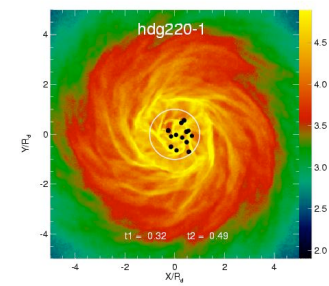
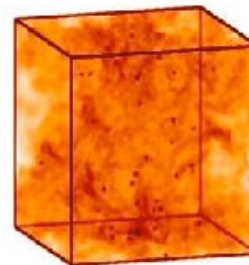
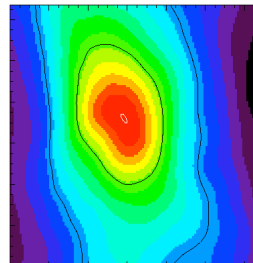
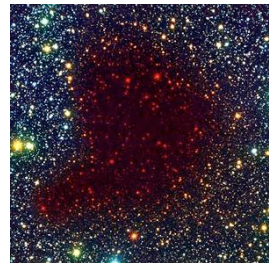


# Molecular Cloud Turbulence and Star Formation



Javier Ballesteros-Paredes<sup>1</sup>, Ralf Klessen<sup>2</sup>, Mordecai-  
Mark Mac Low<sup>3</sup>, Enrique Vazquez-Semadeni<sup>1</sup>

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AIP



DFG

# Overview

- concept of gravoturbulent star formation
- three „steps“ of star formation:
  1. *formation of molecular clouds in the disk of our galaxy*
    - intermezzo:  
*properties of molecular cloud turbulence*
  2. *formation of protostellar cores*
  3. *formation of stars: protostellar collapse and the stellar mass spectrum*
- summary

the idea

# Gravoturbulent star formation

- Idea:

*Star formation is controlled  
by interplay between  
gravity and  
supersonic turbulence!*

- Dual role of turbulence:

- *stability on large scales*
- *initiating collapse on small scales*

(e.g., Larson, 2003, Rep. Prog. Phys, 66, 1651;  
or Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125)

# Gravoturbulent star formation

- Idea:

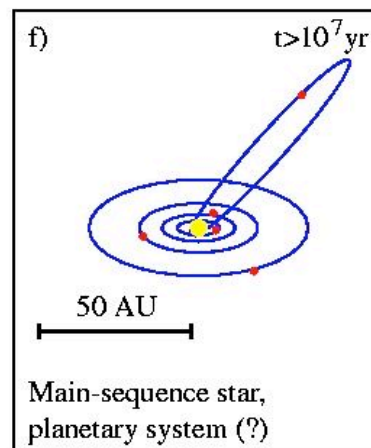
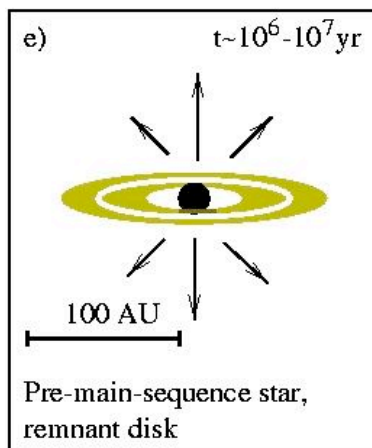
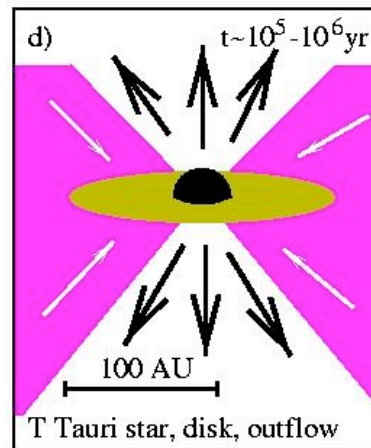
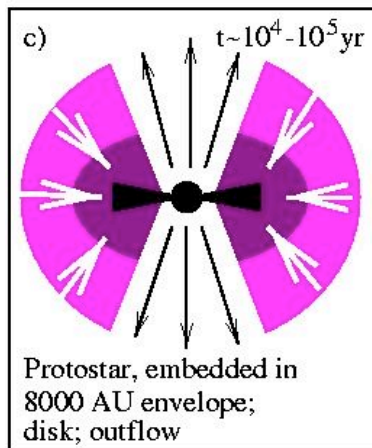
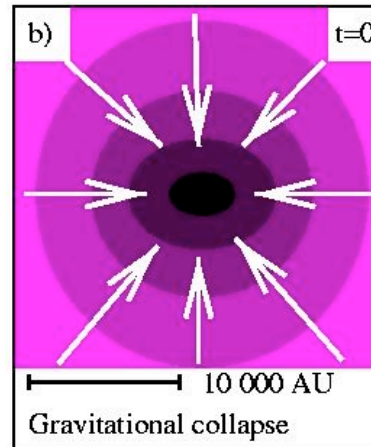
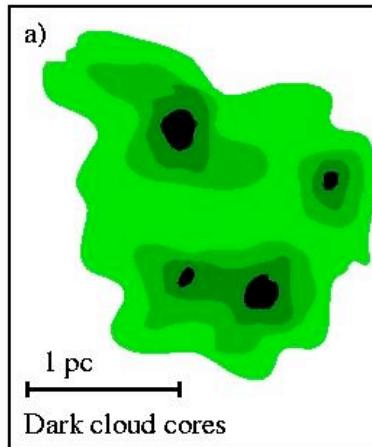
*Star formation is controlled  
by interplay between  
gravity and  
supersonic turbulence!*

- Validity:

This hold on *all* scales and applies to build-up of stars and star clusters within molecular clouds as well as to the formation of molecular clouds in galactic disk.

(e.g., Larson, 2003, Rep. Prog. Phys, 66, 1651;  
or Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125)

# Competing approaches in SF theory:



quasistatic theories:  
magnetically mediated  
star formation

Shu, Adams, & Lizano (1987, ARAA)

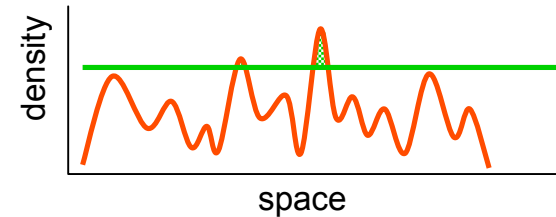


dynamical theories:  
turbulent control of star  
formation

Larson (2003, Prog. Rep. Phys.)  
Mac Low & Klessen (2004, RMP, 76, 125)  
Elmegreen & Scalo (2004, ARAA)  
Scalo & Elmegreen (2004, ARAA)

# Gravoturbulent star formation

- interstellar gas is highly *inhomogeneous*
  - *thermal instability*
  - *gravitational instability*
  - *turbulent compression* (in shocks  $\delta\rho/\rho \propto M^2$ ; in atomic gas:  $M \approx 1\dots3$ )
- cold *molecular clouds* can form rapidly in high-density regions at *stagnation points of convergent large-scale flows*
  - chemical *phase transition*: atomic  $\rightarrow$  molecular
  - process is *modulated* by large-scale *dynamics* in the galaxy
- inside *cold clouds*: turbulence is highly supersonic ( $M \approx 1\dots20$ )  
 $\rightarrow$  *turbulence* creates large density contrast,  
*gravity* selects for collapse  
  
 $\longrightarrow$  **GRAVOTUBULENT FRAGMENTATION**
- *turbulent cascade*: local compression *within* a cloud provokes collapse  
 $\rightarrow$  formation of individual *stars* and *star clusters*



cloud  
formation

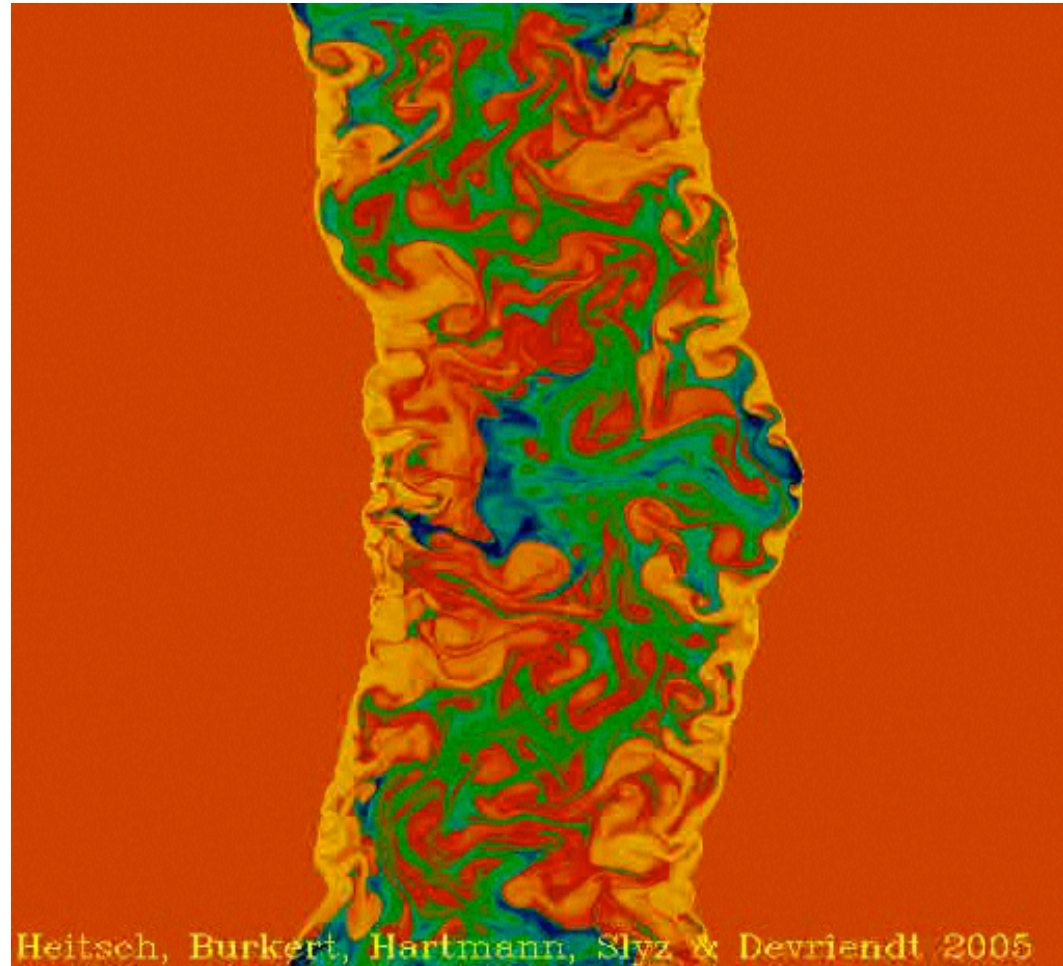


# Molecular cloud formation

... in *convergent large-scale flows*

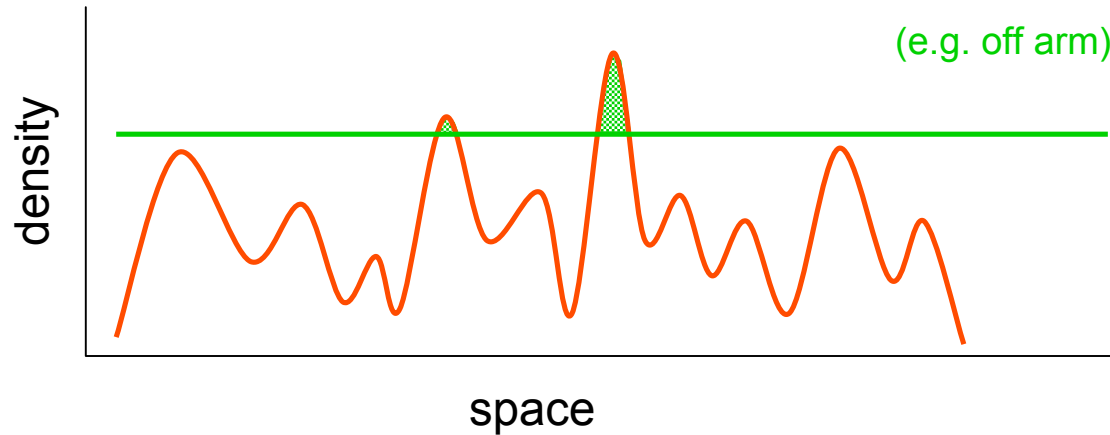
... setting up the *turbulent cascade*

- Mach 3 colliding flow
- Vishniac instability + thermal instability
- compressed sheet *breaks up* and builds up *cold, high-density „blobs“* of gas
- --> molecular cloud formation
- cold cloud motions correspond to supersonic turbulence

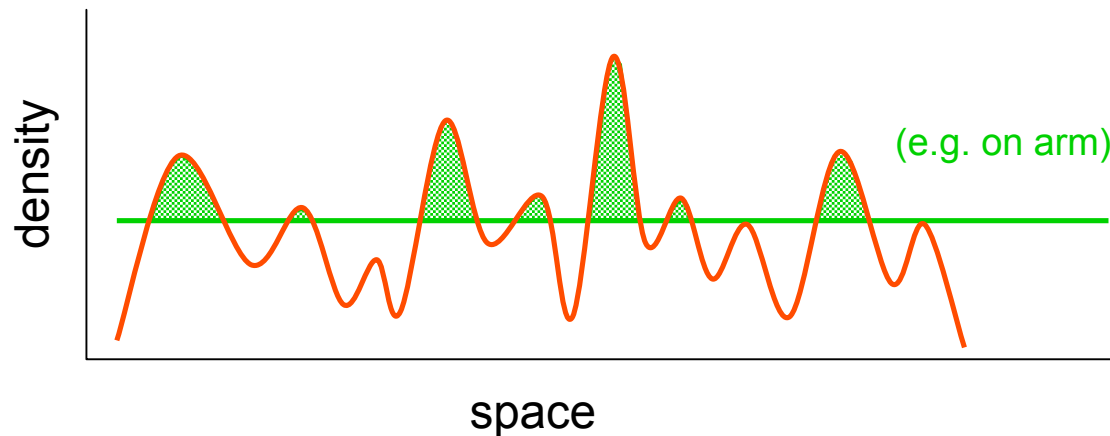


(e.g. Koyama & Inutsuka 2002, Heitsch et al., 2005, Vazquez-Semadeni et al. 2004; also posters 8577, 8302)

# Correlation with large-scale perturbations



*density/temperature fluctuations* in warm atomic ISM are caused by *thermal/gravitational instability* and/or *supersonic turbulence*



some fluctuations are *dense* enough to *form  $H_2$*  within "*reasonable time*"

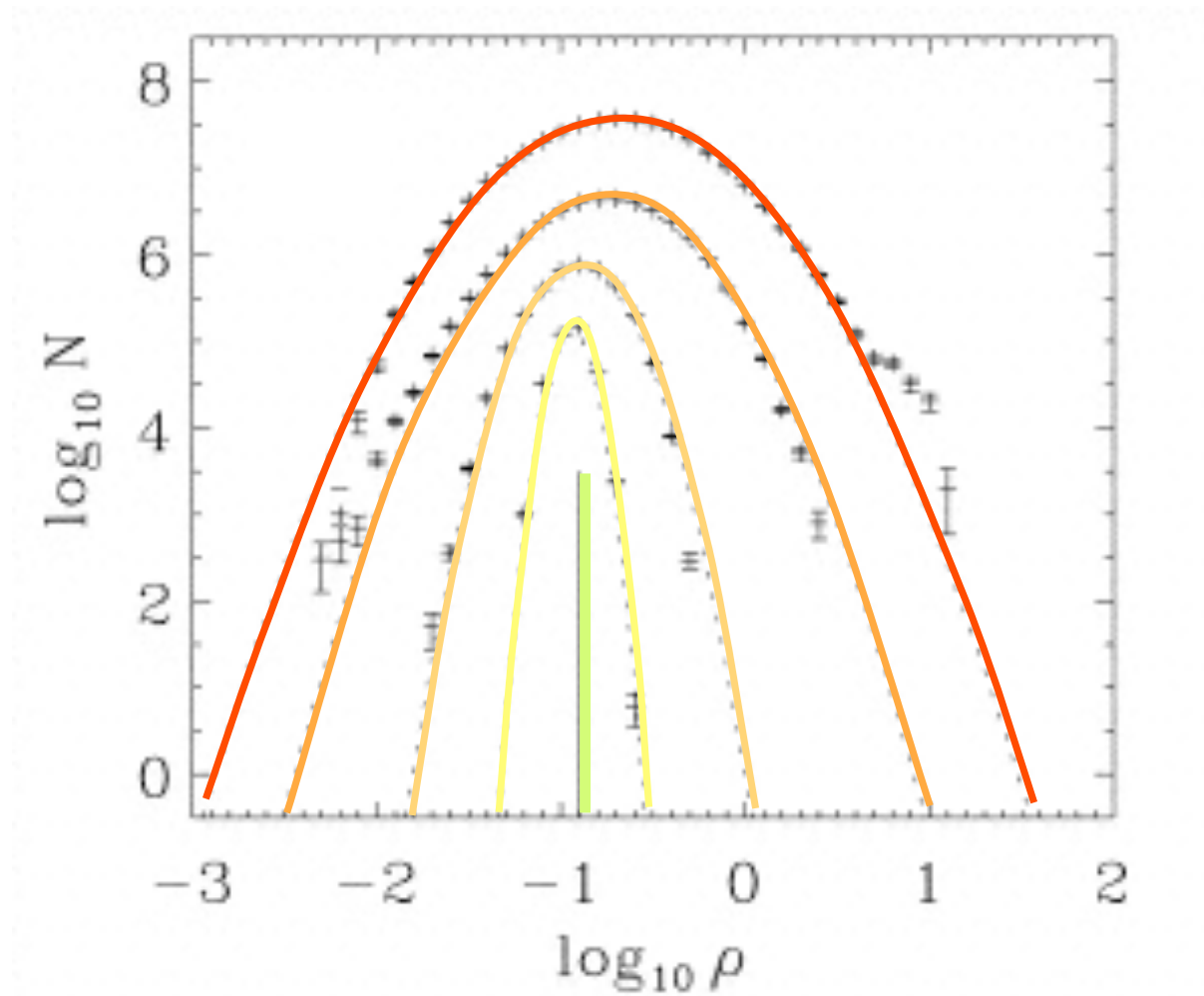
→ *molecular cloud*

(poster 8577: Glover & Mac Low)

*external perturbations* (i.e. potential changes) *increase* likelihood

(poster 8170: Dobbs & Bonnell)

# Star formation on *global scales*



probability distribution  
function of the density  
( $\rho$ -pdf)

