Modeling ISM Dynamics: Star Formation



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Structure

ISM overview

what do we need to study the ISM?

concept of gravoturbulent

fragmentation and star formation.

applications to star cluster formation



















Why study ISM physics?

physical processes

- turbulence theory
- ISM: laboratory for plasm physics
- ISM: laboratory for extreme chemistry

planets

- initial conditions for planet
- formation (chemical composition)
- diversity of planetary systems
- habitability (life)

stars & star clusters

- ISM: environment for star formation
- IMF
- feedback from stars (winds, radiation, SN)
- MC turbulence

extreme environments

- galactic center
- starburst galaxies
- primorial universe

cosmology & galaxy formation

- cooling properties of high-z halos
- primordial star formation
- relation between visible and dark matter





galactic structure & evolution

- chemical enrichment
- global star formation history (Milky Way)
- interrelation between SF and galactic structure





Early dynamical theory

- Jeans (1902): Interplay between self-gravity and thermal pressure
 - stability of homogeneous spherical density enhancements against gravitational collapse
 - dispersion relation:

$$\omega^2 = c_s^2 k^2 - 4\pi G \rho_0$$



• minimal mass:

$$M_J = \frac{1}{6}\pi^{-5/2}G^{-3/2}\rho_0^{-1/2}c_s^3 \propto \rho_0^{-1/2}T^{3/2}$$

full detail in Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125-194 Ralf Klessen: IMPRS School, 25.09.2006



Sir James Jeans, 1877 - 1946

First approach to turbulence

- von Weizsäcker (1943, 1951) and Chandrasekhar (1951): concept of MICROTURBULENCE
 - BASIC ASSUMPTION: separation of scales between dynamics and turbulence

 $\ell_{\rm turb} \ll \ell_{\rm dyn}$

 then turbulent velocity dispersion contributes to effective soundspeed:



S. Chandrasekhar, 1910 - 1995

- $C_c^2 \mapsto C_c^2 + \sigma_{rms}^2$
- \rightarrow Larger effective Jeans masses \rightarrow more stability
- BUT: (1) turbulence depends on k: $\sigma_{rms}^2(k)$

(2) supersonic turbulence \rightarrow usually $\sigma_{rms}^2(k) >> c_s^2$

(full detail in Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125-194) Ralf Klessen: IMPRS School, 25.09.2006

Properties of IMS turbulence

ISM turbulence is:

- Supersonic (rms velocity dispersion >> sound speed)
- Anisotropic (shocks & magnetic field)
- Driven on large scales (power in mol. clouds always dominated by largest-scale modes)

Microturbulent approach is NOT valid in ISM

No closed analytical/statistical formulation known --> necessity for numerical modeling

Problems of early dynamical theory

- Molecular clouds are *highly Jeans-unstable* Yet, they do *NOT* form stars at high rate and with high efficiency.
 (the observed global SFE in molecular clouds is ~5%)
 → something prevents large-scale collapse.
- All throughout the early 1990's, molecular clouds had been thought to be long-lived quasi-equilitrium entities.
- Molecular clouds are *magnetized*.

Magnetic star formation

- Mestel & Spitzer (1956): Magnetic fields can prevent collapse!!!
 - Critical mass for gravitational collapse in presence of B-field

$$M_{cr} = \frac{5^{3/2}}{48\pi^2} \frac{B^3}{G^{3/2}\rho^2}$$

 Critical mass-to-flux ratio (Mouschovias & Spitzer 1976)

$$\left[\frac{M}{\Phi}\right]_{cr} = \frac{\zeta}{3\pi} \left[\frac{5}{G}\right]^{1/2}$$



Lyman Spitzer, Jr., 1914 - 1997

Ambipolar diffusion can initiate collapse
 Ambipolar diffusion can initiate collapse
 Full detail in Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125-194)Ralf Klessen: IMPRS School, 25.09.2006

The "standard theory" of star formation:

