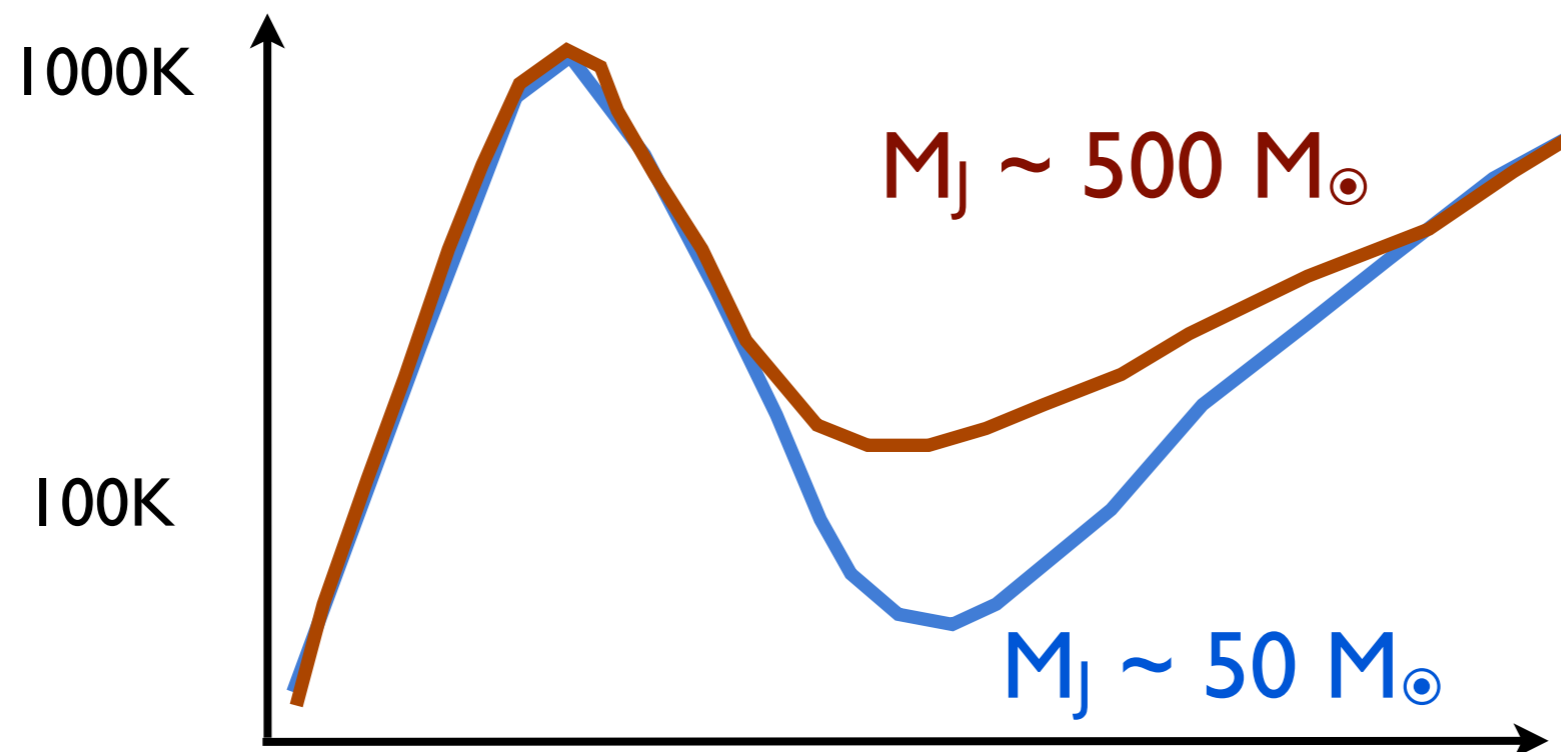
$$\frac{x_{\text{HD}}}{x_{\text{H}_2}} > \frac{x_{\text{D}}}{x_{\text{H}}} \quad (x_{\text{D}} = 2.6 \times 10^{-5})$$