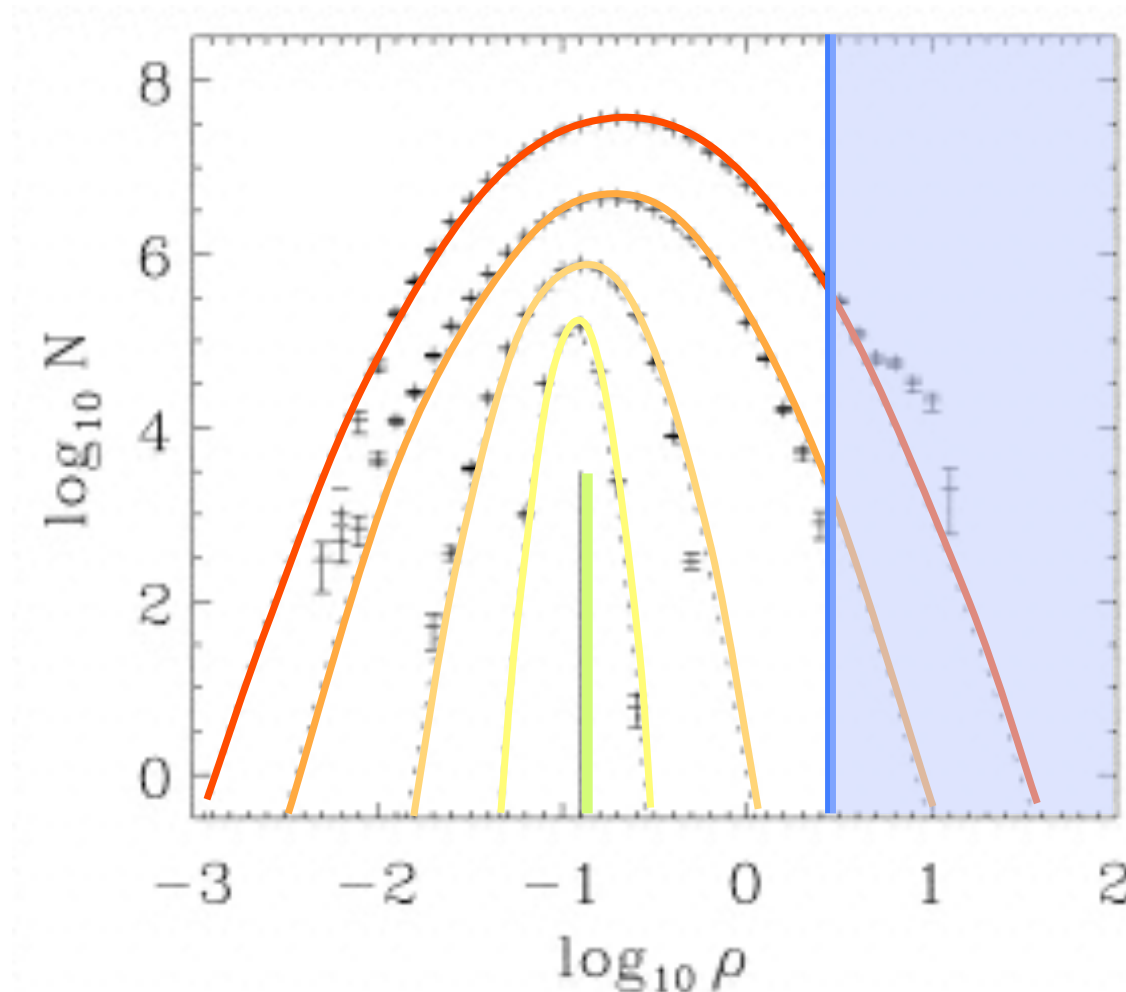
varying rms Mach  
numbers:

**M1** > **M2** >  
**M3** > **M4** > 0

mass weighted  $\rho$ -pdf, each shifted by  $\Delta \log N = 1$

(from Klessen, 2001; also Gazol et al. 2005, Mac Low et al. 2005)

# Star formation on *global scales*



mass weighted  $\rho$ -pdf, each shifted by  $\Delta \log N = 1$

(rate from Hollenback, Werner, & Salpeter 1971, see also poster 8577)

H<sub>2</sub> formation rate:

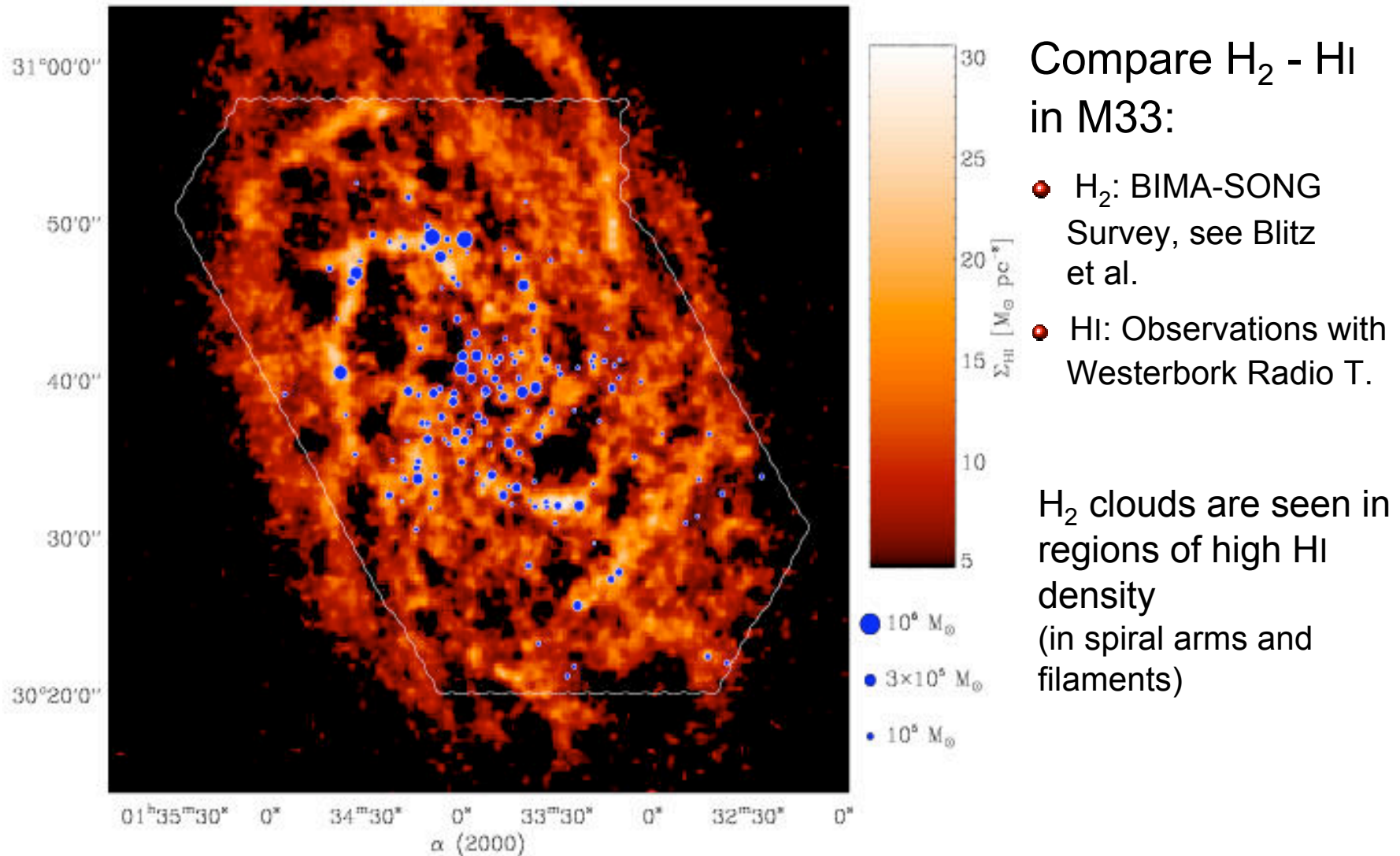
$$\tau_{\text{H}_2} \approx \frac{1.5 \text{ Gyr}}{n_{\text{H}} / 1 \text{ cm}^{-3}}$$

for  $n_{\text{H}} \geq 100 \text{ cm}^{-3}$ , H<sub>2</sub> forms within 10 Myr, this is about the lifetime of typical MC's.

in turbulent gas, the H<sub>2</sub> fraction can become very high on short timescale

(for models with coupling between cloud dynamics and time-dependent chemistry, see Glover & Mac Low 2005)

# Correlation between H<sub>2</sub> and HI



(Deul & van der Hulst 1987, Blitz et al. 2004)

turbulence

# Properties of turbulence

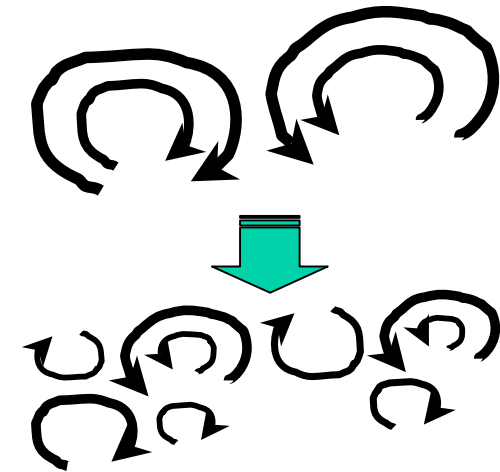
- laminar flows turn *turbulent* at *high Reynolds* numbers

$$\text{Re} = \frac{\text{advection}}{\text{dissipation}} = \frac{VL}{\nu}$$

$V$  = typical velocity on scale  $L$ ,  $\nu$  = viscosity,  $\text{Re} > 1000$

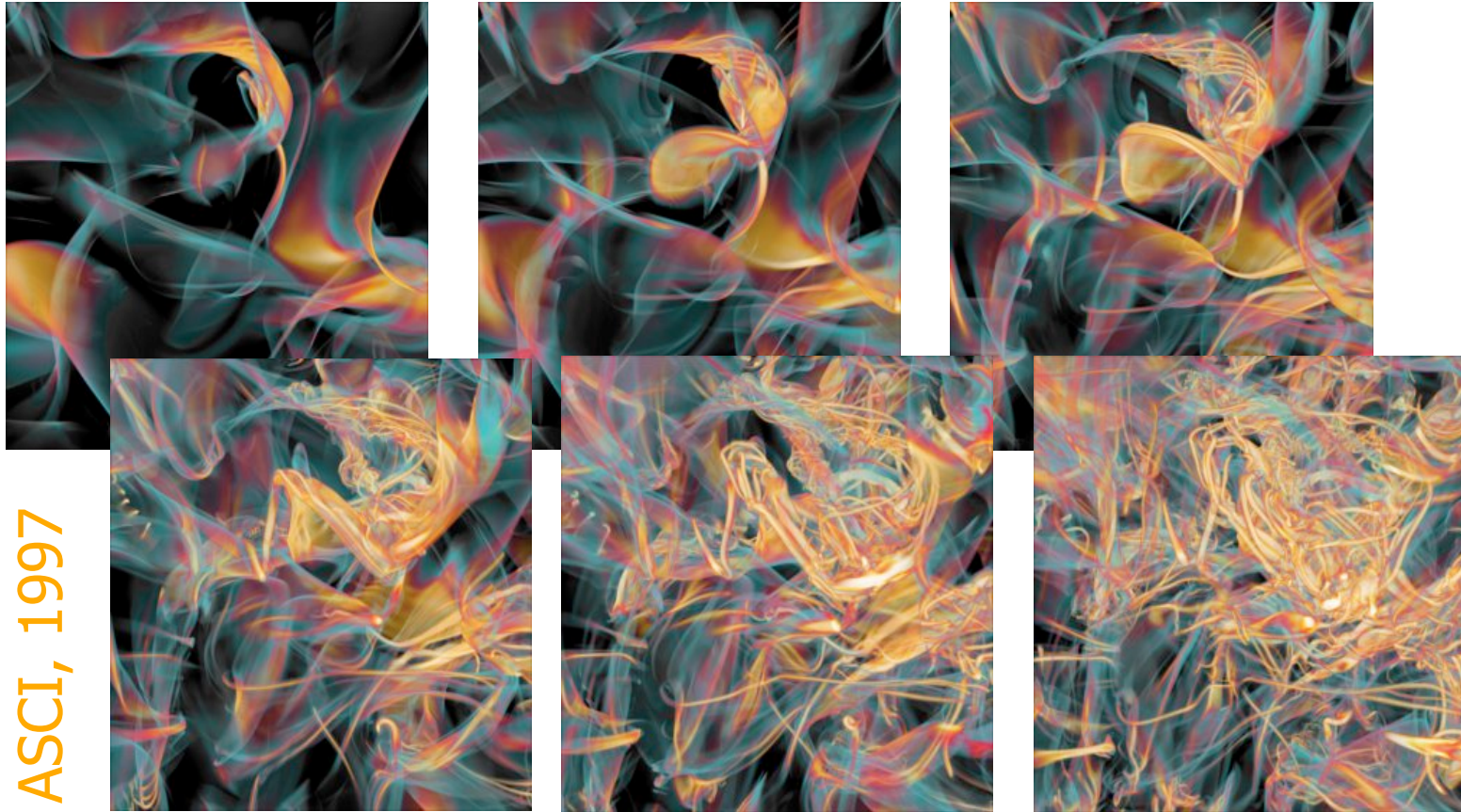
- *vortex stretching* --> turbulence is *intrinsically anisotropic* (only on large scales you *may* get homogeneity & isotropy in a statistical sense; see Landau & Lifschitz, Chandrasekhar, Taylor, etc.)

(ISM turbulence: shocks & B-field cause additional inhomogeneity)



# Vortex Formation

Porter et al.  
ASCI, 1997

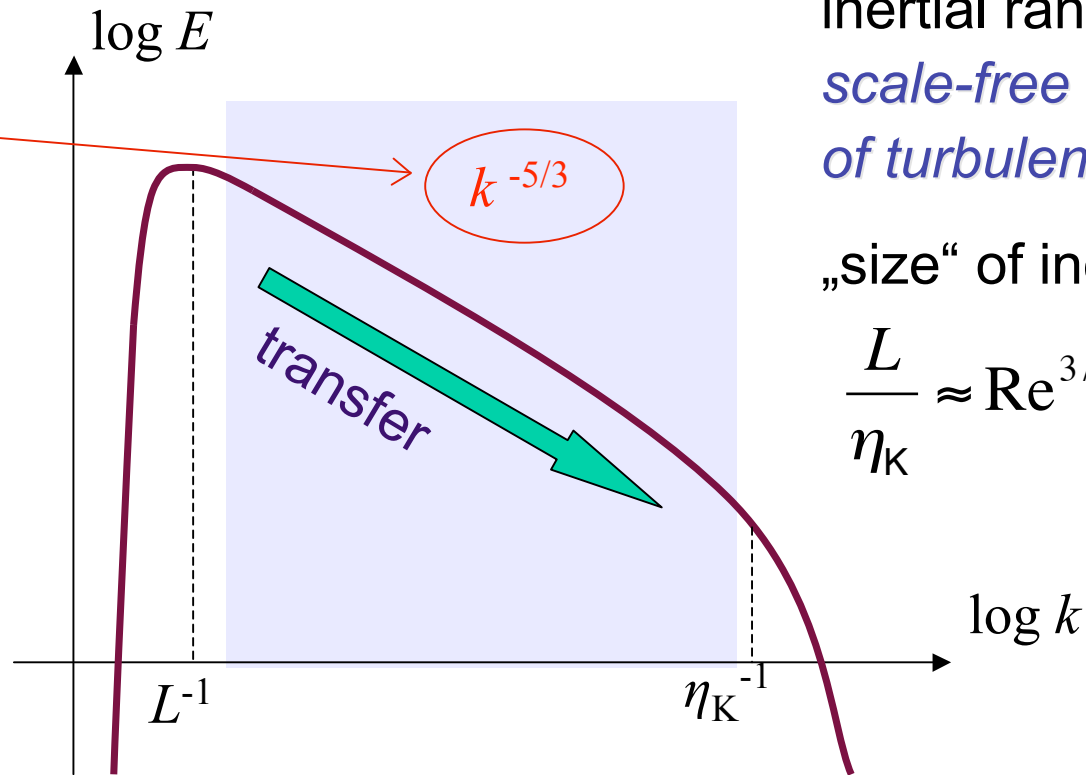


Vortices are stretched and folded in **three dimensions**



# Turbulent cascade

Kolmogorov (1941) theory  
incompressible turbulence



inertial range:  
*scale-free behavior  
of turbulence*

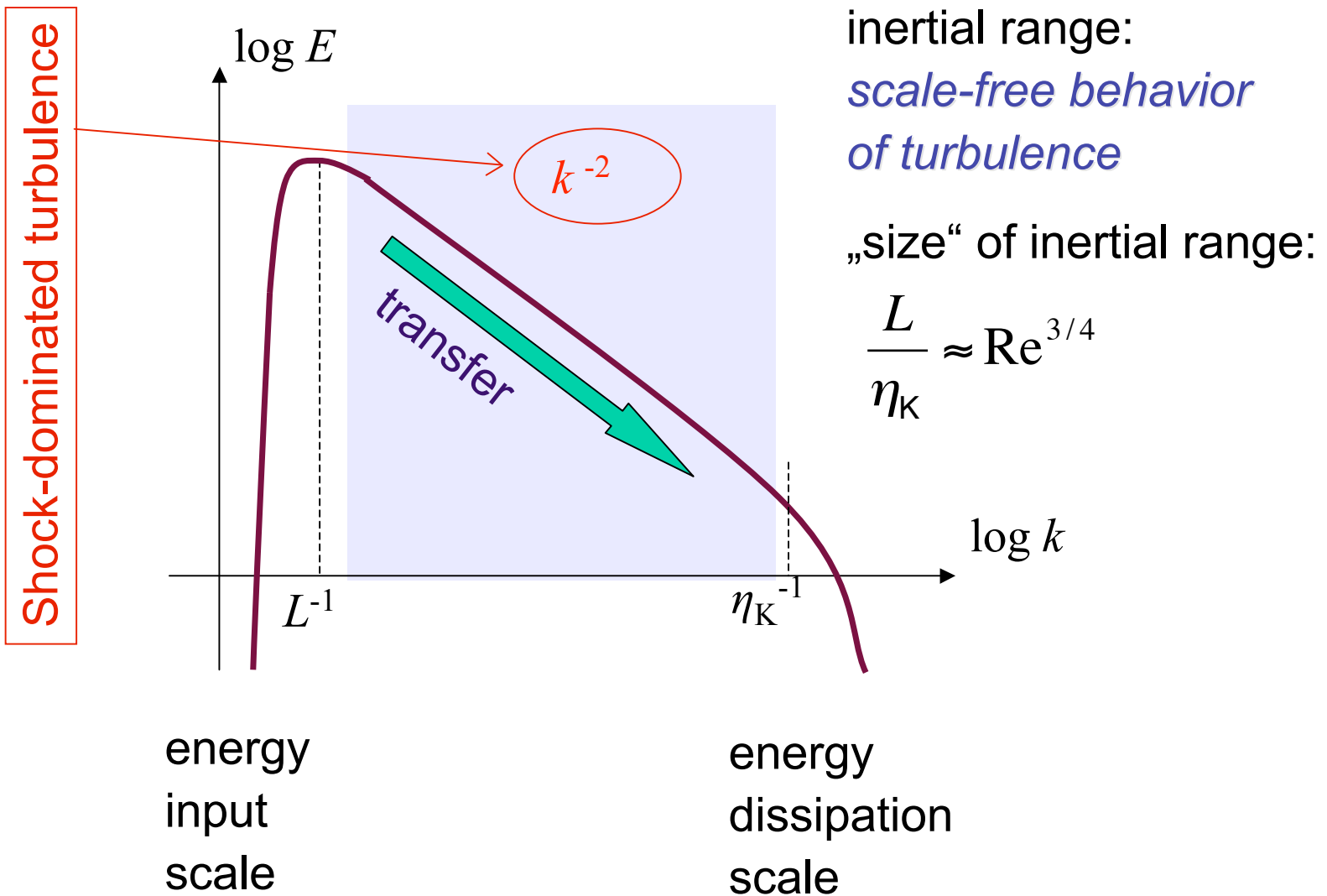
„size“ of inertial range:

$$\frac{L}{\eta_K} \approx \text{Re}^{3/4}$$

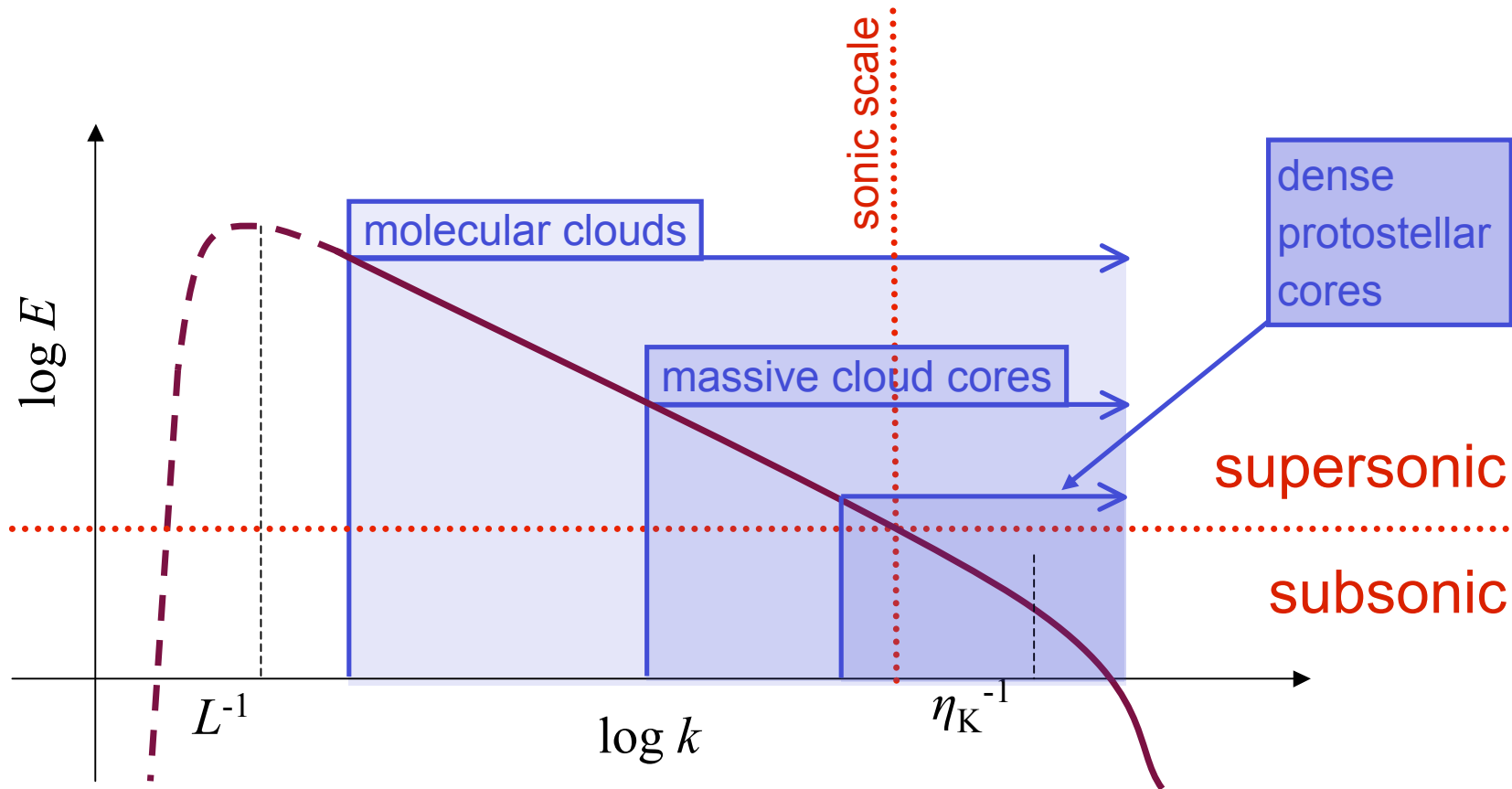
energy  
input  
scale

energy  
dissipation  
scale

# Turbulent cascade



# Turbulent cascade in ISM



energy source & scale  
*NOT known*  
 (supernovae, winds,  
 spiral density waves?)