- BASIC ASSUMPTION: Stars form from magnetically highly subcritical cores
- Ambipolar diffusion slowly increases (M/ Φ): $\tau_{AD} \approx 10 \tau_{ff}$
- Once $(M/\Phi) > (M/\Phi)_{crit}$: dynamical collapse of SIS
 - Shu (1977) collapse solution
 - $dM/dt = 0.975 c_s^3/G = const.$
- Was (in principle) only intended for isolated, low-mass stars





Problems of magnetic SF

- Observed B-fields are weak, at most marginally critical (Crutcher 1999, Bourke et al. 2001)
- Magnetic fields cannot prevent decay of turbulence (Mac Low et al. 1998, Stone et al. 1998, Padoan & Nordlund 1999)
- Structure of prestellar cores
 (Bacman et al. 2000, e.g. Barnard 68 from Alves et al. 2001)
- Strongly time varying dM/dt (e.g. Hendriksen et al. 1997, André et al. 2000)
- More extended infall motions than predicted by the standard model (Williams & Myers 2000, Myers et al. 2000)



Problems of magnetic SF

- As many prestellar cores as protostellar cores in SF regions (e.g. André et al 2002)
- Molecular cloud clumps seem to be chemically young (Bergin & Langer 1997, Pratap et al 1997, Aikawa et al 2001)
- Stellar age distribution small ($\tau_{\rm ff} << \tau_{\rm AD}$) (Ballesteros-Paredes et al. 1999, Elmegreen 2000, Hartmann 2001)
- Strong theoretical criticism of the SIS as starting condition for gravitational collapse

(e.g. Whitworth et al 1996, Nakano 1998, as summarized in Klessen & Mac Low 2004)

Most stars form as binaries



Observed B-fields are weak

 B versus N(H₂) from
 Zeeman measurements.
 (from Bourke et al. 2001)
 → cloud cores are marginally magnetically supercritical!!!



 $(\Phi/M)_n < 1$ collapse

08/11/04

 $(\Phi/M)_n > 1$ no collapse

Molecular cloud dynamics

• Timescale problem: Turbulence decays on timescales comparable to the free-fall time $\tau_{\rm ff}$

(E \propto t^{- η} with $\eta \approx 1$).

(Mac Low et al. 1998, Stone et al. 1998, Padoan & Nordlund 1999)

 Magnetic fields (static or wavelike) cannot prevent loss of energy.

21. Jan. 2003



Problems of magnetic SF

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Gravoturbulent star formation

• more recent suggestions:

Star formation is controlled by interplay between gravity and supersonic turbulence + B-field !

• dual role of turbulence:

- stability on large scales
- initiating collapse on small scales



(e.g., Larson, 2003, Rep. Prog. Phys, 66, 1651; or Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125)

- B-field:
 - adds stability
 - on all scales

Gravoturbulent star formation

- interstellar gas is highly inhomogeneous
 - thermal instability
 - gravitational instability



- *turbulent compression* (in shocks $\delta \rho / \rho \propto M^2$; in atomic gas: $M \approx 1...3$)
- cold molecular clouds can form rapidly in high-density regions at stagnation points of convergent large-scale flows
 - chemical phase transition: atomic → molecular
 - process is *modulated* by large-scale *dynamics* in the galaxy
- inside *cold clouds:* turbulence is highly supersonic ($M \approx 1...20$)
 - → turbulence creates large density contrast, gravity selects for collapse

GRAVOTUBULENT FRAGMENTATION

 turbulent cascade: local compression within a cloud provokes collapse → formation of individual stars and star clusters



(e.g. Mac Low & Klessen, 2004, Rev. Mod. Phys., 76, 125-194)



Properties of turbulence

 laminar flows turn *turbulent* at *high Reynolds* numbers

 $\mathsf{Re} = \frac{\mathsf{advection}}{\mathsf{dissipation}} = \frac{VL}{v}$



V= typical velocity on scale L, $\nu = viscosity$, Re > 1000

vortex streching --> turbulence
 is intrinsically anisotropic
 (only on large scales you may get
 homogeneity & isotropy in a statistical sense;
 see Landau & Lifschitz, Chandrasekhar, Taylor, etc.)