# Different path to Pop III?

## HD cooling:

- HD can cool the gas down to the CMB temperature,

$$T_{\text{CMB}} = 2.728 \text{ K} (1 + z)$$

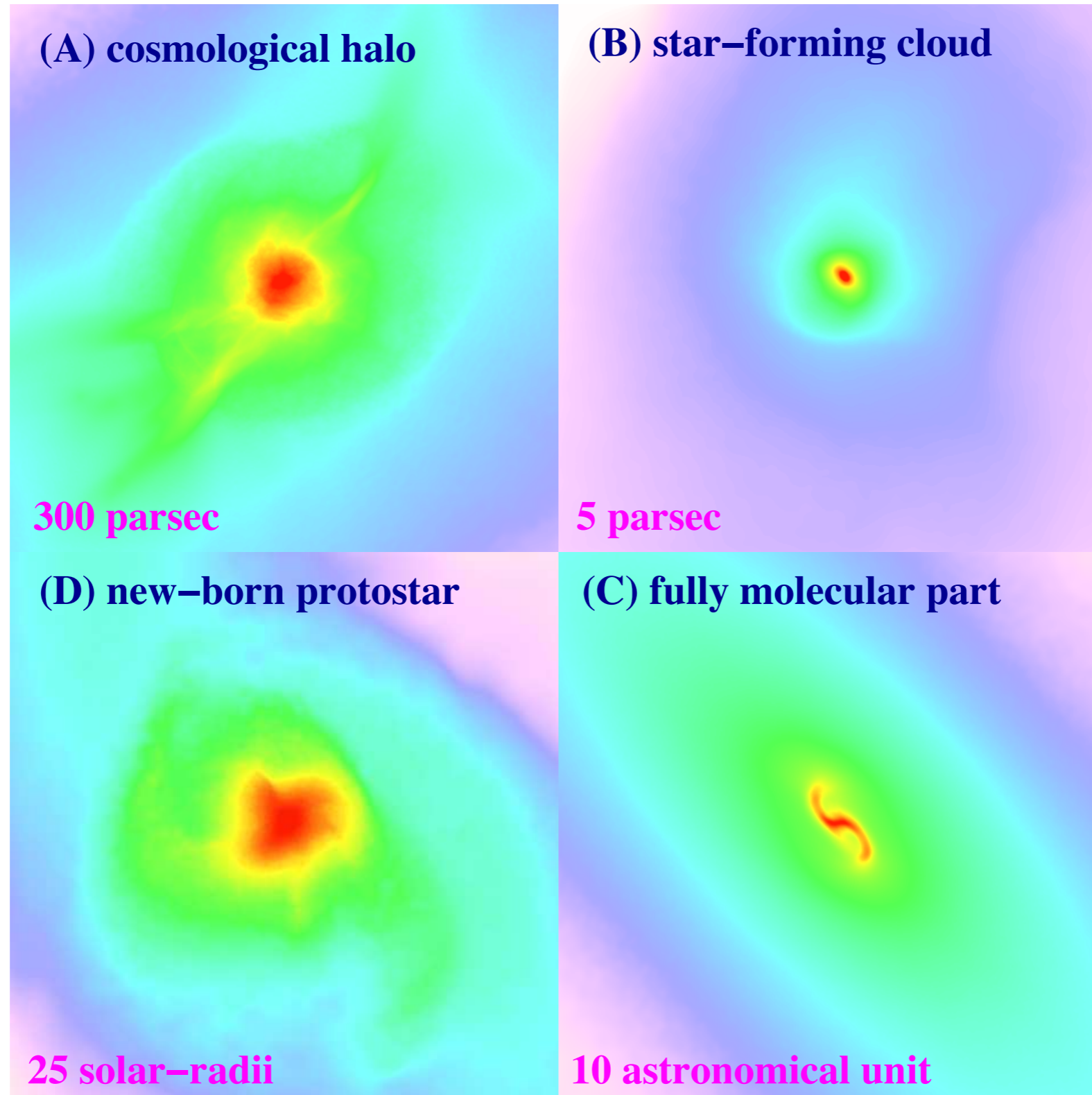


# Different path to Pop III?

## How to excite the HD cooling channel?

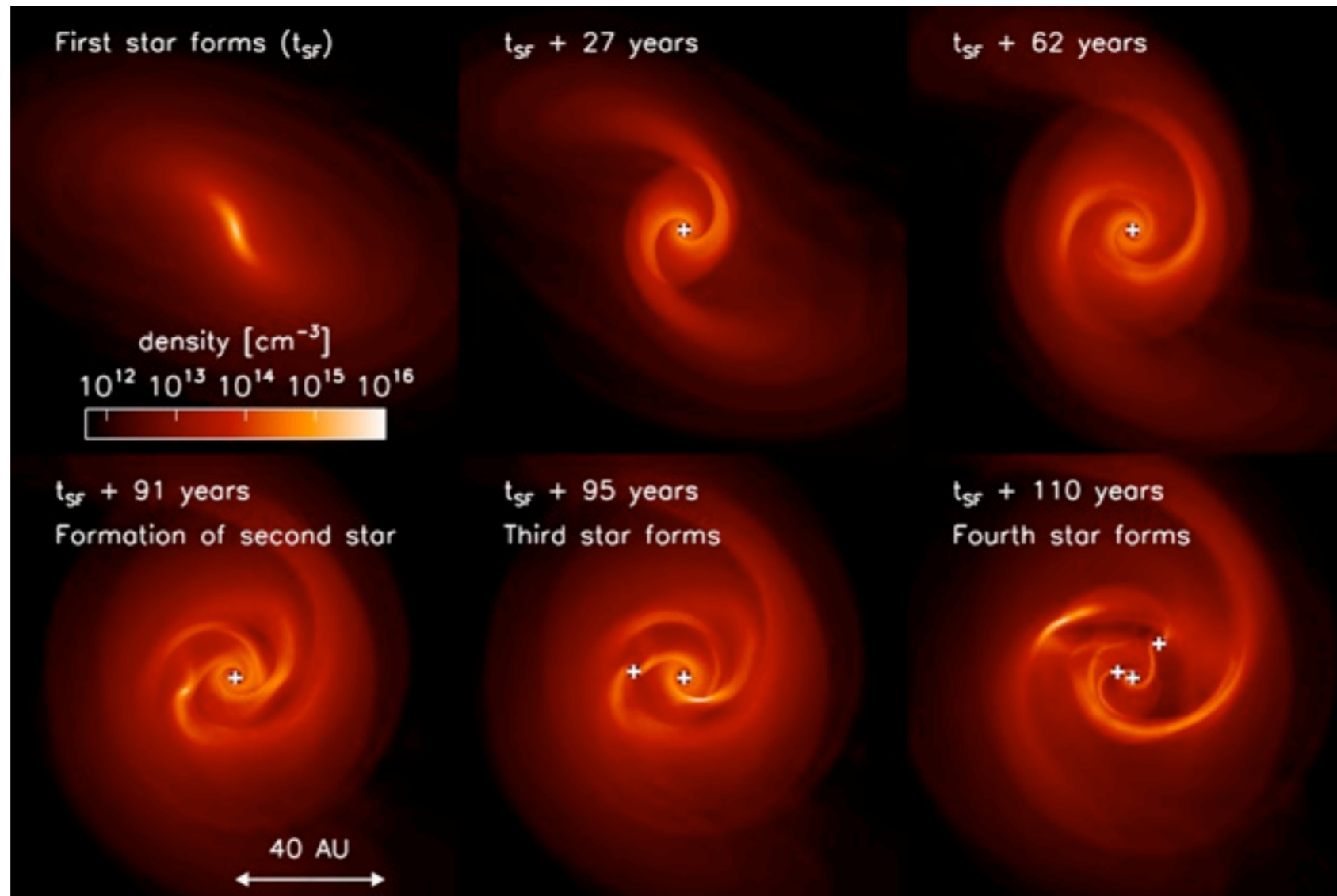
- Recall that the H<sub>2</sub> fraction depends on the ionisation fraction in the minihalo.
- Halos with higher ionisation fractions can form more H<sub>2</sub>, and cool to slightly lower temperatures.
- Can then produce the HD fractionation, and the HD cooling.
- Can be triggered by nearby star formation or via collisional ionisation in slightly more massive halos.