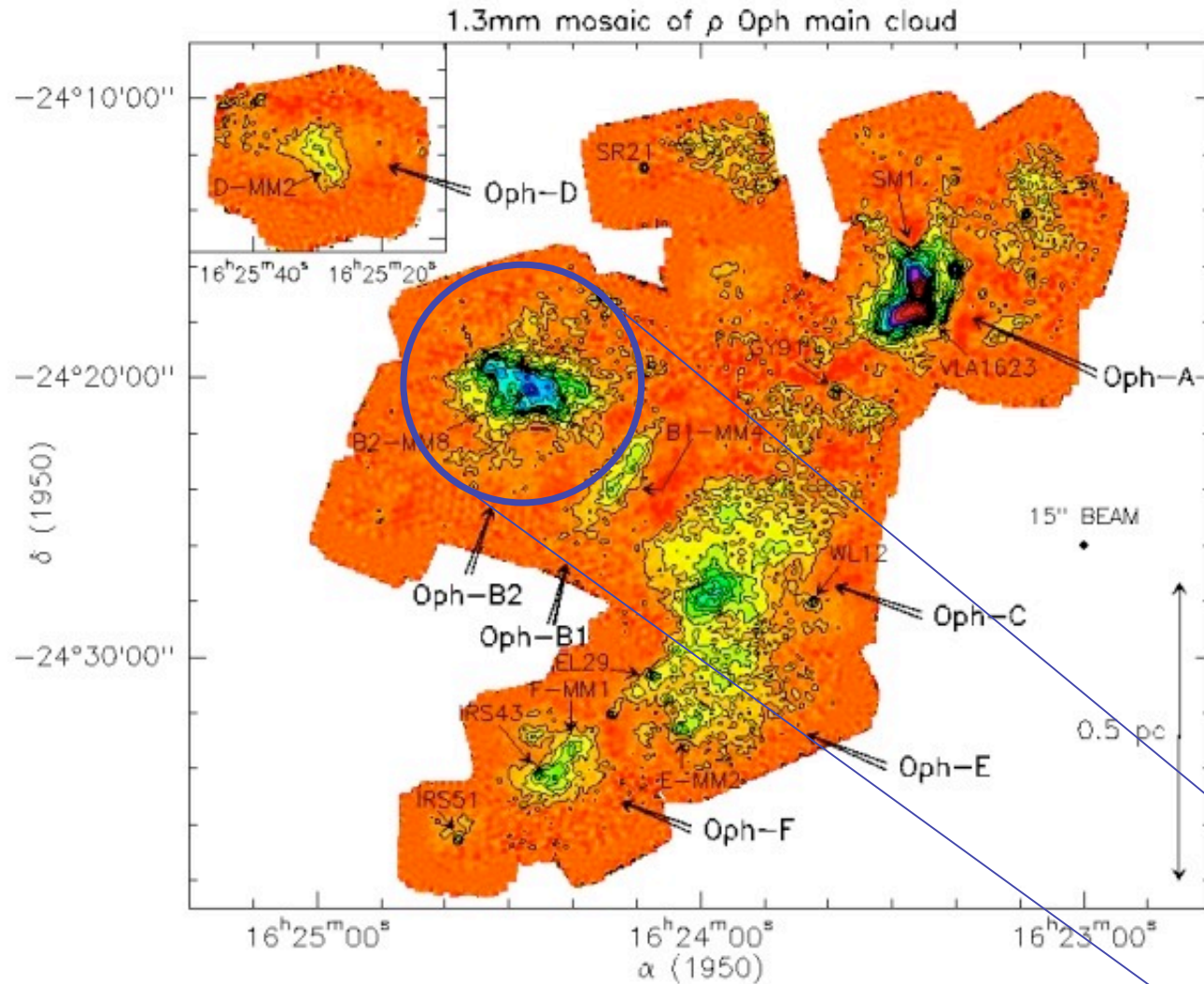
$$\sigma_{\text{rms}} \ll 1 \text{ km/s}$$

$$M_{\text{rms}} \leq 1$$

$$L \approx 0.1 \text{ pc}$$

dissipation scale not known  
 (ambipolar diffusion,  
 molecular diffusion?)

# Density structure of MC's



molecular clouds  
are highly  
inhomogeneous

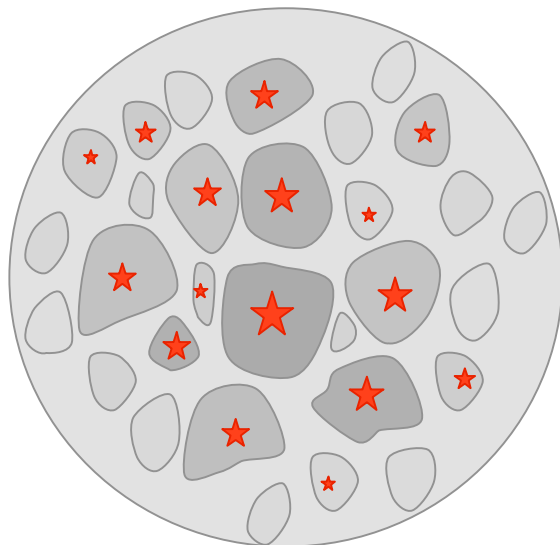
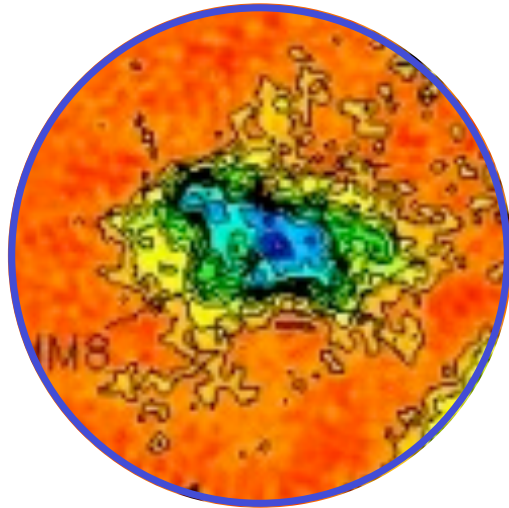
stars form in the  
densest and  
coldest parts of  
the cloud

$\rho$ -Ophiuchus  
cloud seen in dust  
emission

let's focus on  
a cloud core  
like this one

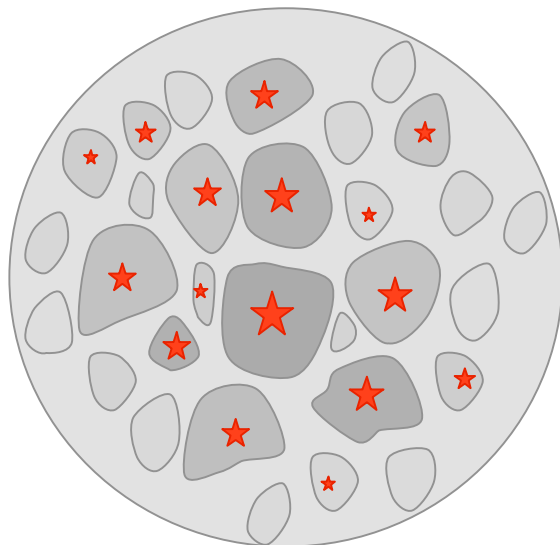
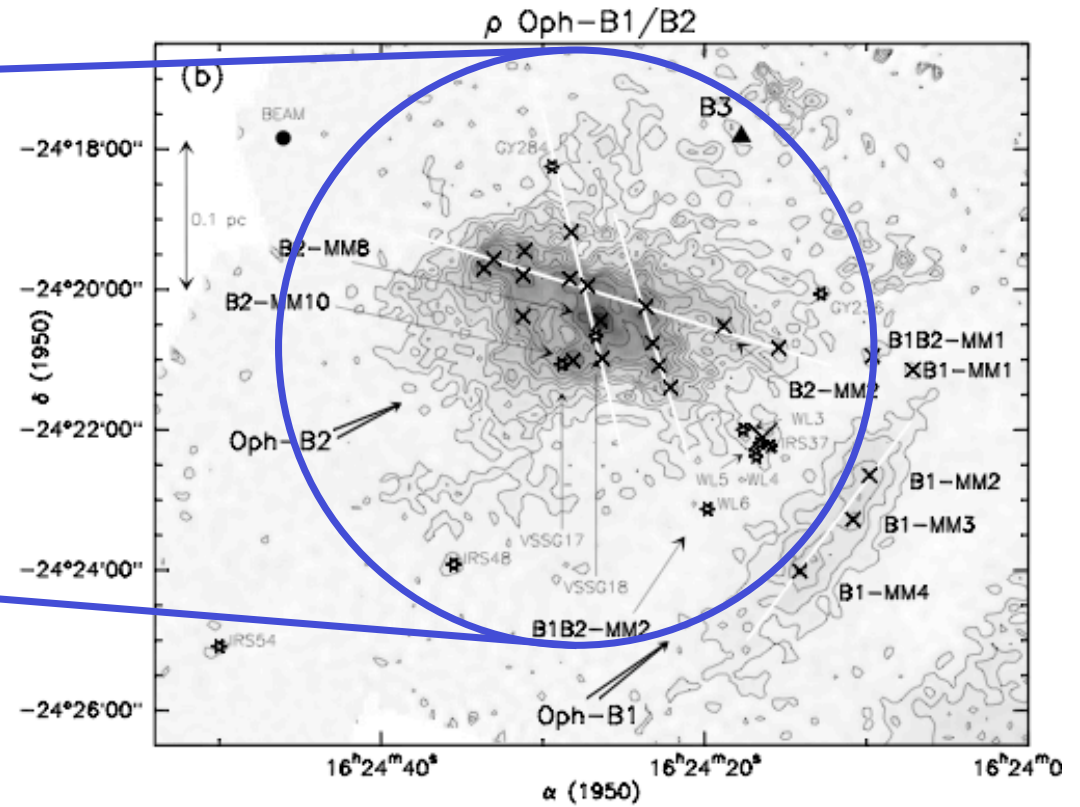
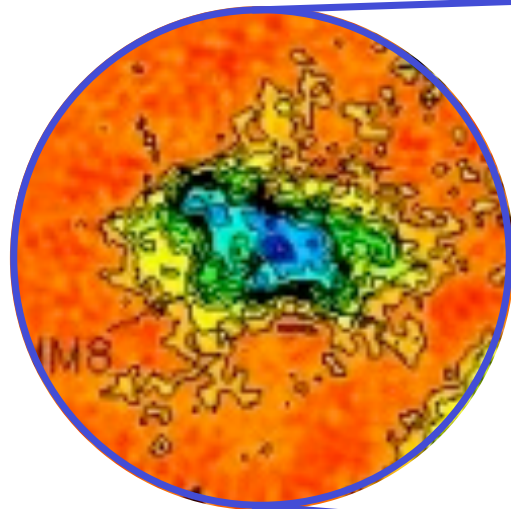
(Motte, André, & Neri 1998)

# Evolution of cloud cores



- Does core form single massive star or cluster with mass distribution?
- Turbulent cascade „goes through“ cloud core
  - > NO *scale separation* possible
  - > NO *effective sound speed*
- Turbulence is supersonic!
  - > produces strong density contrasts:  
 $\delta\rho/\rho \approx M^2$
  - > with typical  $M \approx 10$  -->  $\delta\rho/\rho \approx 100!$
- many of the shock-generated fluctuations are Jeans unstable and go into collapse
- --> *core breaks up and forms a cluster of stars*

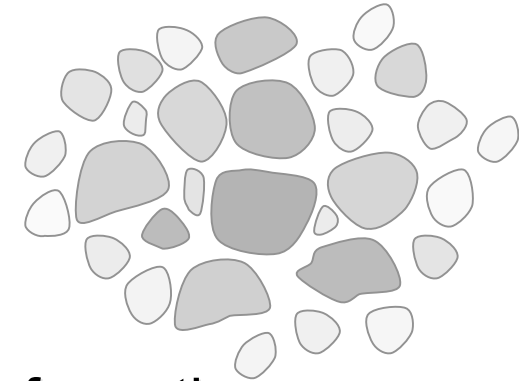
# Evolution of cloud cores



indeed  $\rho$ -Oph B1/2 contains several cores (“starless” cores are denoted by  $\times$ , cores with embedded protostars by  $\star$ )

(Motte, André, & Neri 1998)

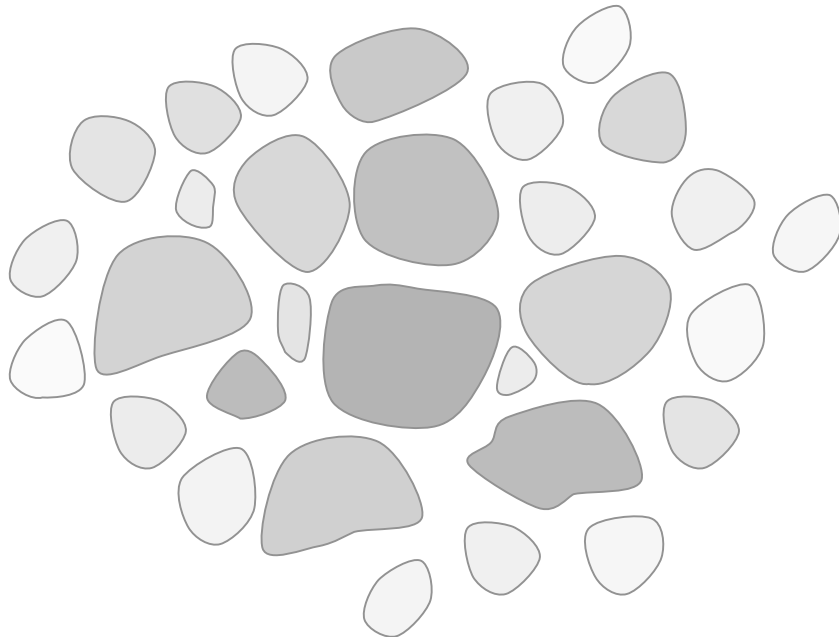
# Formation and evolution of cores



- protostellar cloud cores form at the *stagnation points* of *convergent turbulent flows*
- if  $M > M_{\text{Jeans}} \propto \rho^{-1/2} T^{3/2}$ : collapse and star formation
- if  $M < M_{\text{Jeans}} \propto \rho^{-1/2} T^{3/2}$ : reexpansion after external compression fades away  
(e.g. Vazquez-Semadeni et al 2005)
- typical timescales:  $t \approx 10^4 \dots 10^5$  yr
- because *turbulent* ambipolar diffusion time is *short*, this time estimate still holds for the presence of magnetic fields, in *magnetically critical cores*  
(e.g. Fatuzzo & Adams 2002, Heitsch et al. 2004)

# Formation and evolution of cores

What happens to distribution of cloud cores?



Two extreme cases:

(1) turbulence dominates energy budget:

$$\alpha = E_{\text{kin}} / |E_{\text{pot}}| > 1$$

--> individual cores do *not* interact

--> *collapse of individual cores dominates stellar mass growth*

--> *loose cluster of low-mass stars*

(2) turbulence decays, i.e. gravity dominates:  $\alpha = E_{\text{kin}} / |E_{\text{pot}}| < 1$