(ISM turbulence: shocks & B-field cause additional inhomogeneity)





Vortex Formation



Vortices are streched and folded in three dimensions



Turbulent cascade



Turbulent cascade





Turbulent cascade in ISM



L ≈ 0.1 pc

(supernovae, winds, spiral density waves?)

molecular diffusion?)



Density structure of MC's



molecular clouds are highly inhomogeneous

stars form in the densest and coldest parts of the cloud

 $\rho\text{-Ophiuchus}$ cloud seen in dust emission

let's focus on a cloud core like this one

Evolution of cloud cores





- Does core form single massive star or cluster with mass distribution?
- Turbulent cascade "goes through" cloud core --> NO scale separation possible
 - --> NO effective sound speed
- Turbulence is supersonic!
 - --> produces strong density contrasts: $\delta \rho / \rho \approx M^2$
 - --> with typical $M \approx 10 --> \delta \rho / \rho \approx 100!$
- many of the shock-generated fluctuations are Jeans unstable and go into collapse
- --> core breaks up and forms a cluster of stars

Evolution of cloud cores



Formation and evolution of cores

 protostellar cloud cores form at the stagnation points of convergent turbulent flows



- if M > M_{Jeans} $\propto \rho^{-1/2} T^{3/2}$: collapse and star formation
- if M < $M_{Jeans} \propto \rho^{-1/2} T^{3/2}$: reexpansion after external compression fades away (e.g. Vazquez-Semadeni et al 2005)
- typical timescales: $t \approx 10^4 \dots 10^5$ yr
- because *turbulent* ambipolar diffusion time is *short*, this time estimate still holds for the presence of magnetic fields, in *magnetically critical cores*

(e.g. Fatuzzo & Adams 2002, Heitsch et al. 2004)



Formation and evolution of cores

What happens to distribution of cloud cores?



Two externe cases:

- (1) turbulence dominates energy budget: $\alpha = E_{kin} / |E_{pot}| > 1$
 - --> individual cores do not interact
 - --> collapse of individual cores dominates stellar mass growth
 - --> loose cluster of low-mass stars
- (2) turbulence decays, i.e. gravity dominates: $\alpha = E_{kin} / |E_{pot}| < 1$
 - --> global contraction
 - --> core do interact while collapsing
 - --> *competition* influences *mass growth*
 - --> dense cluster with high-mass stars





turbulence creates a hierarchy of clumps





as turbulence decays locally, contraction sets in





as turbulence decays locally, contraction sets in




while region contracts, individual clumps collapse to form stars





while region contracts, individual clumps collapse to form stars













individual clumps collapse to form stars



 $\alpha = E_{kin} / |E_{pot}| < 1$

in *dense clusters*, clumps may merge while collapsing --> then contain multiple protostars





in *dense clusters*, clumps may merge while collapsing --> then contain multiple protostars





in *dense clusters*, clumps may merge while collapsing --> then contain multiple protostars











in dense clusters, N-body effects influence mass growth











result: star cluster, possibly with HII region











magneto-hydrodynamics

(multi-phase, non-ideal MHD, turbulence)

chemistry (gas + dust, heating + cooling)

radiation (continuum + lines)

stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)

stellar evolution (feedback: radiation, winds, SN)

laboratory work

(reaction rates, cross sections, dust coagulation properties, etc.)

massive parallel codes

- particle-based: SPH with improved algorithms (XSPH with turb. subgrid model, GPM, particle splitting, MHD-SPH?)
- grid-based: AMR (FLASH, ENZO, RAMSES, Nirvana3, etc), subgrid-scale models (FEARLESS)

BGK methods

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magneto-hydrodynamics (multi-phase, non-ideal MHD, turbulence)