# What happens next?



Yoshida, Omukai, Hernquist (2008)

# What happens at high densities?



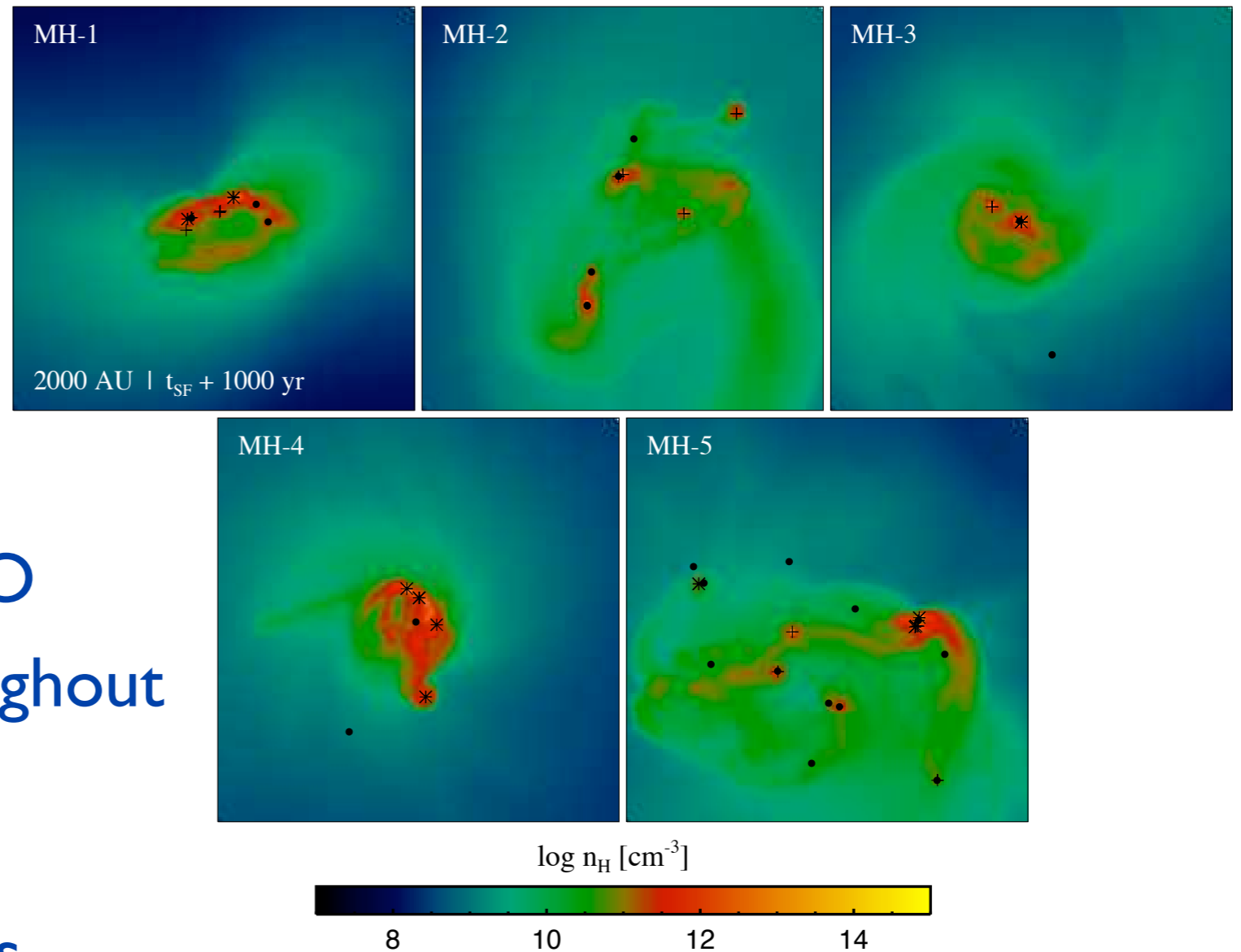
Clark et al. (2011b)

- Disc builds up which fragments to form a small N system
- Central protostar mass  $\sim 0.5 M_{\odot}$
- Fragmentation timescale  $\sim 100$  yr