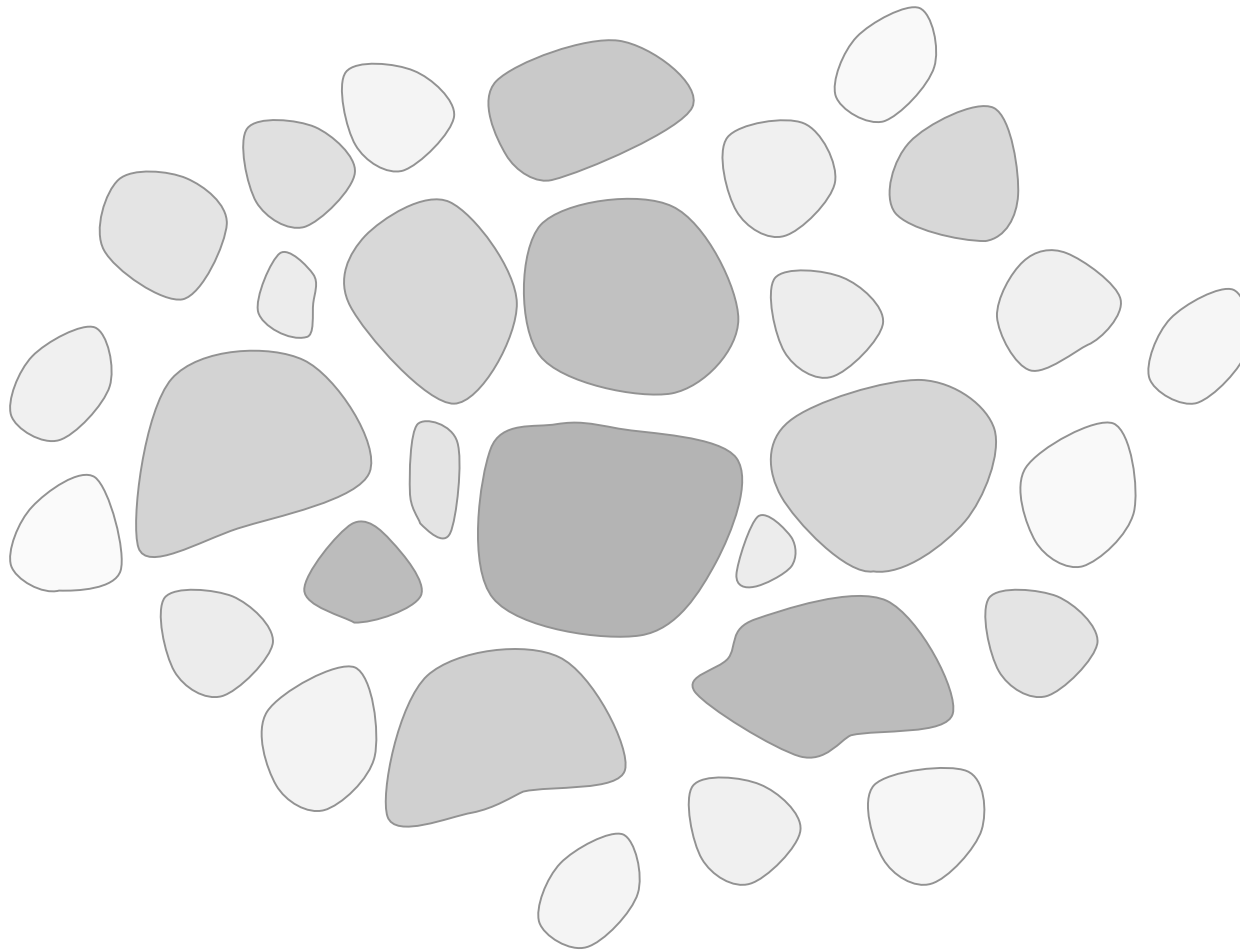
--> *global contraction*

--> *core do interact while collapsing*

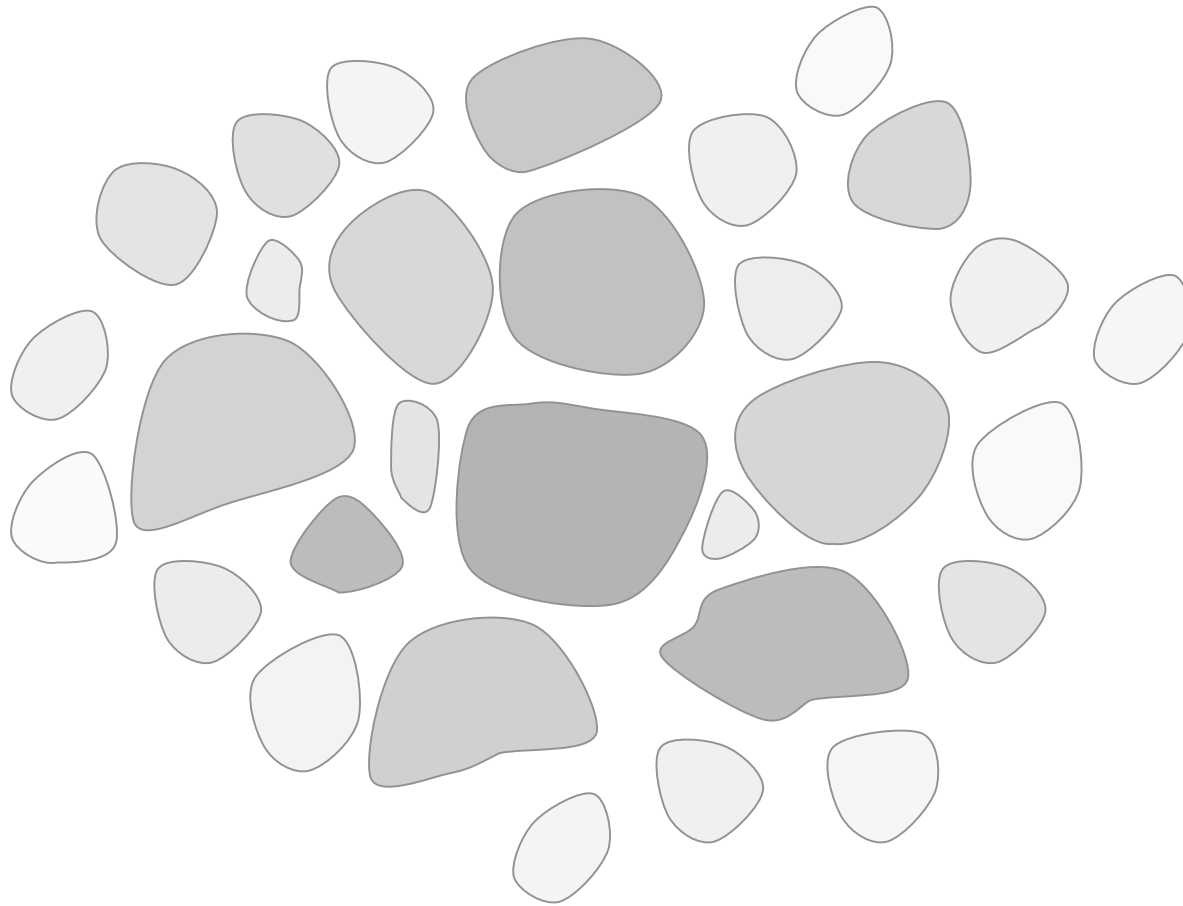
--> *competition influences mass growth*

--> *dense cluster with high-mass stars*

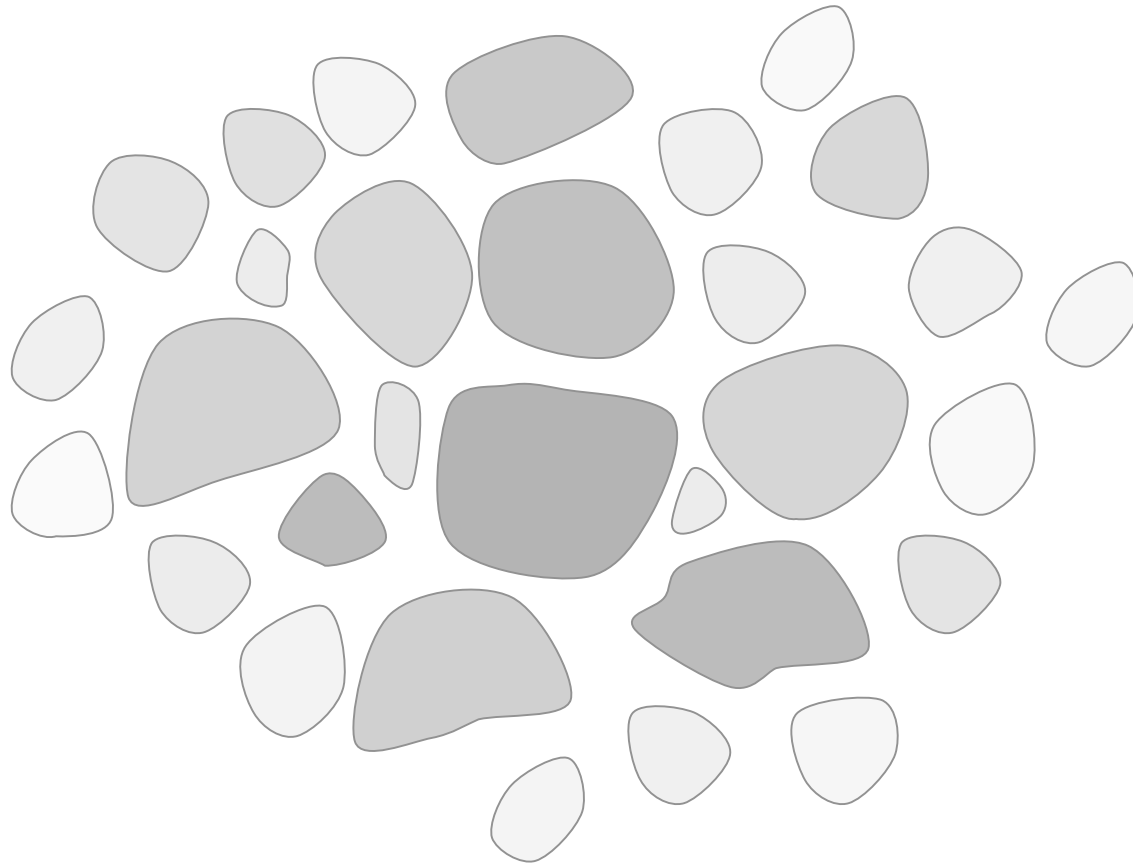




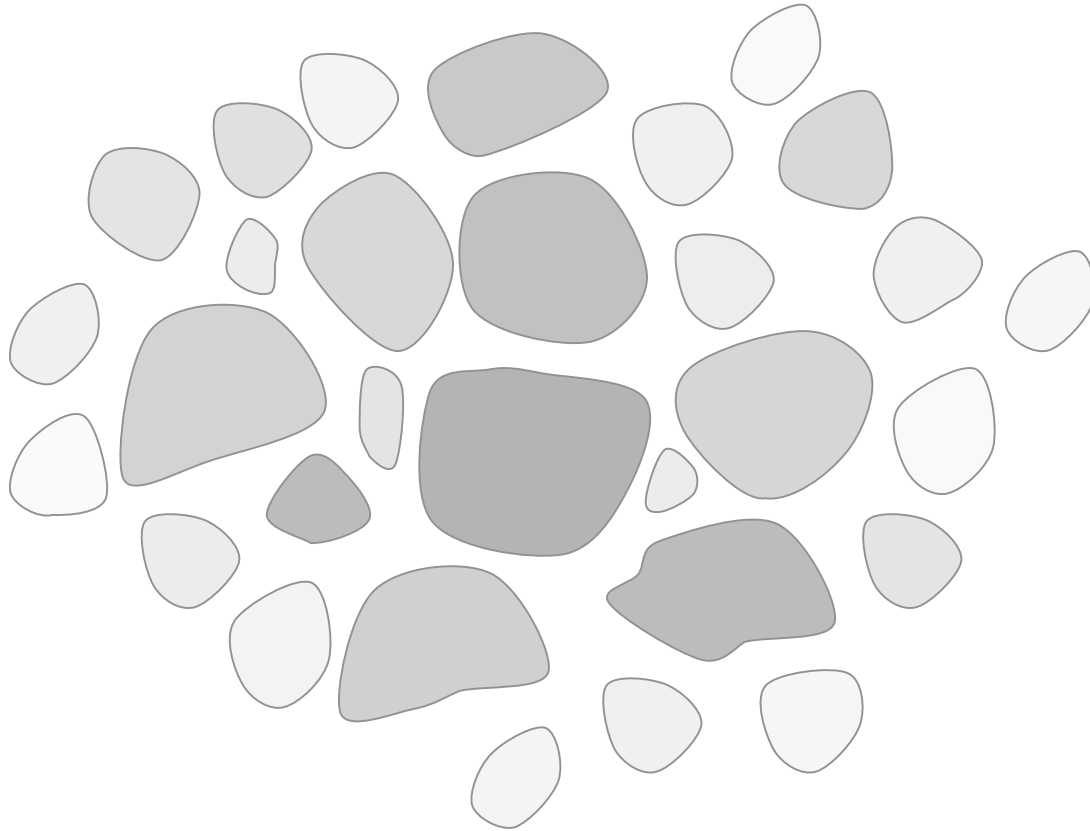
turbulence creates a hierarchy of clumps



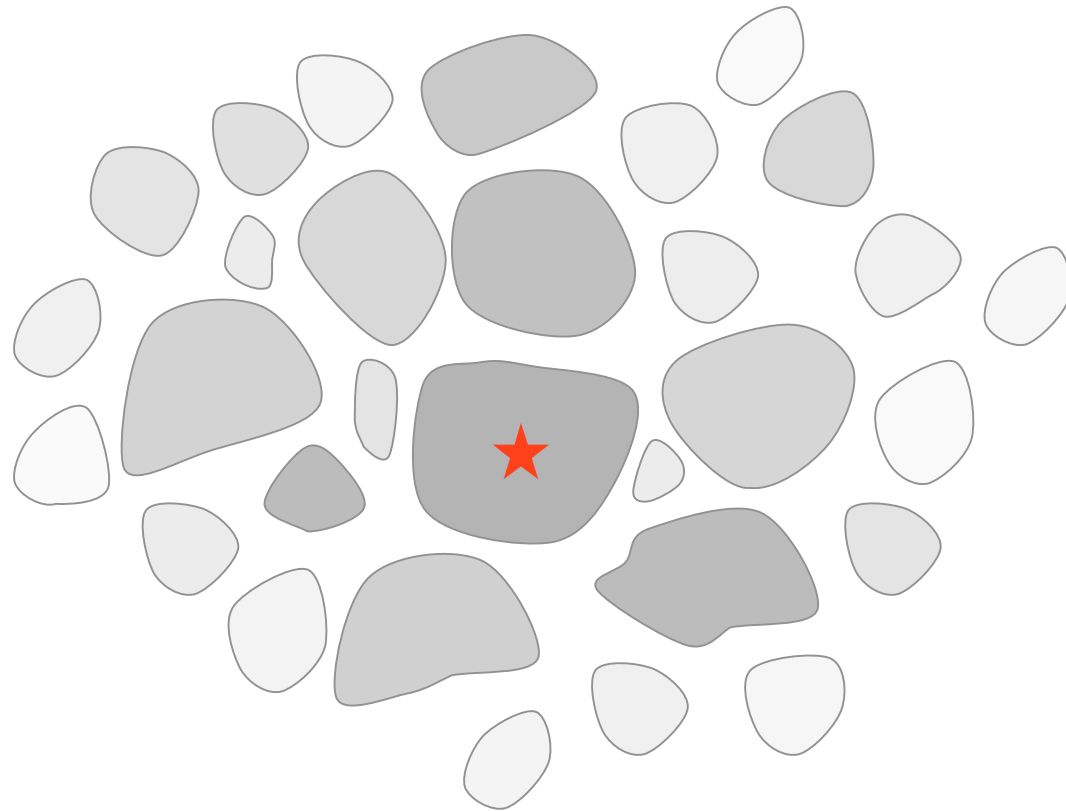
as turbulence decays locally, contraction sets in



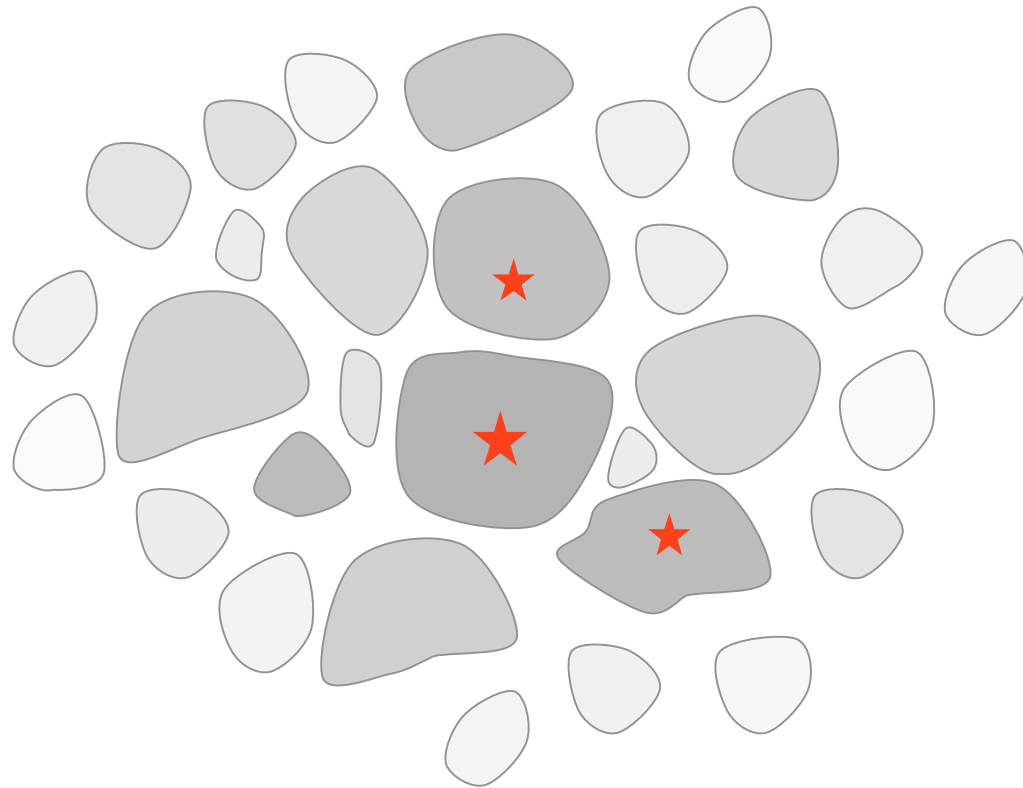
as turbulence decays locally, contraction sets in



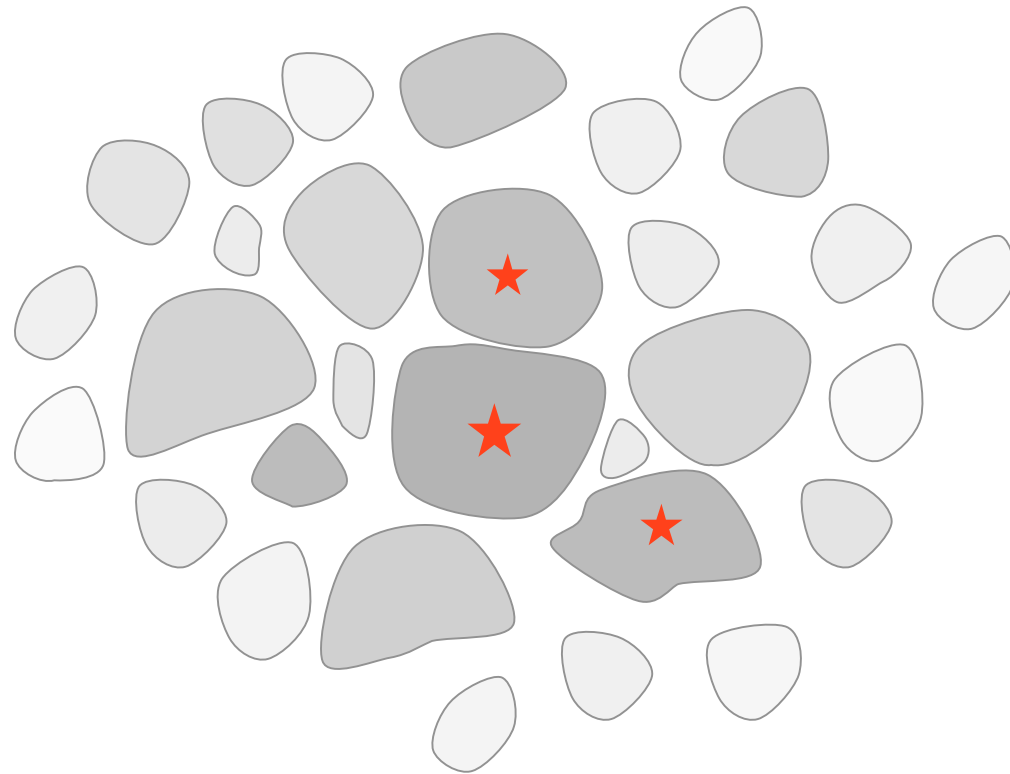
while region contracts, individual clumps collapse to form stars