- ever increasing chemical networks
- working reduced networks for time-dependent chemistry in combination with hydrodynamics
- improved data on reaction rates (laboratory + quantum mechanical calculations)

chemistry (gas + dust, heating + cooling)

radiation (continuum + lines)

stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)



magneto-hydrodynamics (multi-phase, non-ideal MHD, turbulence)

chemistry (gas + dust, heating + cooling)

• continuum vs. lines

- Monte Carlo, characteristics
- approximative methods

• combine with hydro

radiation (continuum + lines)

stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)



magneto-hydrodynamics (multi-phase, non-ideal MHD, turbulence)

chemistry (gas + dust, heating + cooling)

 statistics: number of stars (collisional: 10⁶,

- collisionless: 10¹⁰)
- transition from gas to stars
- binary orbits
- long-term integration

K

radiation (continuum + lines)

stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)



magneto-hydrodynamics (multi-phase, non-ideal MHD, turbulence)

chemistry (gas + dust, heating + cooling)

 very early phases (pre main sequence tracks)

- massive stars at late phases
- role of rotation
- primordial star formation

radiation (continuum + lines)

stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)





magneto-hydrodynamics

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chemistry (gas + dust, heating + cooling)

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stellar dynamics (collisional: star clusters, collisionless: galaxies, DM)

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Three examples















modeling star formation in galactic disk

(hydrodynamics, stellar dynamics) (Schmidt law, star-formation history, relation between global dynamics and SF)

molecular cloud formation

(hydrodynamics, chemistry, feedback [radiation, outflows]) (molecular cloud dynamics, star cluster formation)

modeling properties of prestellar cores

(MHD, chemistry, radiation) (initial conditions of star formation, IMF, multiplicity, planet formation, etc.)

Correlation between H_2 and HI



⁽Deul & van der Hulst 1987, Blitz et al. 2004)

Correlation with large-scale perturbations



density/temperature fluctuations in warm atomar ISM are caused by *thermal/gravitational instability* and/or *supersonic turbulence*

some fluctuations are dense enough to form H₂ within "reasonable time" → molecular cloud

external perturbuations (i.e. potential changes) *increase* likelihood



Modeling galactic SF



SPH + stars + DM models of isolated disk galaxies with several million particles

→ begin to resolve
individual molecular
clouds
→ we need to care
about "small-scale"
physics (i.e. *transition*from *atomic gas* to *molecular*)

(simple physics: gravity + hydrodynamics (isothermal EOS) + stellar dynamics [stars + DM])



Result: gravitational instability alone leads to the *Schmidt law* (power-law correlation between star formation and surface denstiy)

 $\Sigma_{
m SFR} \propto \Sigma_{
m gas}^{1.5}$

Molecular cloud formation

- ... in convergent large-scale flows
- ... setting up the turbulent cascade
 - Mach 3 colliding flow
 - Vishniac instability + thermal instability
 - compressed sheet
 breaks up and builds
 up cold, high-density
 "blobs" of gas
 - --> molecular cloud formation
 - cold cloud motions correspond to
- supersonic turbulence



(e.g. Koyama & Inutsuka 2002, Heitsch et al., 2005, Vazquez-Semadeni et al. 2004)









(de Avillez & Breitschwerdt)

n

Ralf Klessen: IMPRS School, 25.09.2006

Τ

ISM: transition HI to H_2

y (pc)

consistent models of ISM dynamics require to go beyond the simple models!

- magnetohydrodynamics (account for large-scale dynamics
 - + turbulence)
- time-dependent chemistry (reduced network, focus on few dominant species, e.g. H_2)
- radiation (currently simple assumptions)

H2 forms rapidly in shocks / transient density fluctuations / H2 gets destroyed slowly in low density regions / result: turbulence greatly enhances H2formation rate



(Glover & Mac Low 2006ab:)



Reduced chemical network

Table 1. The set of chemical reactions that make up our model of non-equilibrium hydrogen chemistry.

Table 2. Processes included in our thermal model.