# Complimentary study

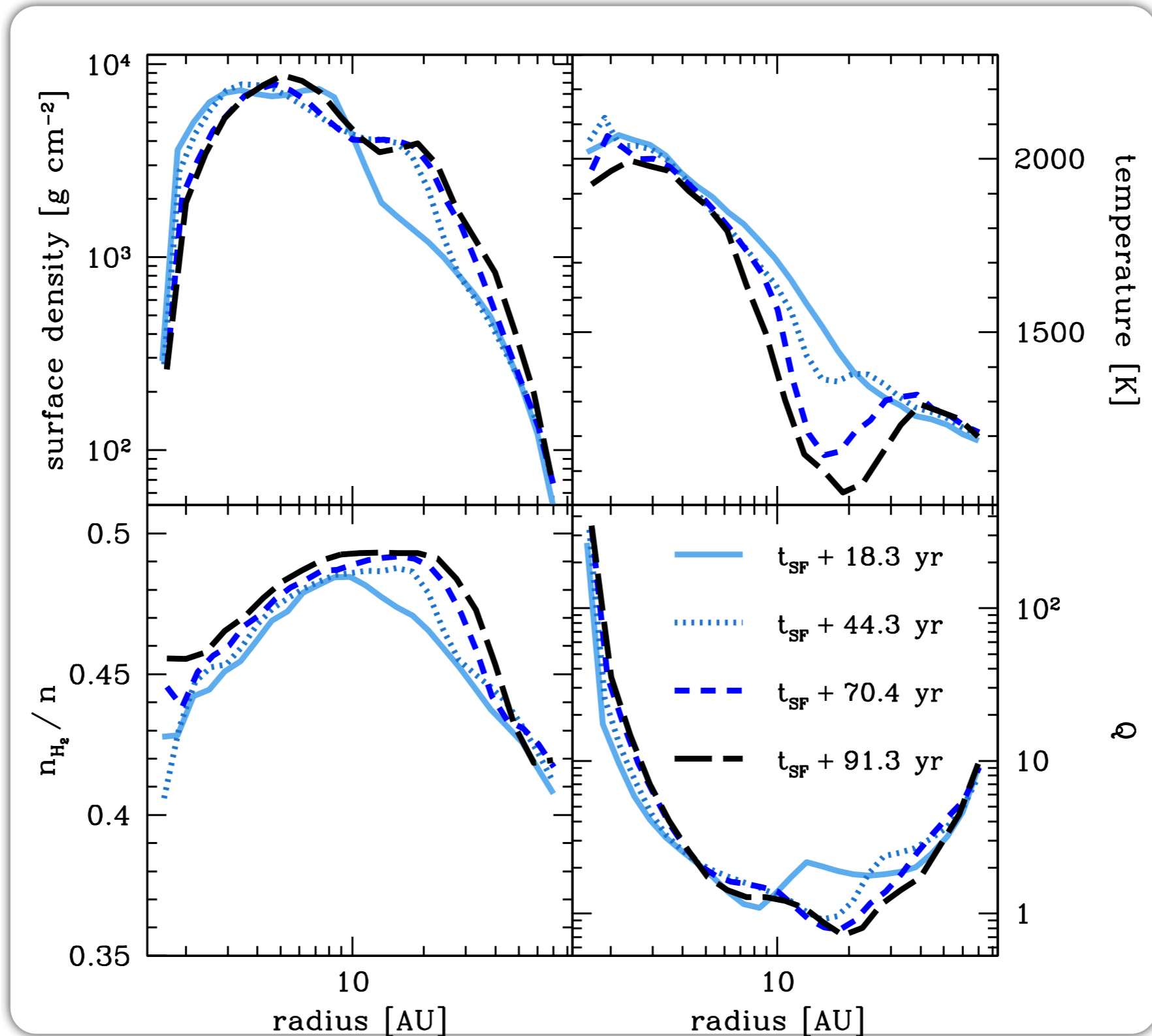
Always good to use a different code!



- Greif et al. 2011 used AREPO
- Had higher resolution throughout collapse.
- Better able to resolve the turbulence in the infalling gas.
- Same basic picture, but much more messy.

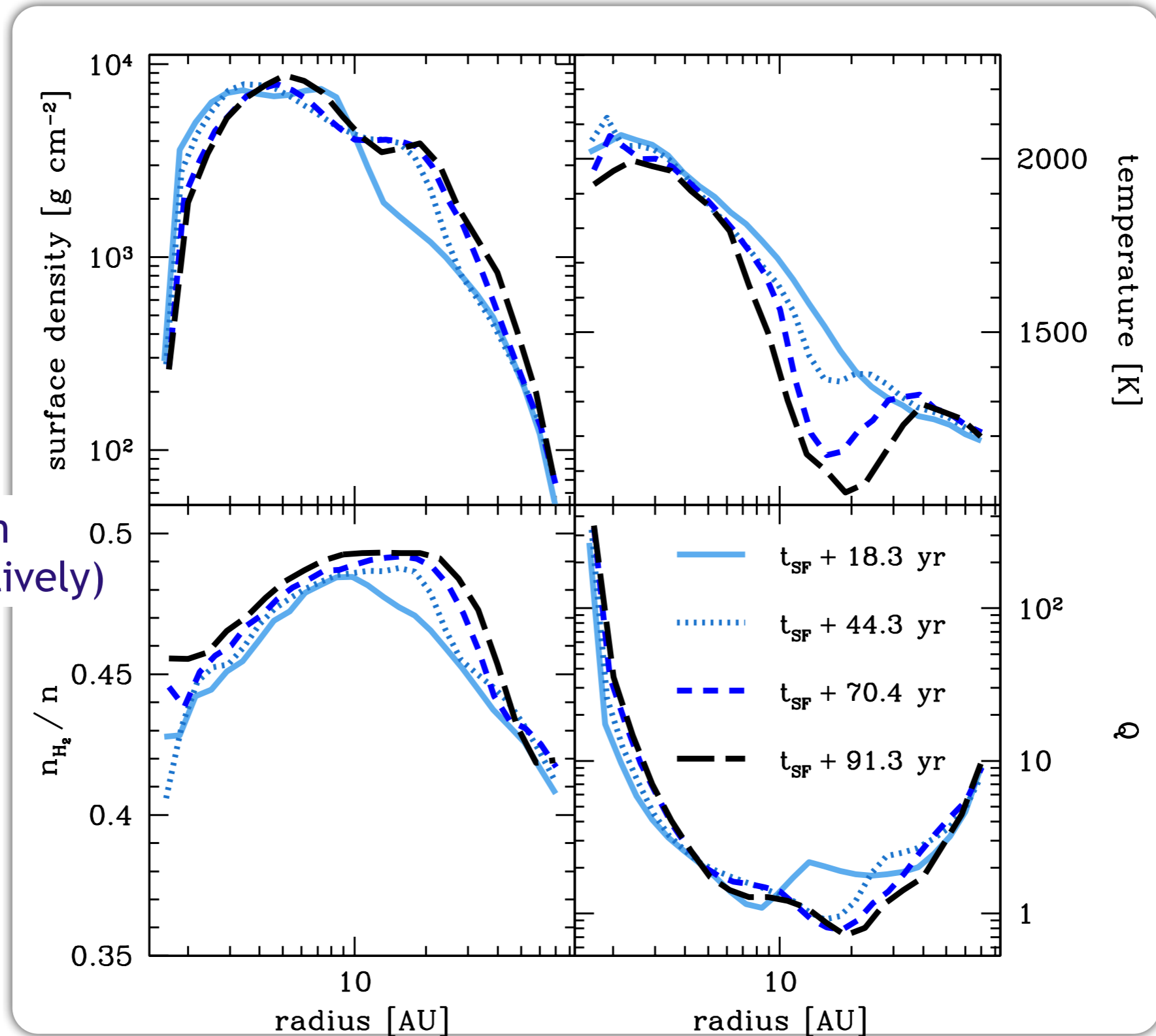
Greif et al. (2011) + Smith et al. (2011)