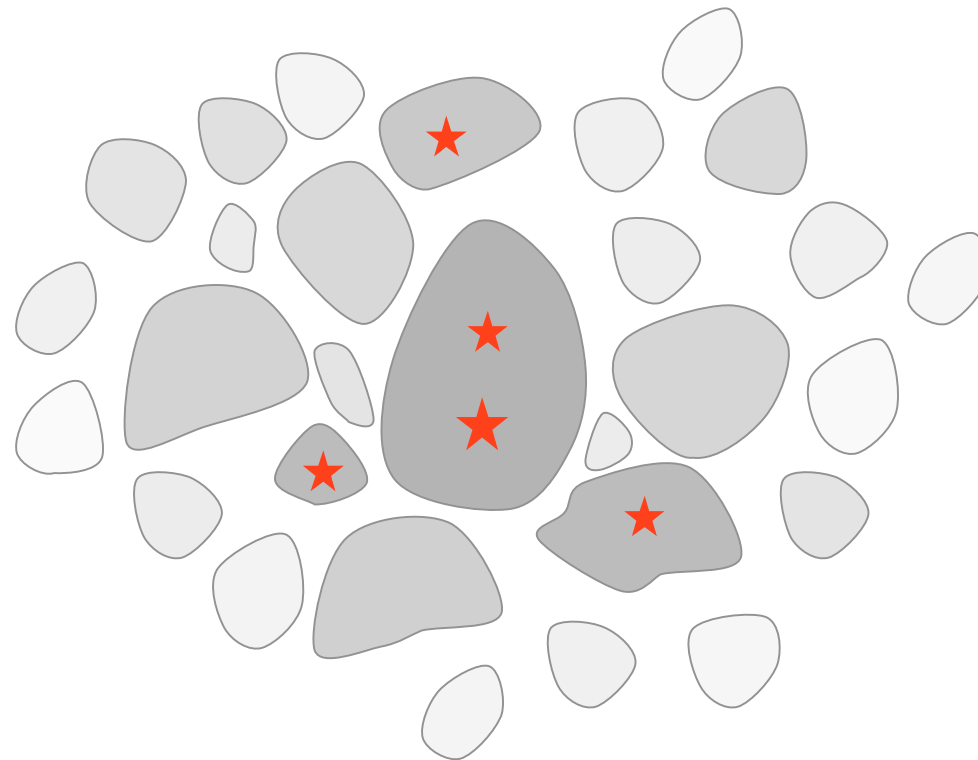
while region contracts, individual clumps collapse to form stars



individual clumps collapse to form stars

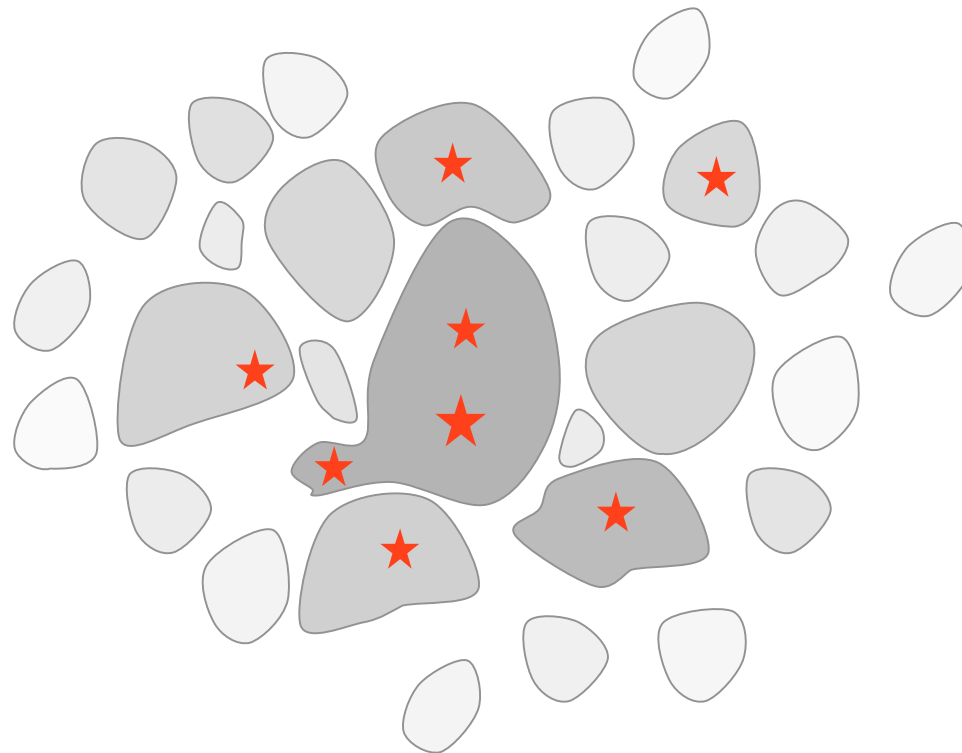


individual clumps collapse to form stars

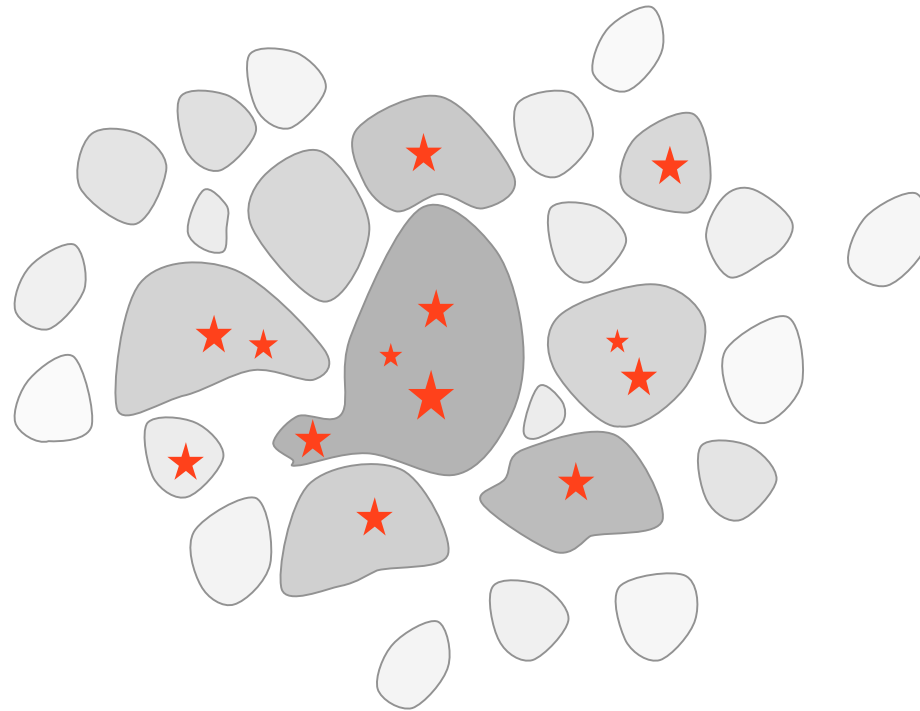


in *dense clusters*, clumps may merge while collapsing  
--> then contain multiple protostars

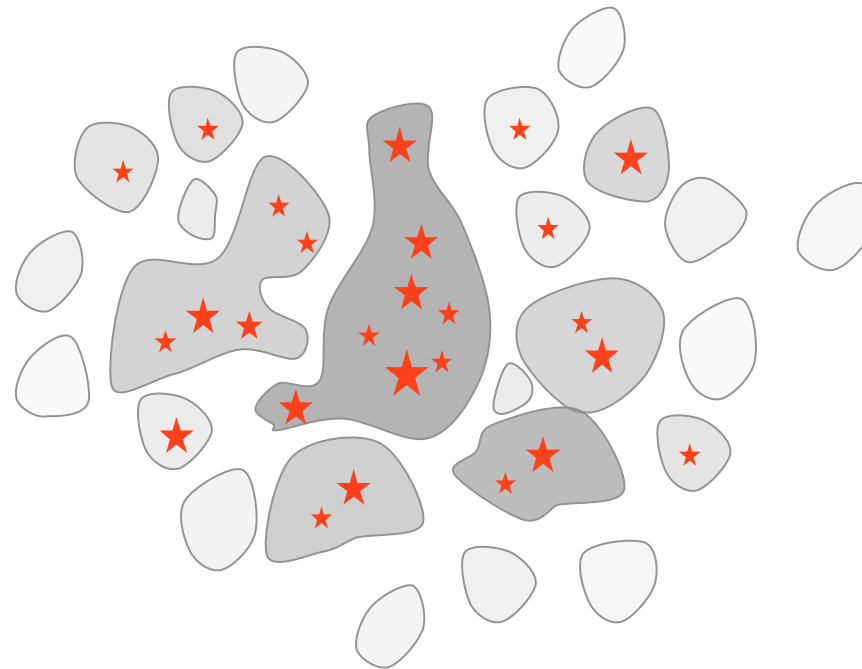




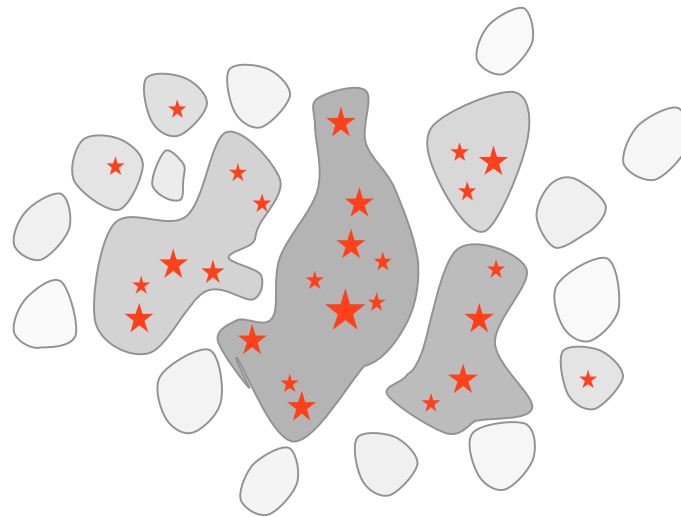
in *dense clusters*, clumps may merge while collapsing  
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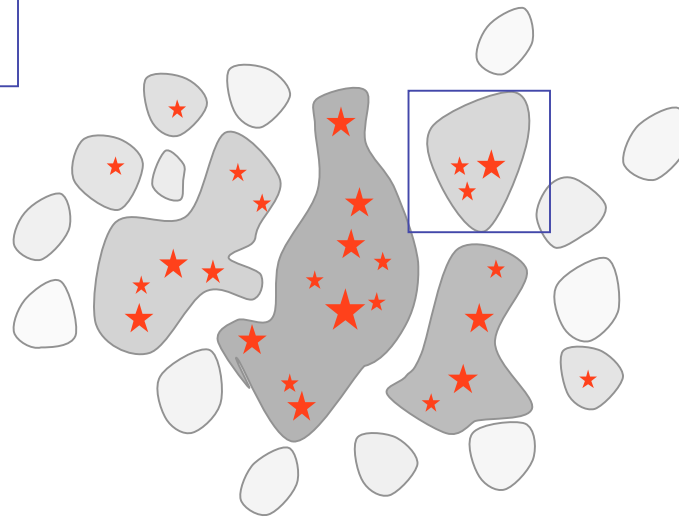
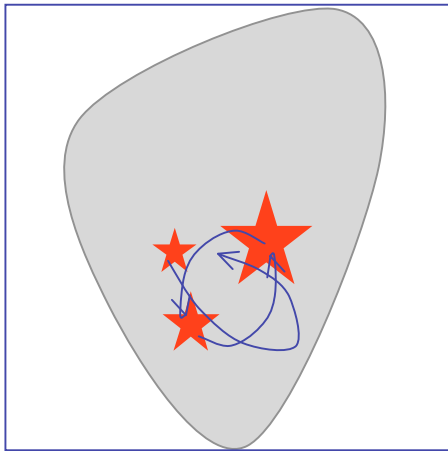
in *dense clusters*, clumps may merge while collapsing  
--> then contain multiple protostars



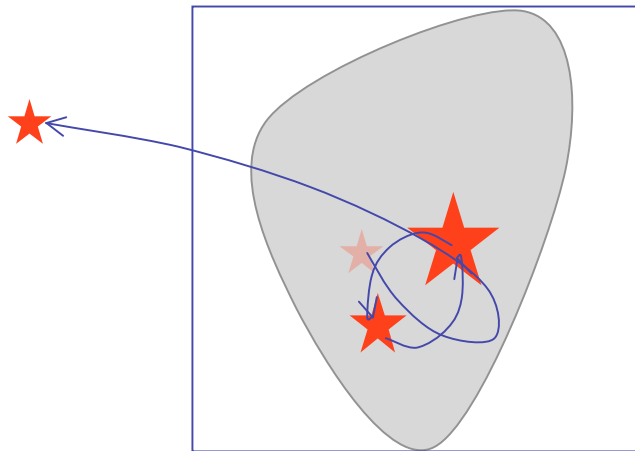
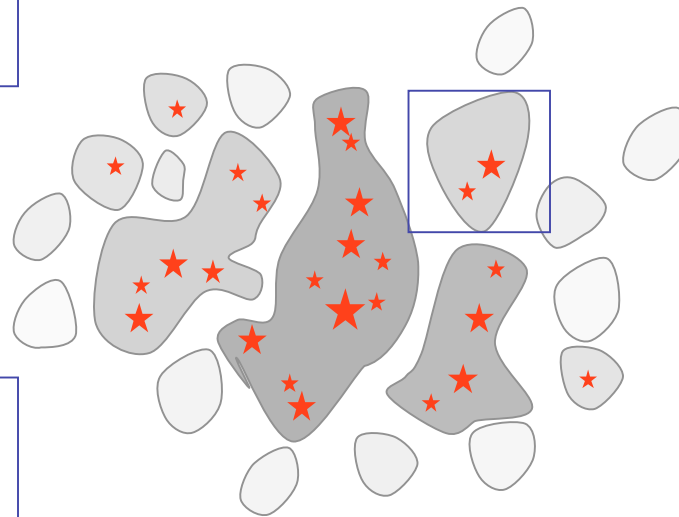
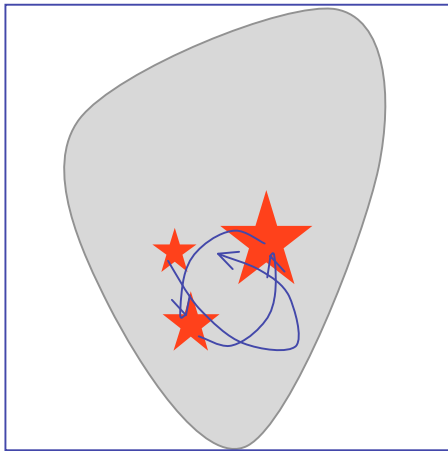
in *dense clusters*, competitive mass growth becomes important



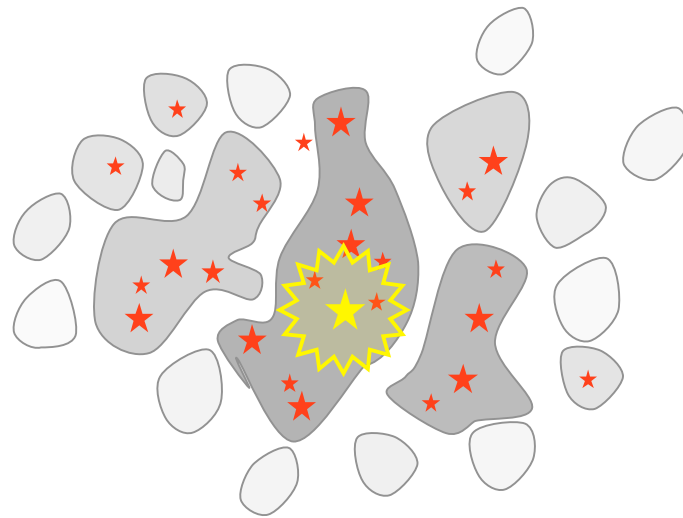
in *dense clusters*, competitive mass growth becomes important



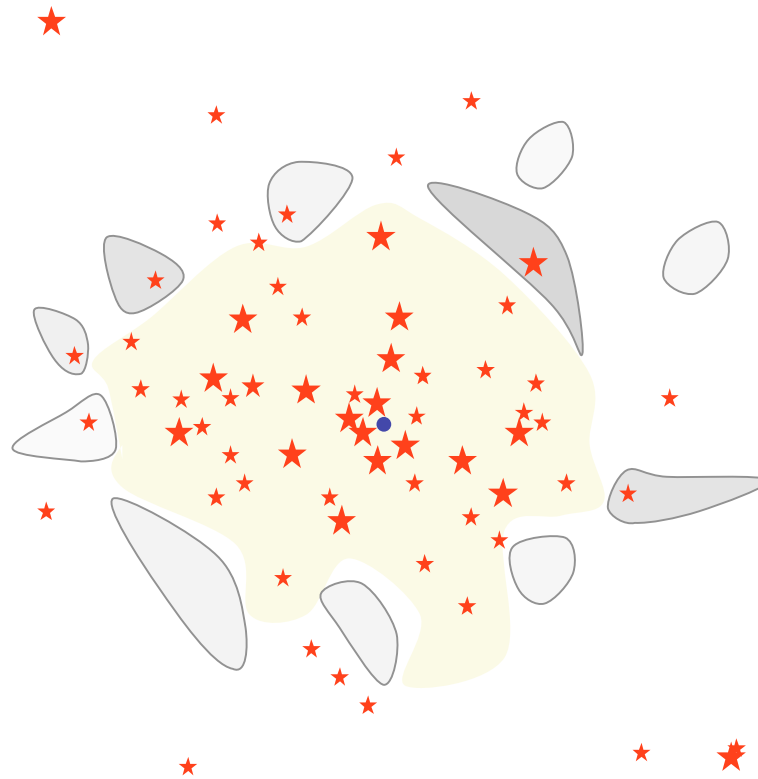
in *dense clusters*,  $N$ -body effects influence mass growth



low-mass objects may  
become ejected --> accretion stops



feedback terminates star formation



result: *star cluster*, possibly with H II region



predictions

# Predictions

- *global properties* (statistical properties)
  - SF efficiency and timescale
  - stellar mass function -- IMF
  - dynamics of young star clusters
  - description of self-gravitating turbulent systems (pdf's,  $\Delta$ -var.)
  - chemical mixing properties
- *local properties* (properties of individual objects)
  - properties of individual clumps (e.g. shape, radial profile, lifetimes)
  - accretion history of individual protostars ( $dM/dt$  vs.  $t$ ,  $j$  vs.  $t$ )
  - binary (proto)stars (eccentricity, mass ratio, etc.)
  - SED's of individual protostars
  - dynamic PMS tracks:  $T_{\text{bol}}-L_{\text{bol}}$  evolution

# Examples and predictions

*example 1:* transient structure of turbulent clouds

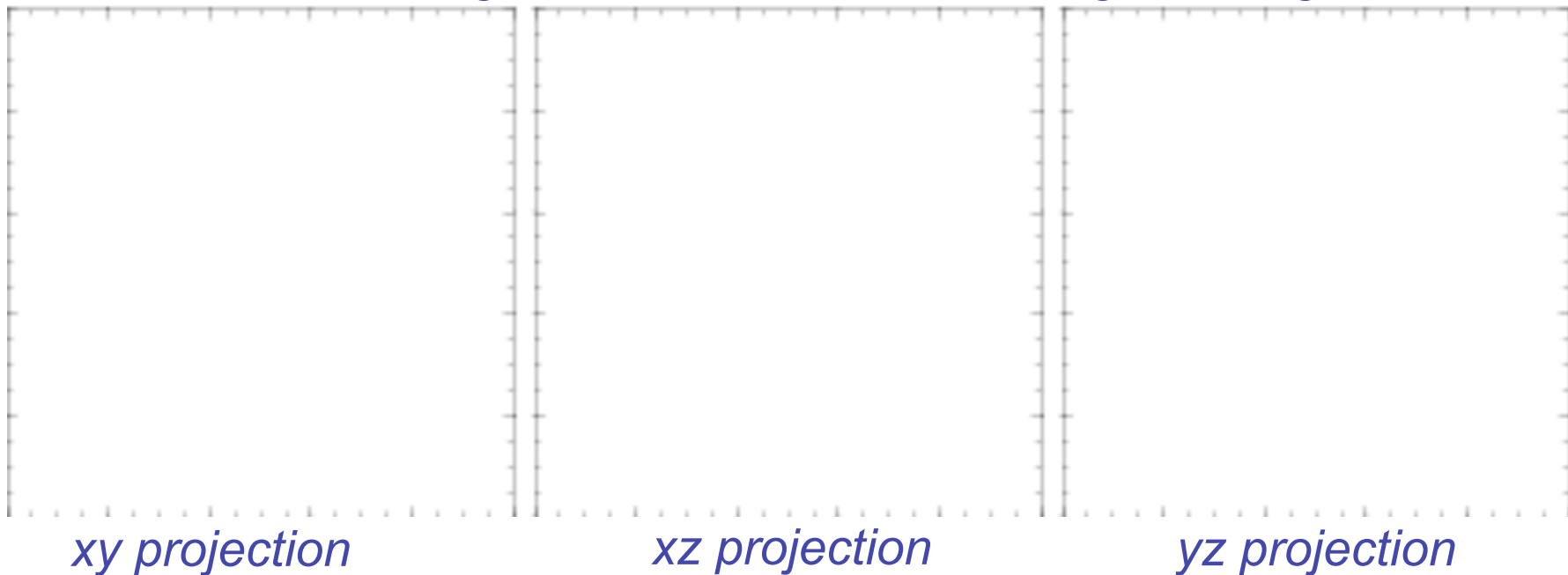
*example 2:* quiescent and coherent appearance of molecular cloud cores

*example 3:* speculations on the origin of the stellar mass spectrum (IMF)

example 1

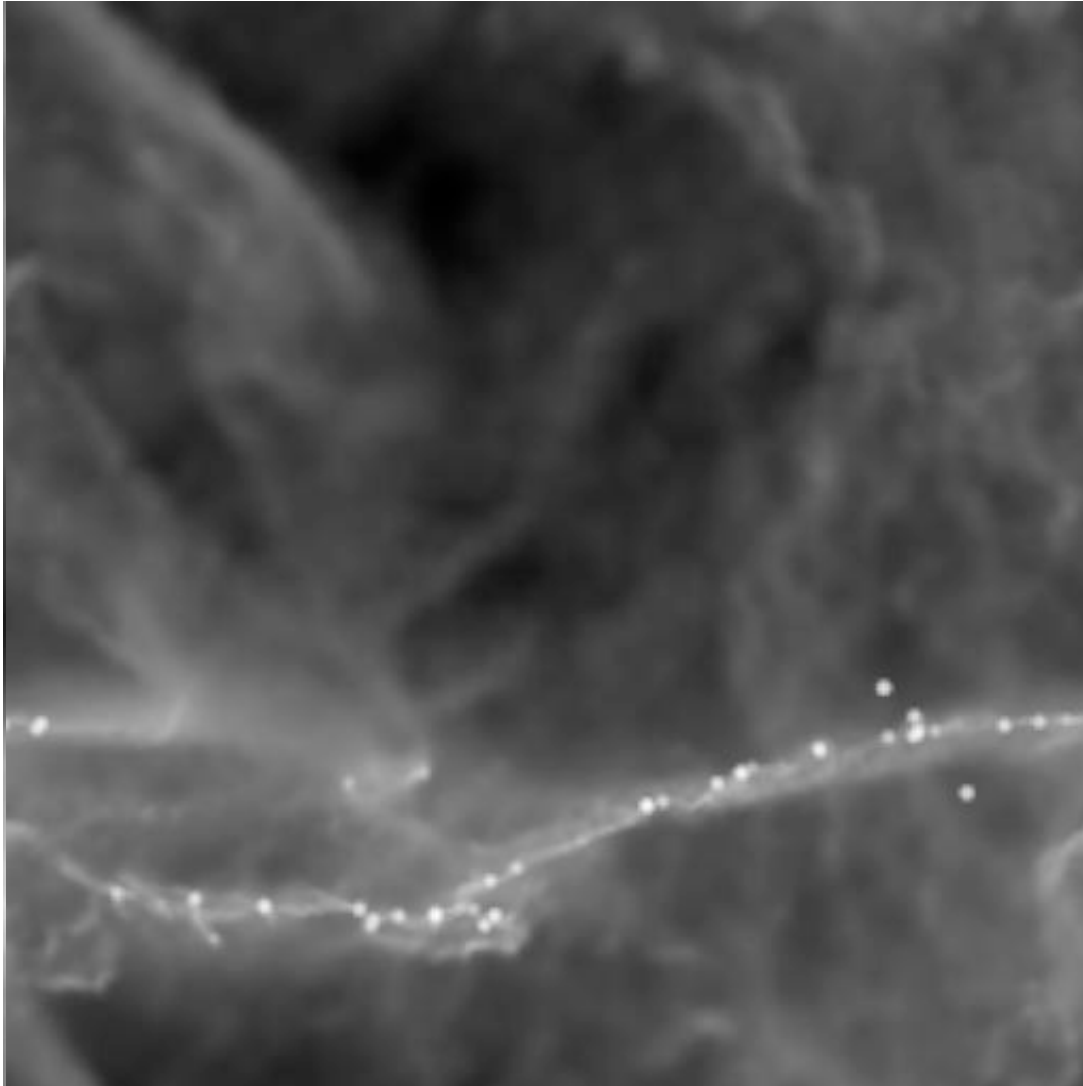
# Transient cloud structure

*Gravoturbulent fragmentation of turbulent self-gravitating clouds*



- SPH model with  $1.6 \times 10^6$  particles
- large-scale driven turbulence
- Mach number  $\mathcal{M} = 6$
- periodic boundaries
- physical scaling: “Taurus”

# Gravoturbulent fragmentation

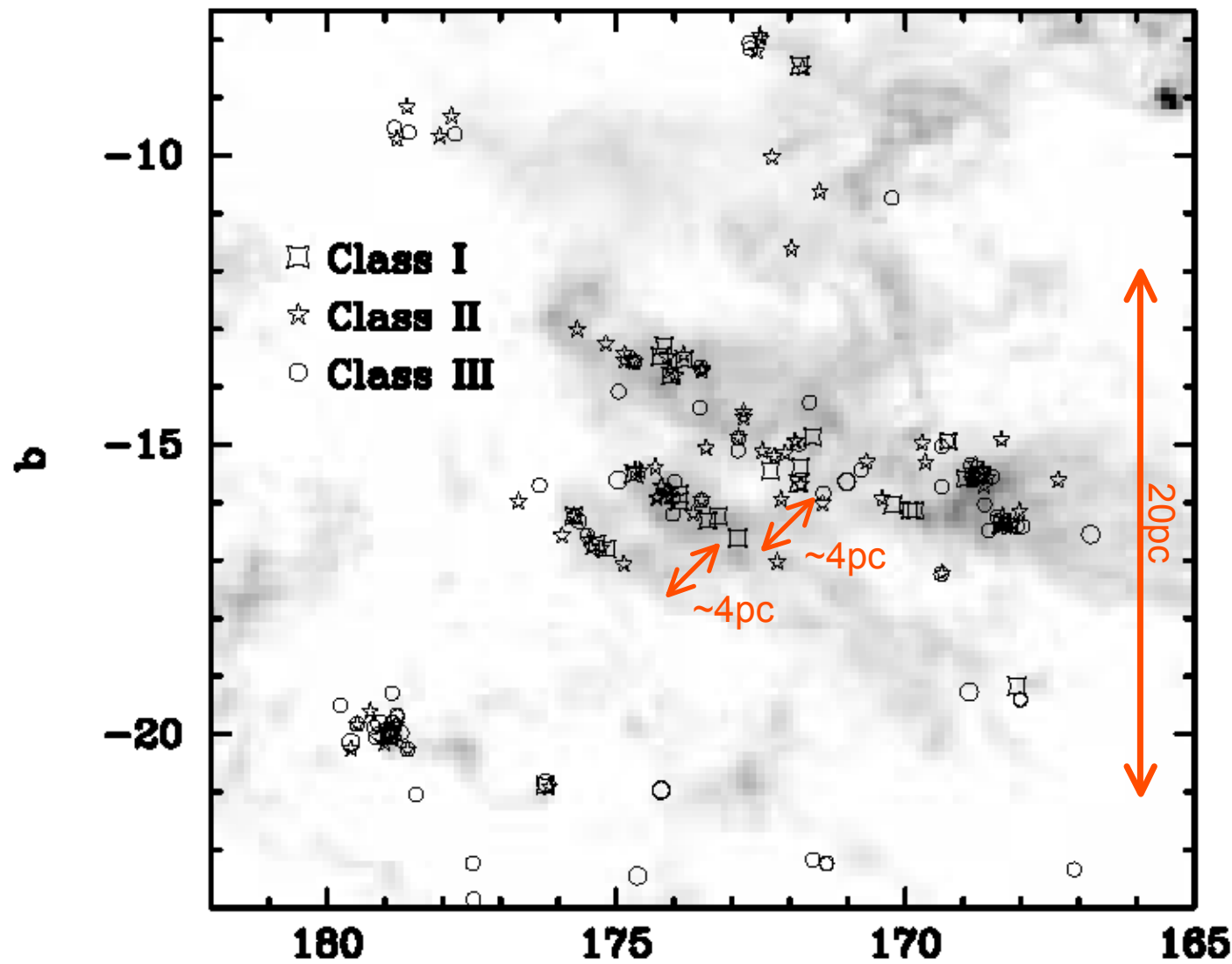


## Gravoturbulent fragmentation in molecular clouds:

- SPH model with  $1.6 \times 10^6$  particles
- large-scale driven turbulence
- Mach number  $\mathcal{M} = 6$
- periodic boundaries
- physical scaling:

### “Taurus”:

- density  $n(\text{H}_2) \approx 10^2 \text{ cm}^{-3}$
- $L = 6 \text{ pc}$ ,  $M = 5000 M_{\odot}$



# Taurus cloud

Star-forming filaments in the *Taurus* molecular cloud

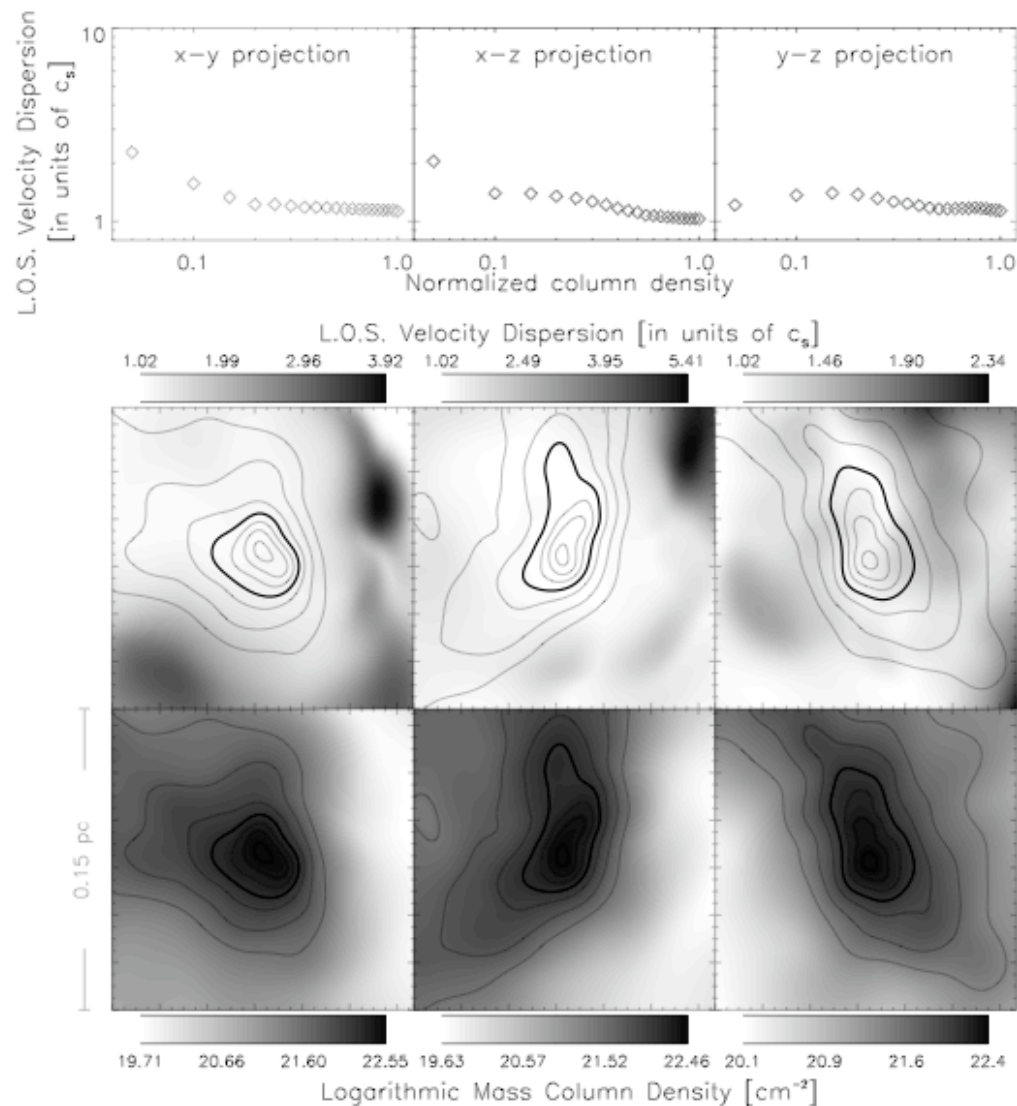
(from Hartmann 2002, ApJ)

# example 2



# Quiescent & coherent cores

(from Klessen et al. 2005, ApJ, 620, 768 - 794; also poster 8415)



correlation between linewidth and column density

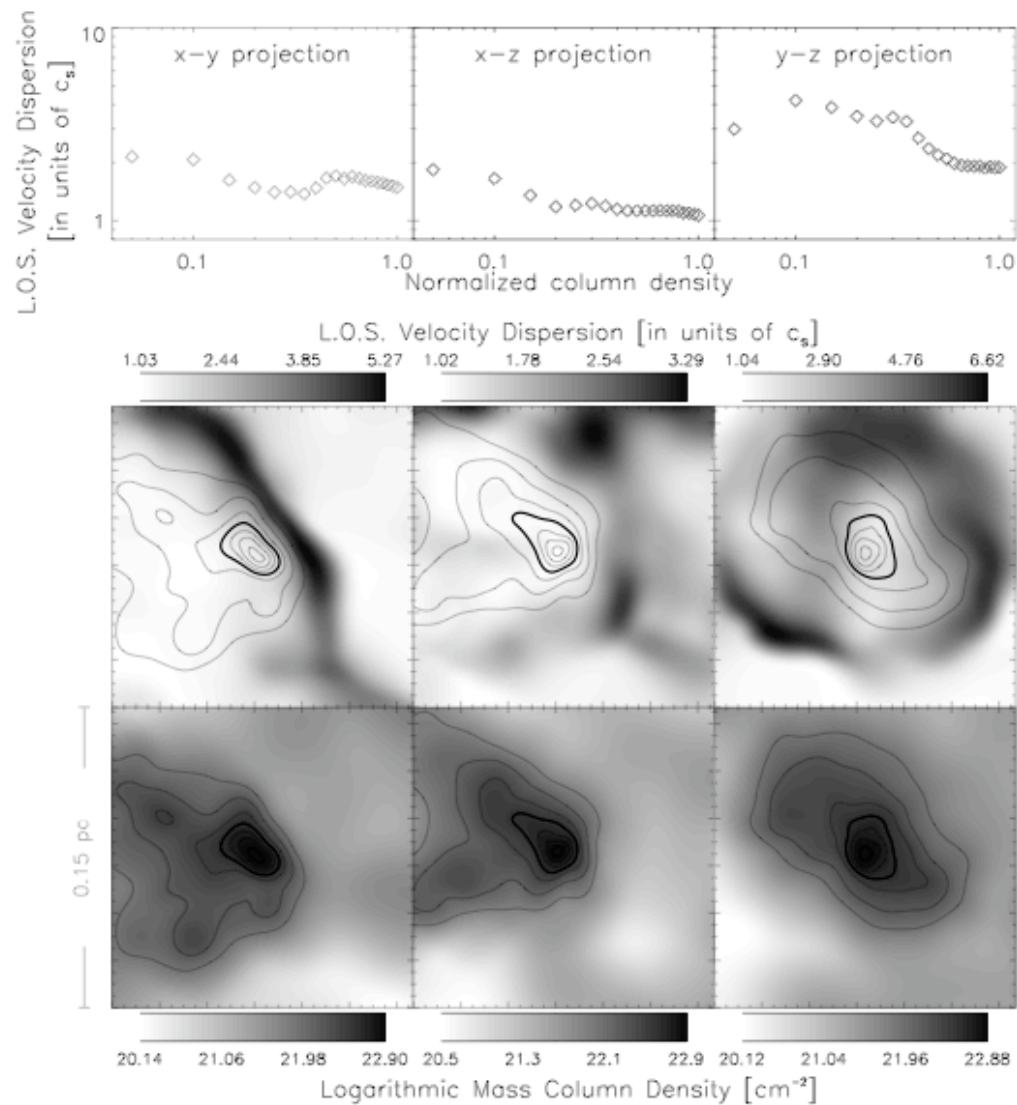
(e.g. Goodman et al. 1998; Barranco & Goodman 1998 Caselli et al. 2002; Tafalla et al. 2004)

map of linewidth (contours column density)

column density map (contours column density)

# Quiescent & coherent cores

(from Klessen et al. 2005, ApJ, 620, 768 - 794; also poster 8415)



correlation between linewidth and column density

(e.g. Goodman et al. 1998; Barranco & Goodman 1998 Caselli et al. 2002; Tafalla et al. 2004)

map of linewidth (contours column density)

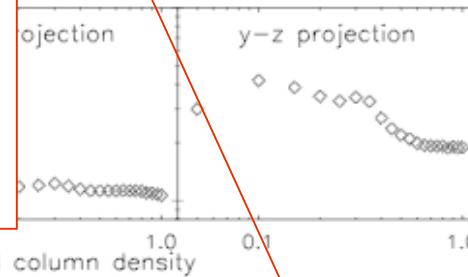
column density map (contours column density)

cores form at stagnation points of convergent large-scale flows

--> often are bounded by ram pressure

--> velocity dispersion highest at boundary

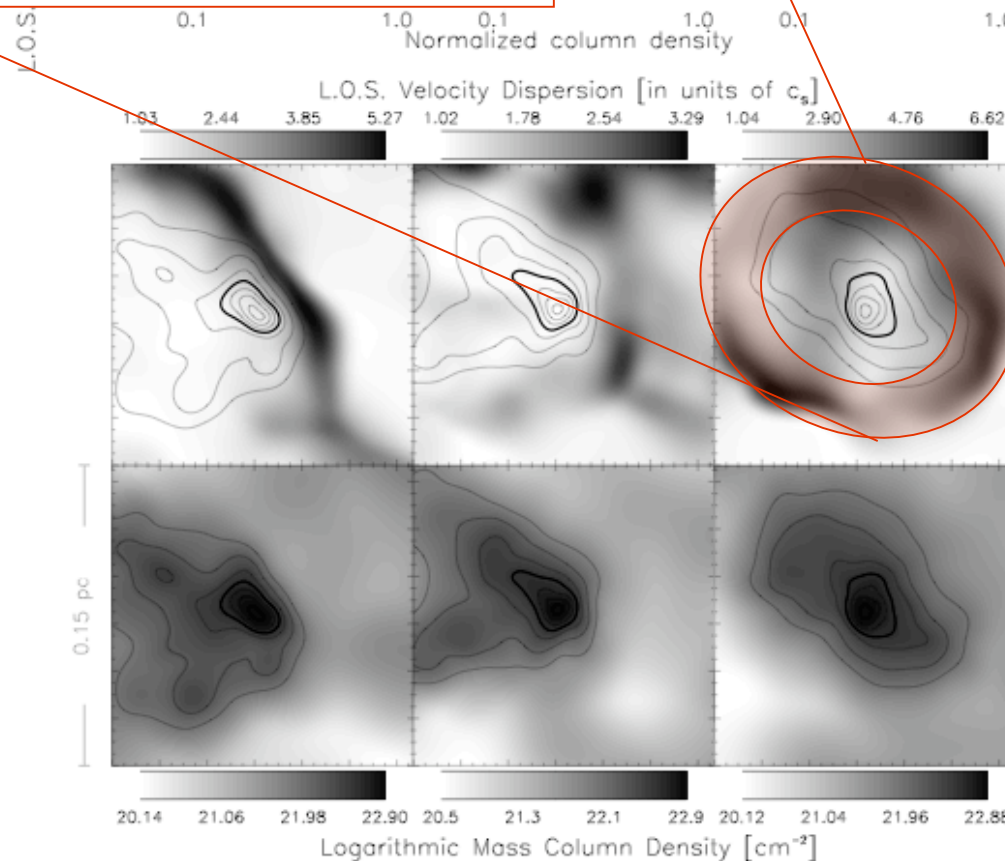
# t & coherent cores



correlation between linewidth and column density

(e.g. Goodman et al. 1998; Barranco & Goodman 1998 Caselli et al. 2002; Tafalla et al. 2004)

(from Klessen et al. 2005, ApJ, 620, 768 - 794; also p...