# Why does the gas fragment?



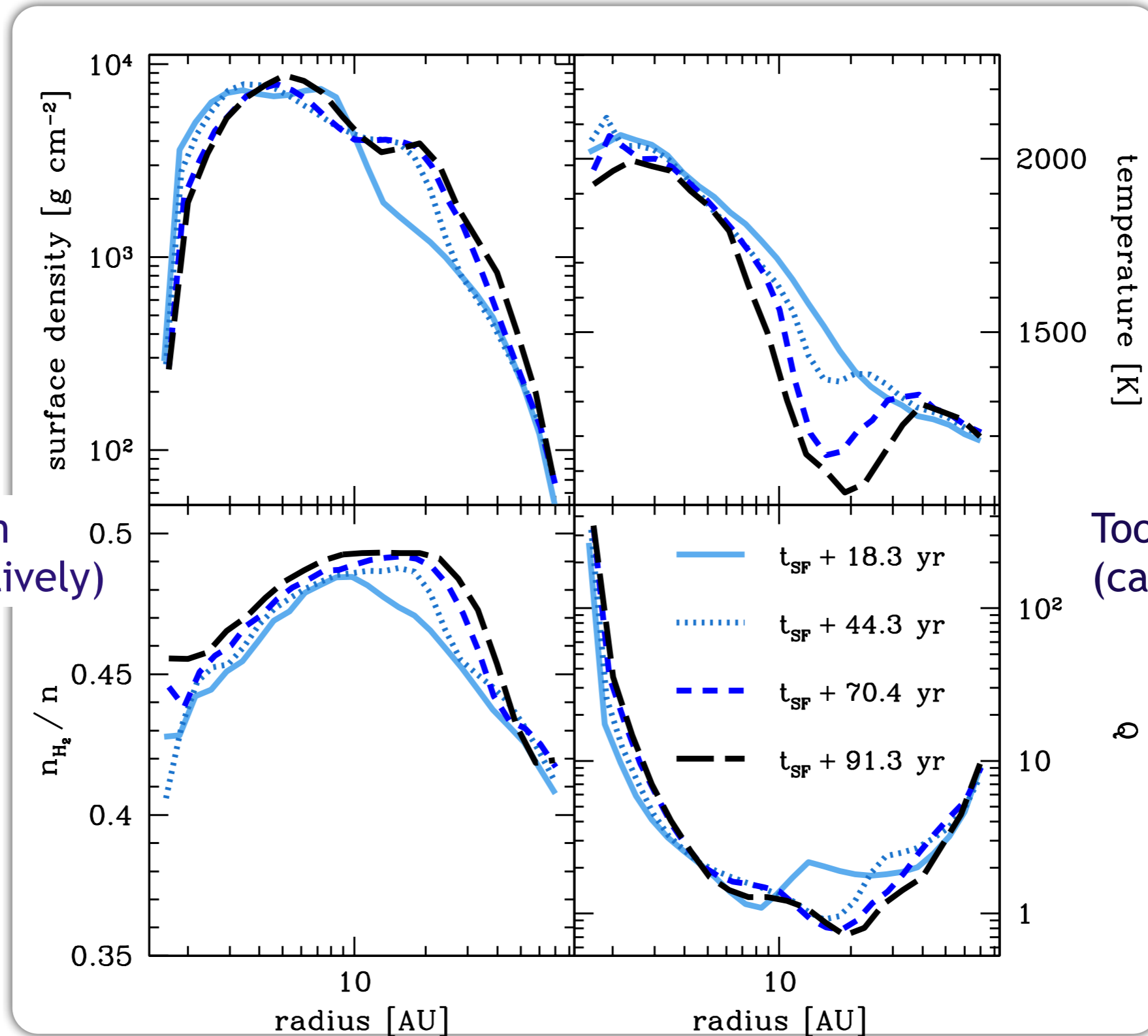
Clark et al. (2011b)

# Why does the gas fragment?



High  $\text{H}_2$  fraction  
(can cool effectively)