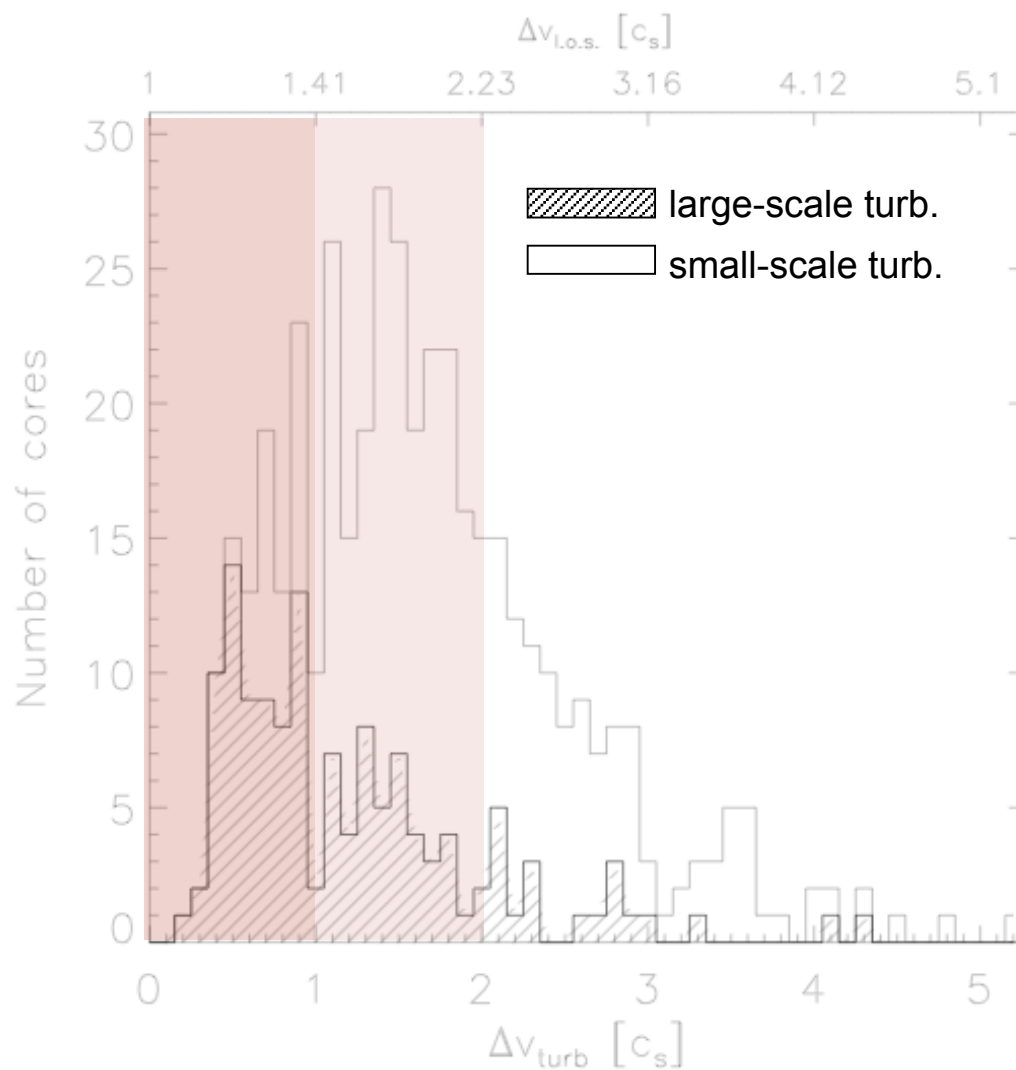


map of linewidth  
(contours column density)

column density map  
(contours column density)

# Quiescent & coherent cores

(from Klessen et al. 2005, ApJ, 620, 768 - 794; also poster 8415)



## Statistics:

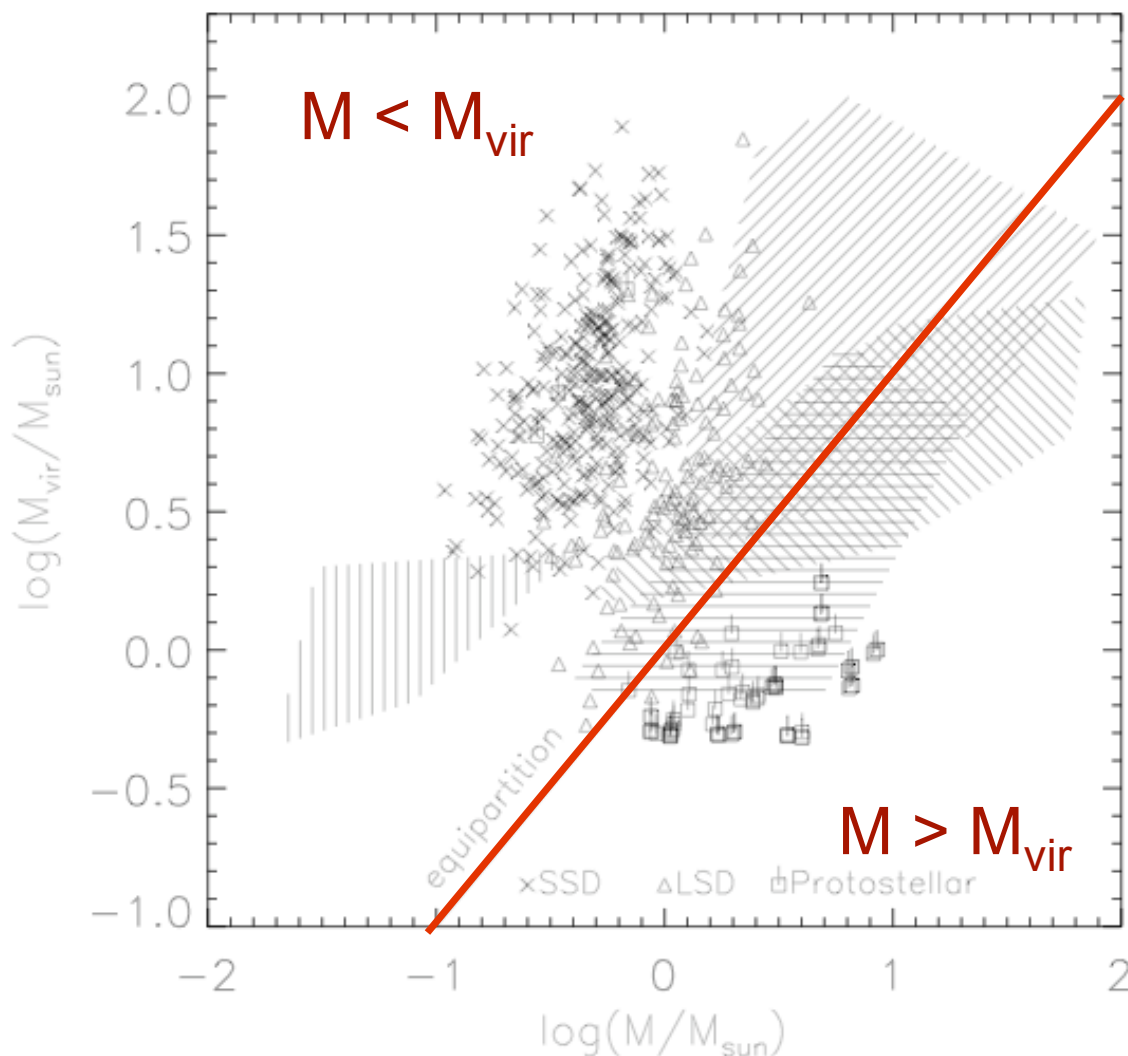
23% of our cores are *quiescent* (i.e. with  $\sigma_{\text{rms}} \leq c_s$ )

48% of our cores are *transonic* (i.e. with  $c_s \leq \sigma_{\text{rms}} \leq 2c_s$ )

half of our cores are *coherent* (i.e. with  $\sigma_{\text{rms}}$  independent of N)

# Quiescent & coherent cores

(from Klessen et al. 2005, ApJ, 620, 768 - 794; also poster 8415)



## Statistics:

most cores have masses smaller than  $M_{\text{vir}}$   
(should reexpand once external compression fades)

some core have more mass than  $M_{\text{vir}}$   
(should collapse)  
(indeed all cores with protostars have  $M > M_{\text{vir}}$ )

# example 3

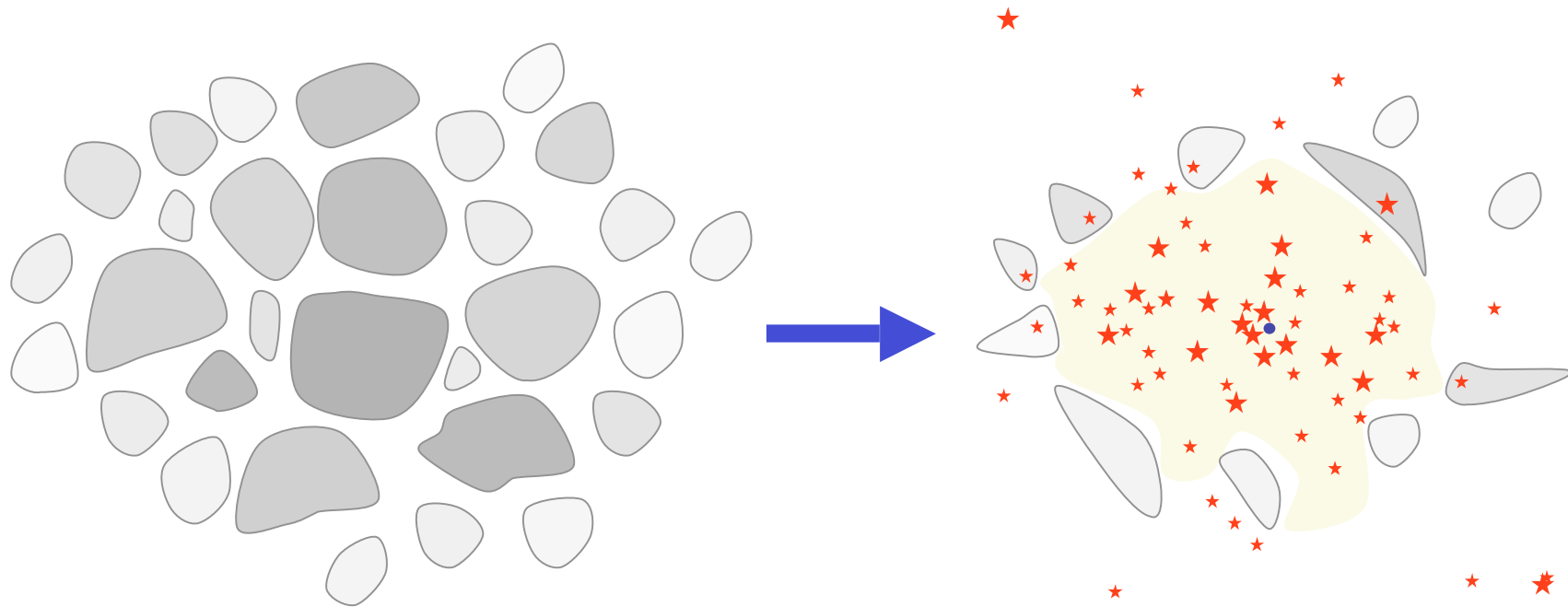
# IMF

- distribution of stellar masses depends on
  - turbulent initial conditions
    - > mass spectrum of prestellar cloud cores
  - collapse and interaction of prestellar cores
    - > competitive accretion and  $N$ -body effects
  - thermodynamic properties of gas
    - > balance between heating and cooling
    - > EOS (determines which cores go into collapse)
  - (proto) stellar feedback terminates star formation
    - ionizing radiation, bipolar outflows, winds, SN

(e.g. Larson 2003, Prog. Rep. Phys.; Mac Low & Klessen, 2004, Rev. Mod. Phys, 76, 125 - 194)

# Star cluster formation

Most stars form in clusters → *star formation = cluster formation*



How to get from **cloud cores** to **star clusters**?

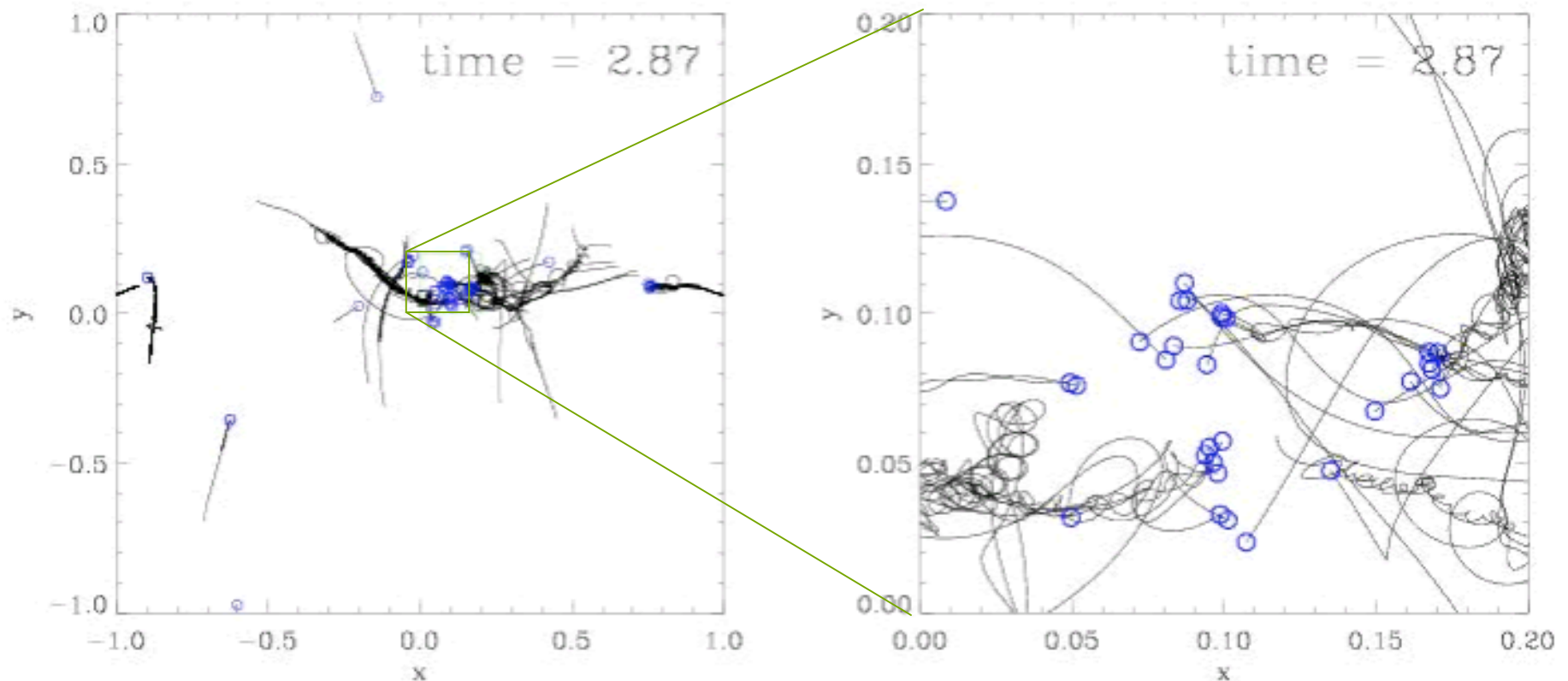
How do the stars **acquire mass**?

(e.g. Larson 2003, Prog. Rep. Phys.; Mac Low & Klessen, 2004, Rev. Mod. Phys, 76, 125 - 194)



# Star cluster formation

Most stars form in clusters → *star formation = cluster formation*

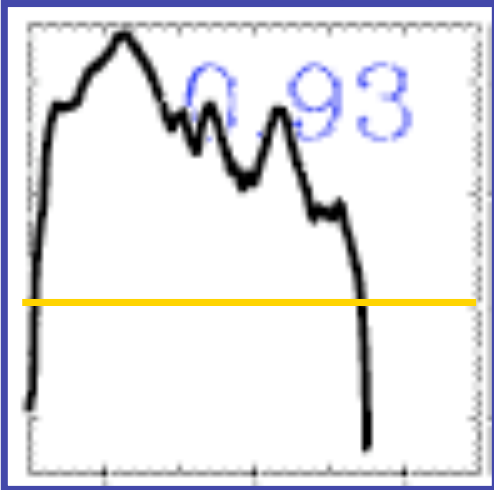
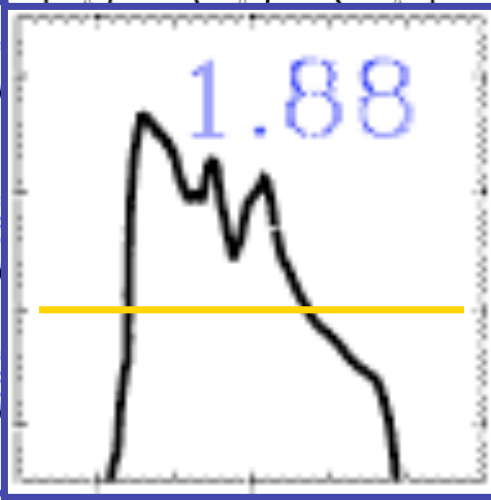
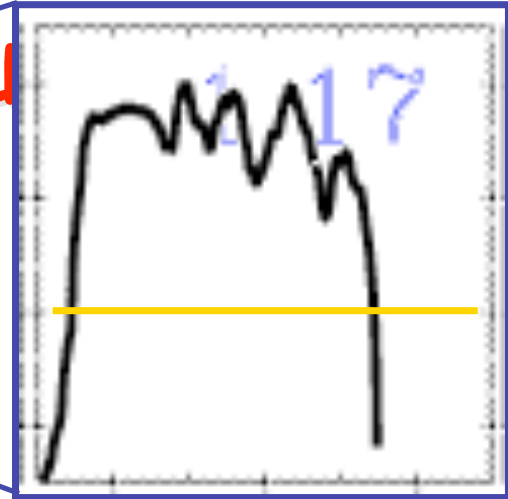
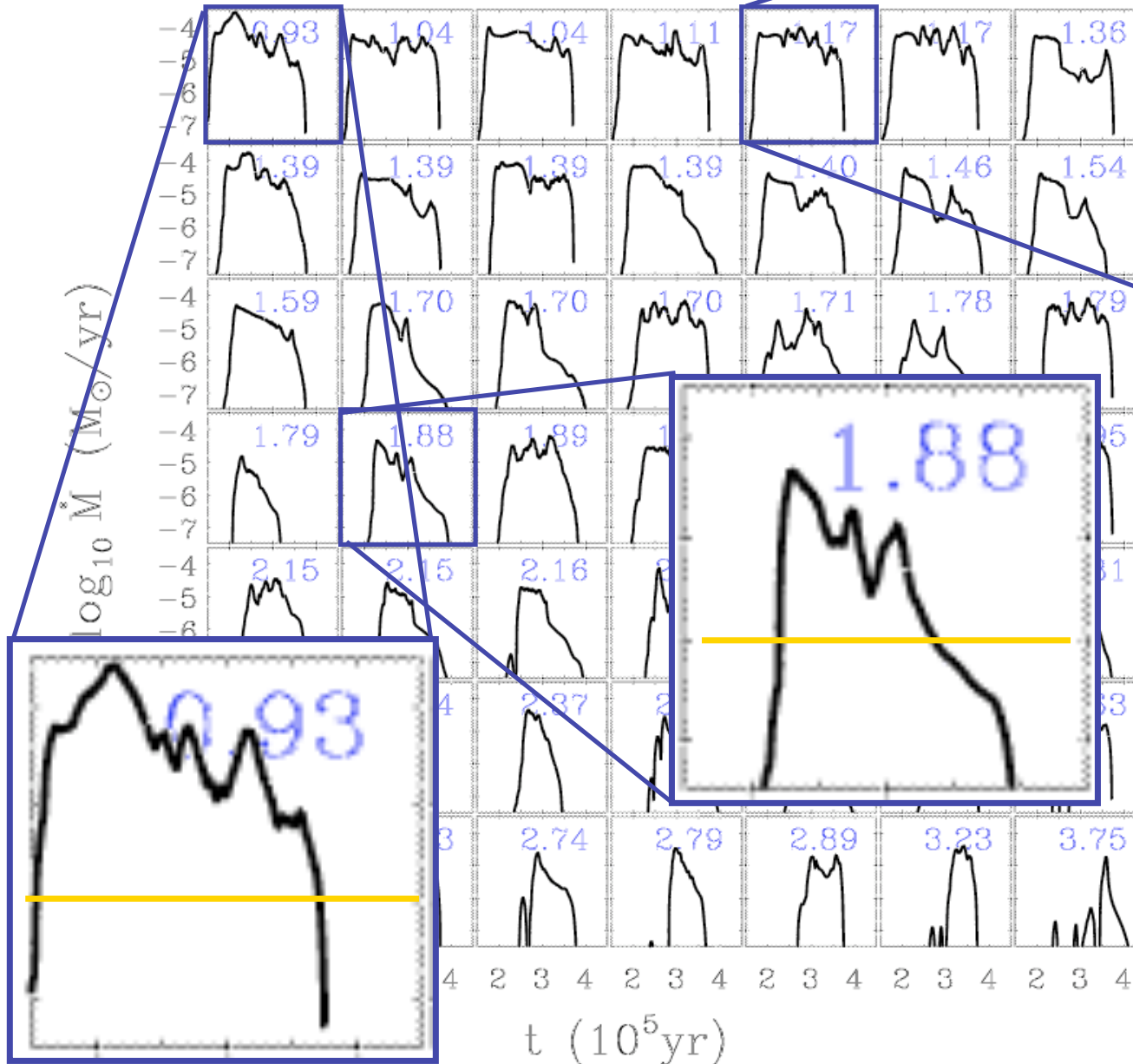


Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation

(from Klessen & Burkert 2000, ApJS, 128, 287)

Ralf Klessen: PPV, Oct. 24, 2005

# Accretion rates in clu



Mass accretion rates *vary with time* and are strongly *influenced* by the *cluster environment*.

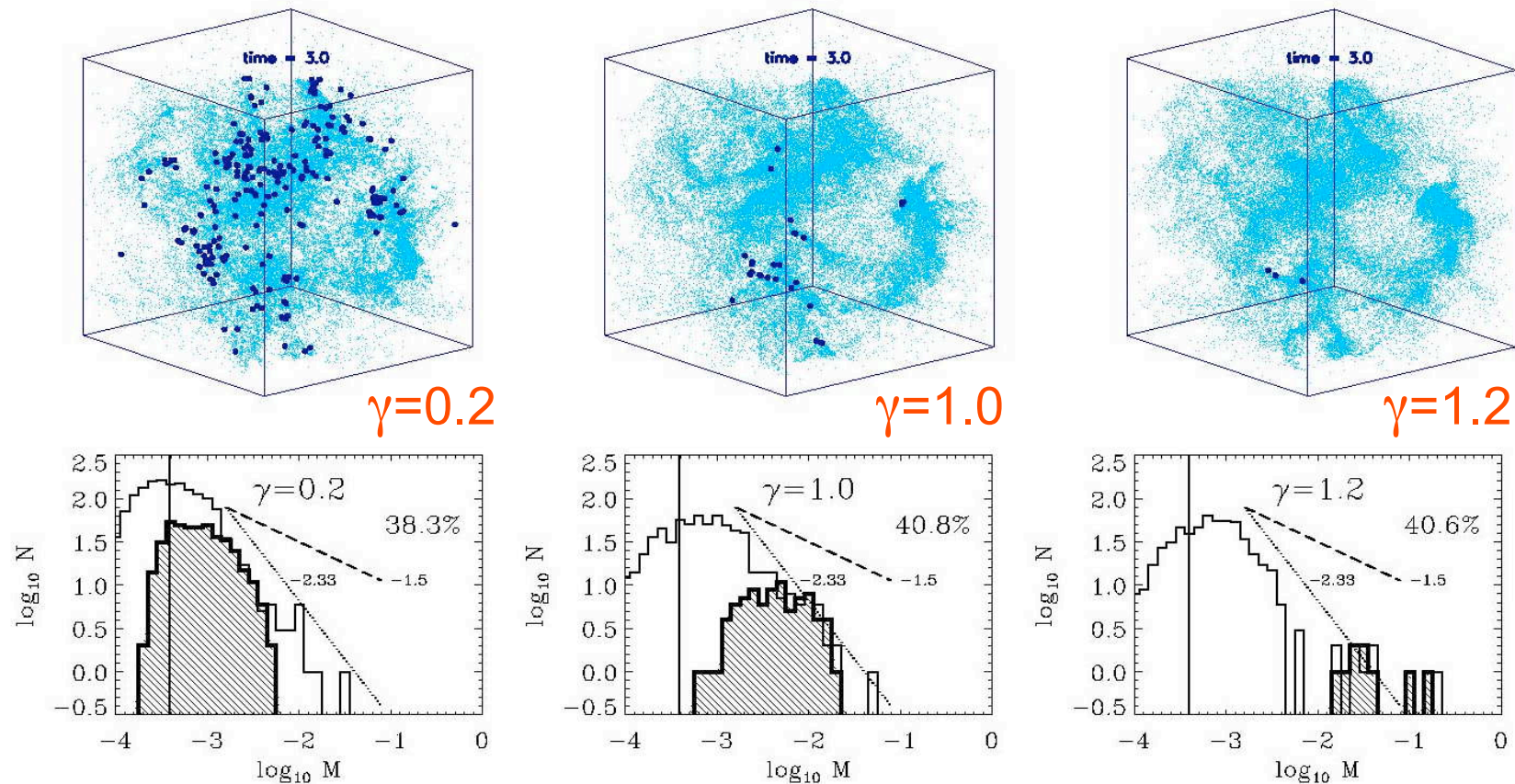
(Klessen 2001, ApJ, 550, L77;  
also Schmeja & Klessen,  
2004, A&A, 419, 405)

# Dependency on EOS

- degree of fragmentation depends on *EOS!*
- polytropic EOS:  $p \propto \rho^\gamma$
- $\gamma < 1$ : dense cluster of low-mass stars
- $\gamma > 1$ : isolated high-mass stars

(see Li, Klessen, & Mac Low 2003, ApJ, 592, 975; also Kawachi & Hanawa 1998, Larson 2003)

# Dependency on EOS



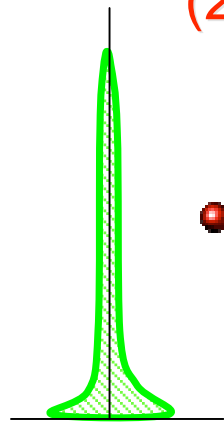
for  $\gamma < 1$  fragmentation is enhanced  $\rightarrow$  *cluster of low-mass stars*  
for  $\gamma > 1$  it is suppressed  $\rightarrow$  formation of *isolated massive stars*

(from Li, Klessen, & Mac Low 2003, ApJ, 592, 975)

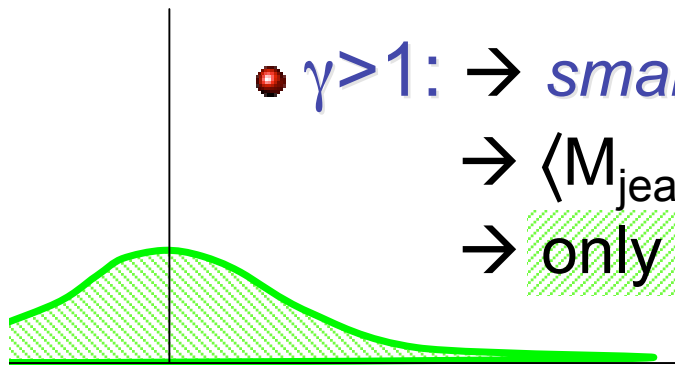
# How does that work?

$$(1) \quad p \propto \rho^\gamma \quad \rightarrow \quad \rho \propto p^{1/\gamma}$$

$$(2) \quad M_{\text{jeans}} \propto \gamma^{3/2} \rho^{(3\gamma-4)/2}$$



- $\gamma < 1$ :  $\rightarrow$  *large* density excursion for given pressure  
 $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  becomes small  
 $\rightarrow$  number of fluctuations with  $M > M_{\text{jeans}}$  is large



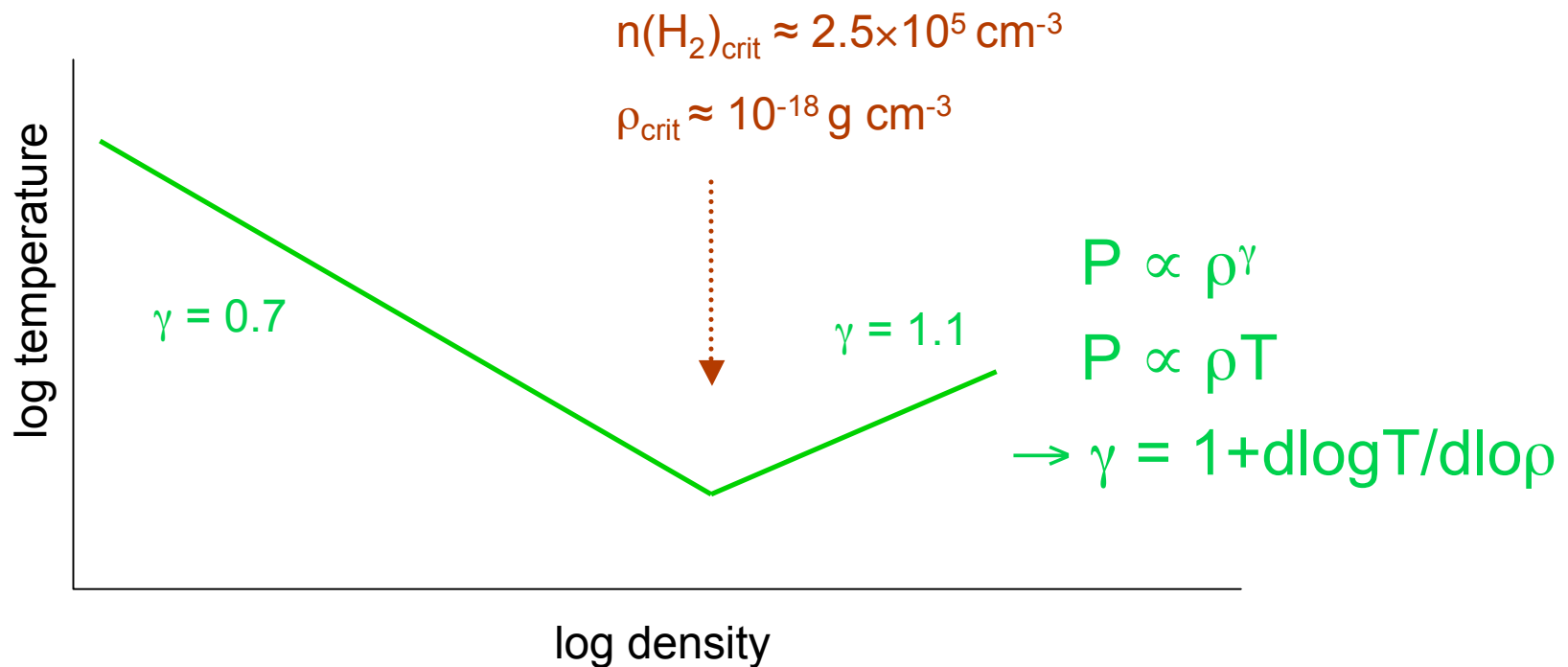
- $\gamma > 1$ :  $\rightarrow$  *small* density excursion for given pressure  
 $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  is large  
 $\rightarrow$  only few and massive clumps exceed  $M_{\text{jeans}}$

# Implications

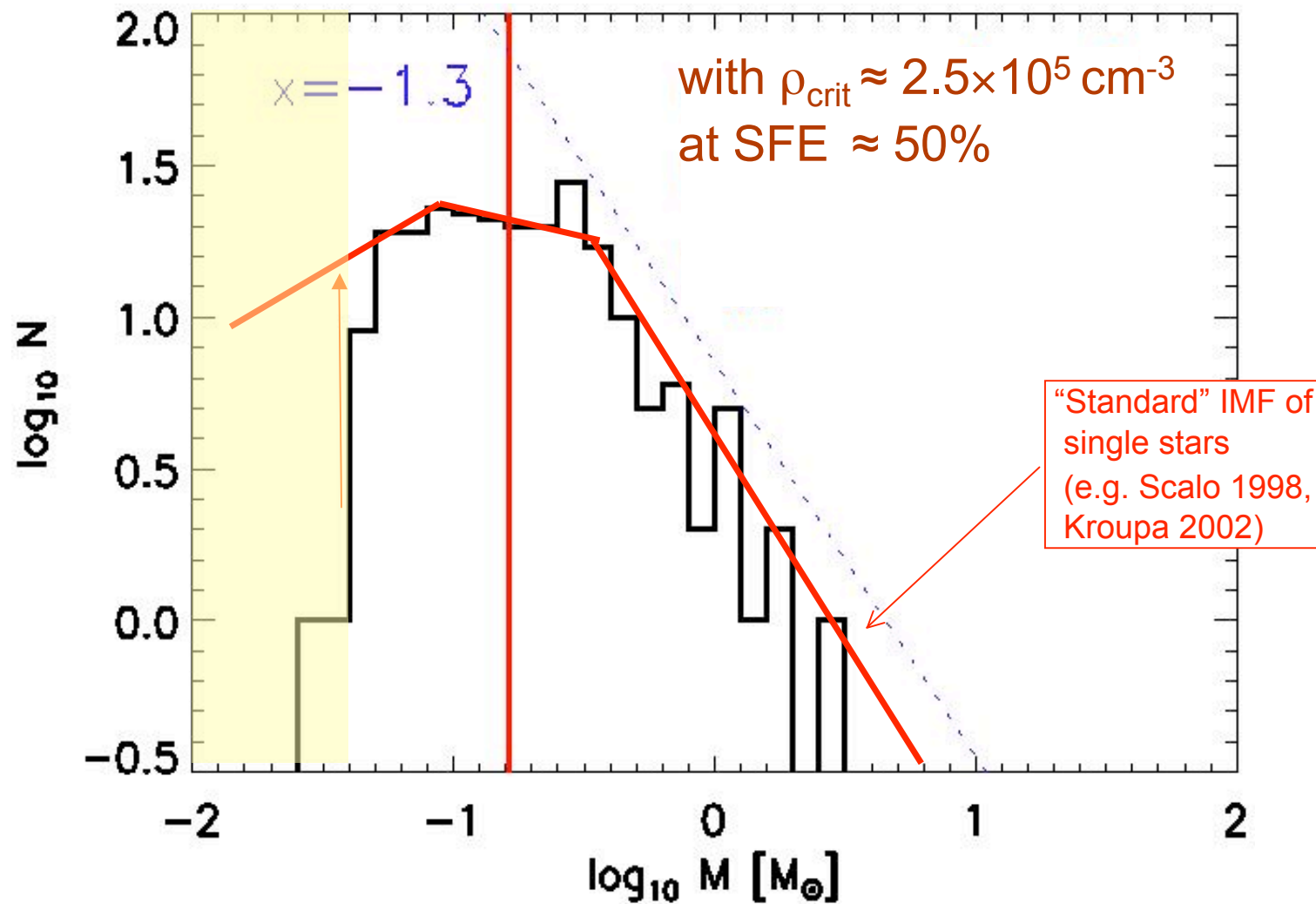
- degree of fragmentation depends on *EOS!*
- polytropic EOS:  $p \propto \rho^\gamma$
- $\gamma < 1$ : dense cluster of low-mass stars
- $\gamma > 1$ : isolated high-mass stars
  - (see Li, Klessen, & Mac Low 2003, ApJ, 592, 975; Kawachi & Hanawa 1998; Larson 2003; also Jappsen, Klessen, Larson, Li, Mac Low, 2005, 435, 611)
- implications for extreme environmental conditions
  - expect Pop III stars to be massive and form in isolation
  - expect IMF variations in warm & dusty starburst regions
    - (Spaans & Silk 2005; Klessen, Spaans, & Jappsen 2005)
- Observational findings: isolated O stars in LMC (and M51)?
  - (Lamers et al. 2002, Massey 2002; see however, de Witt et al. 2005 for Galaxy)

# More realistic EOS

- But EOS depends on *chemical state*, on *balance* between *heating* and *cooling*



# IMF in nearby molecular clouds





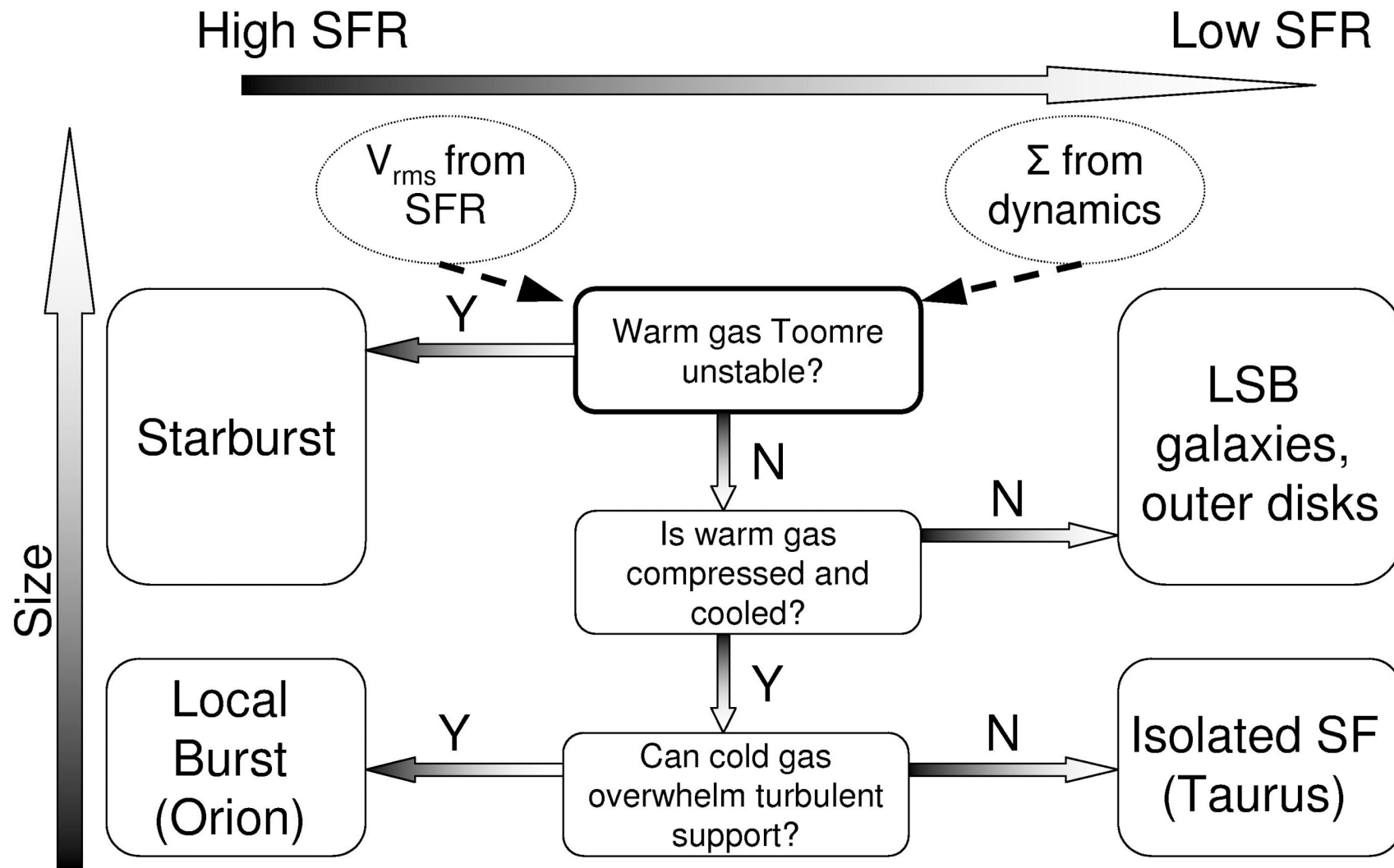
# Summary

# Summary

- interstellar gas is highly inhomogeneous
  - *thermal instability*
  - *gravitational instability*
  - *turbulent compression* (in shocks  $\delta\rho/\rho \approx M^2$ ; in atomic gas:  $M \approx 1...3$ )
- cold *molecular clouds* form rapidly in high-density regions
  - chemical *phase transition*: atomic  $\rightarrow$  molecular
  - process is *modulated* by large-scale *dynamics* in the galaxy
- inside *cold clouds*: turbulence is highly supersonic ( $M \approx 1...20$ )  
 $\rightarrow$  *turbulence* creates density structure, *gravity* selects for collapse  
 $\longrightarrow$  **GRAVOTUBULENT FRAGMENTATION**
- *turbulent cascade*: local compression *within* a cloud provokes collapse
- individual *stars* and *star clusters* form through *sequence* of highly *stochastic* events:
  - *collapse* of cloud cores in turbulent cloud (cores change during collapse)
  - plus mutual *interaction* during collapse (importance depends on ratio of potential energy to turbulent energy) (buzz word: *competitive accretion*)

Thanks!

# SF Flow Chart



(from Mc Low & Klessen, 2004, Rev. Mod. Phys., 76, 125 - 194)