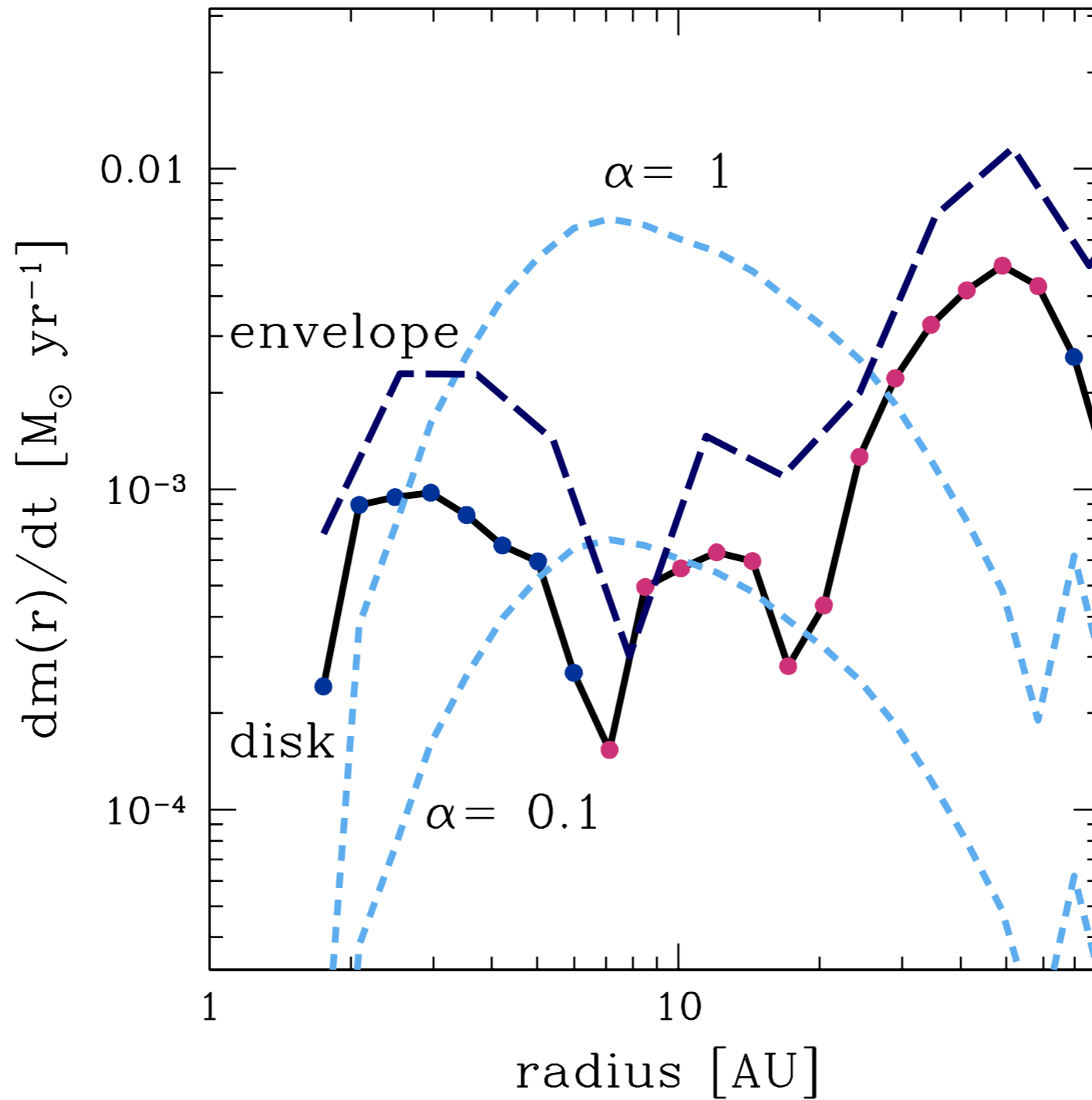
# Why does the gas fragment?



High H<sub>2</sub> fraction  
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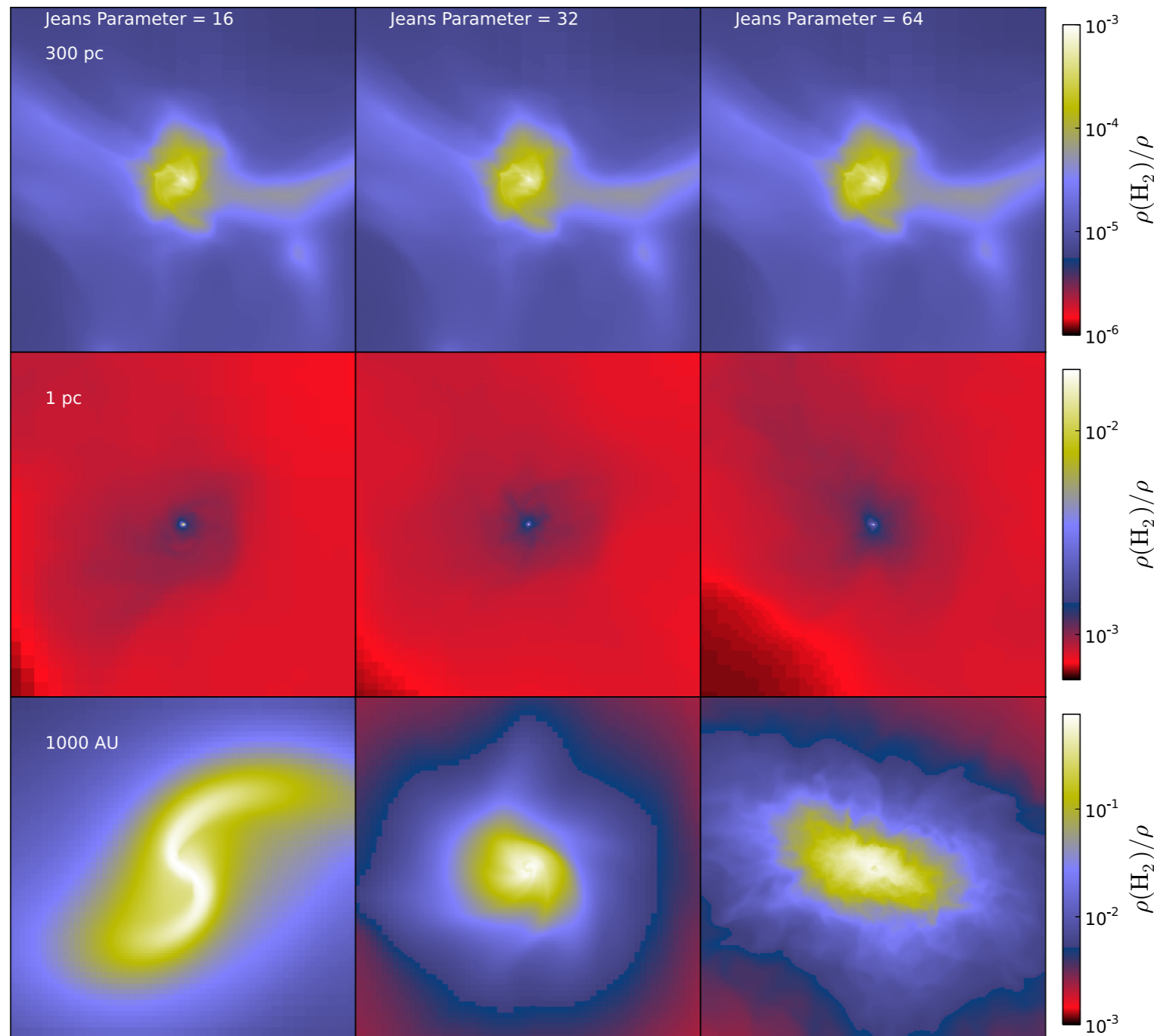
Toomré unstable  
(can fragment)

# Why does the gas fragment?



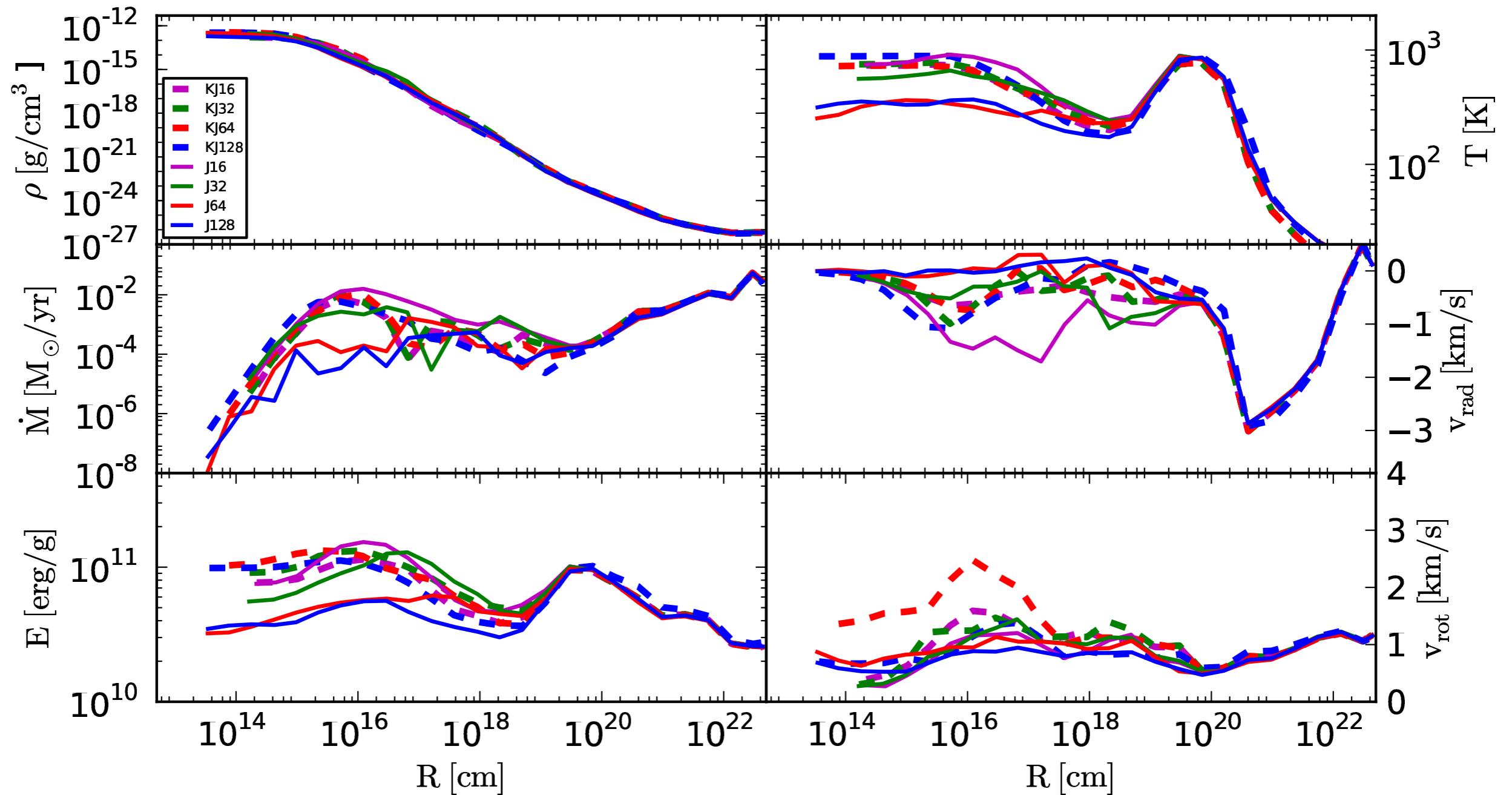
# Problems with convergence?

- Should we form a large, gravitationally unstable disc?



# Problems with convergence

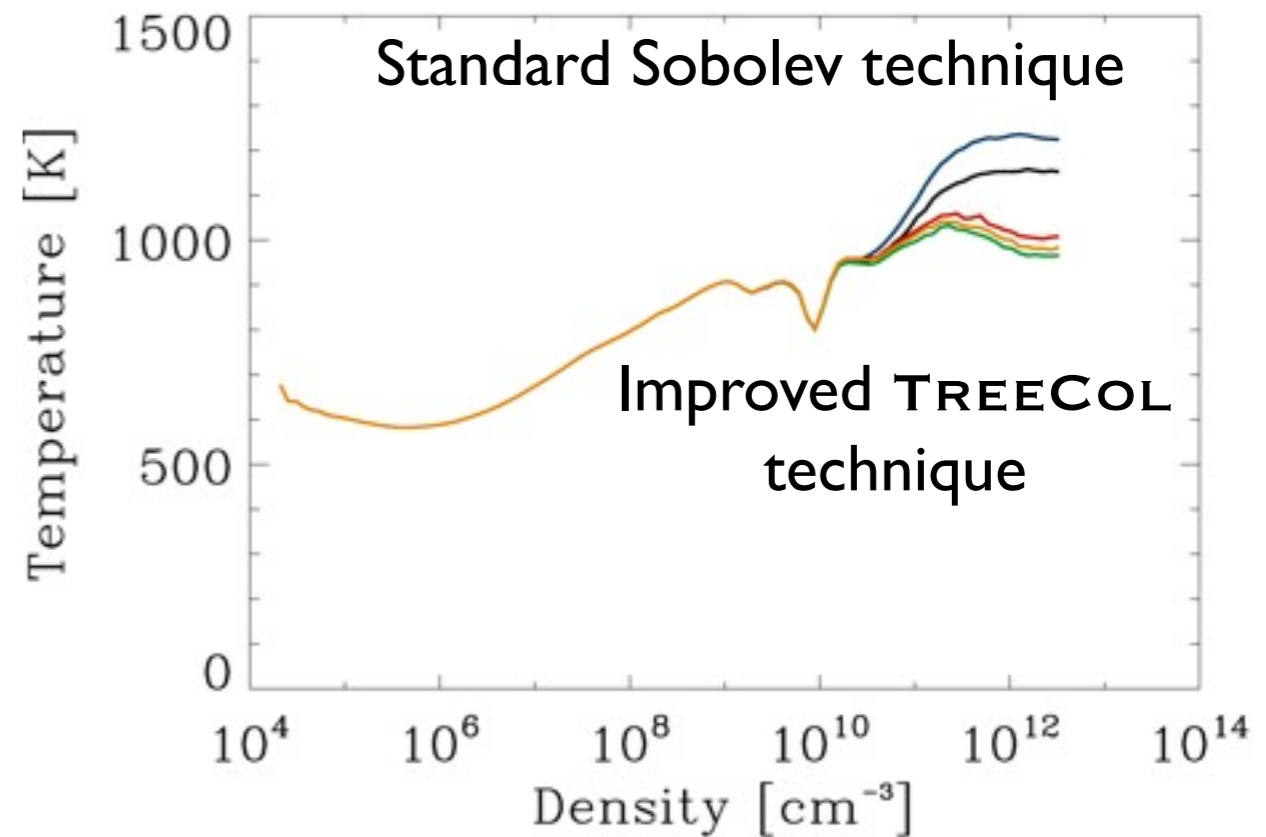
- Need to be careful of how the chemistry/cooling is solved!



# Clearly H<sub>2</sub> line-cooling is important

Tilman Hartwig's Masters Thesis:

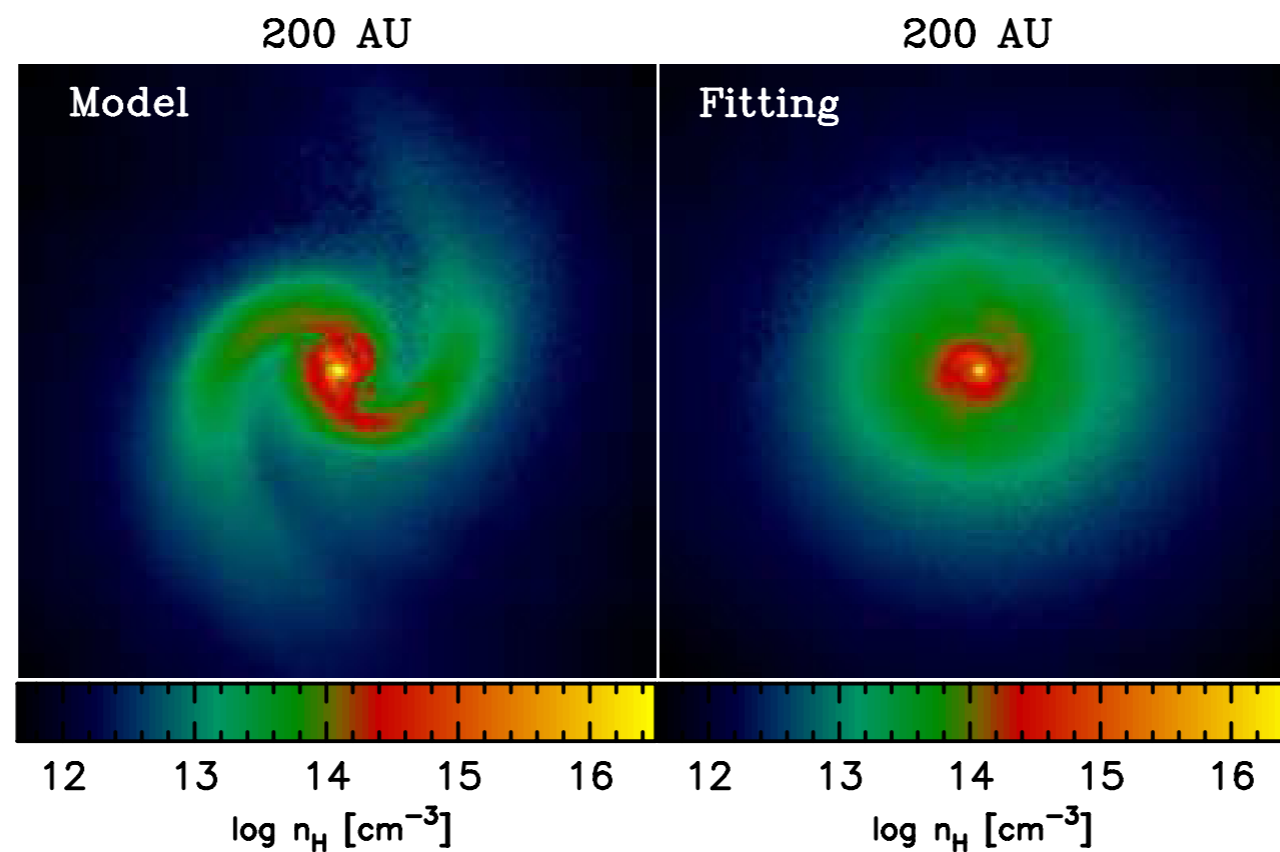
- Working to improve H<sub>2</sub> line cooling using TREECOL.
- Finds that we have been significantly underestimating the amount of cooling that occurs in the 'optically thick' regime.
- TREECOL allows us to better estimate the effective optical depth of the line.





# What about CIE cooling?

- Shingo Hirano has shown the how you treat the CIE cooling can affect the cooling and dynamics.



Hirano et al. (2013)