Reaction	Reference	Process	References
1. $H + H + grain \rightarrow H_2 + grain$	Hollenbach & McKee (1979)	Cooling:	
		CII fine structure lines	Atomic data – Silva & Viegas (2002)
2. $H_2 + H \rightarrow 3H$	Mac Low & Shull (1986) (low density),		Collisional rates (H ₂) – Flower & Launay (1977)
	Lepp & Shull (1983) (high density)		Collisional rates (H, $T < 2000$ K) – Hollenbach & McKee (1989)
			Collisional rates (H, $T > 2000$ K) – Keenan <i>et al.</i> (1986)
3. $H_2 + H_2 \rightarrow 2H + H_2$	Martin, Keogh & Mandy (1998) (low densit		Collisional rates (e ⁻) – Wilson & Bell (2002)
	Shapiro & Kang (1987) (high density)	Offine structure lines	Atomic data – Silva & Viegas (2002)
			Collisional rates (H, H ₂) – Flower, priv. comm.
4. $H_2 + \gamma \rightarrow 2H$	See § 2.2.1		Collisional rates (e ⁻) – Bell, Berrington & Thomas (1998)
			Collisional rates (H ⁺) – Pequignot (1990, 1996)
5. $H + c.r. \rightarrow H^+ + e$	Liszt (2003)	Si II fine structure lines	Atomic data – Silva & Viegas (2002)
			Collisional rates (H) – Roueff (1990)
6. $H + e \rightarrow H^+ + 2e$	Abel et al. (1997)		Collisional rates (e ⁻) – Dufton & Kingston (1991)
		H ₂ rovibrational lines	Le Bourlot, Pineau des Forêts & Flower (1999)
7. $H^+ + e \rightarrow H + \gamma$	Ferland <i>et al.</i> (1992)	Gas-grain energy transfer ¹	Hollenbach & McKee (1989)
		Recombination on grains	Wolfire et al. (2003)
8. $H^+ + e + grain \rightarrow H + grain$	Weingartner & Draine (2001)	Atomic resonance lines	Sutherland & Dopita (1993)
		H collisional ionization	Abel et al. (1997)
		H ₂ collisional dissociation	See Table 1

here: e^- , H^+ , H, H_2 in primordial gas we do: e^- , H^+ , H, H^- , H_2^+ , H_2 , C, C⁺, O, O⁺

Heating: Photoelectric effect Bakes & Tielens (1994); Wolfire et al. (2003) H₂ photodissociation Black & Dalgarno (1977) UV pumping of H₂ Burton, Hollenbach & Tielens (1990) H₂ formation on dust grains Hollenbach & McKee (1989) Cosmic ray ionization Goldsmith & Langer (1978)







L = 40 pc, $n_0 = 100$ cm-3, $B_0 = 5.85$ mG, $v_{rms} = 0.0$ (Glover & Mac Low 2006a)



Gravitational collapse within MCs



Gravitational collapse within MCs

immediate future: SPH with radiation feedback (first validation runs)




IRAM Observations N_2H^+ and $C^{18}O$ -15 arcsecond res. --3000 AU N_2H^+ dense gas tracer -Most SCUBA -Few extinction!





















adaptive mesh refinement:

computational grid gets refined in regions of high interest (e.g. protostellar cores)

formation of 30 M_{sun} core in turbulent molecular cloud (FLASH with appropriate cooling curve module and turbulent driving)





These 2D snapshots show the onset of the **large scale outflow**. After ca. 70.000 years into the collapse a strong toroidal magnetic field builds up whose magnetic pressure reverses the gas flow and drives an outflow (time difference between these snapshots: 1400 years).





Initially a magnetic field aligned with the rotation axis of the cloud core threads the entire simulation box. The field strength varies slightly (3.4 – 14 micro Gauss) along the equatorial plane to maintain a constant plasma b = 8pp/B2. In this configuration, prior to the gravitational collapse the sphere loses a considerable amount of angular momentum from 'magnetic braking' (Mouschovias & Paleologou, 1980).





The 3D structure of the magnetic field line configuration in the **jet** launching 2×10^{13} cm region.

As predicted by analytics the magnetocentrifugally driven disk jet is faster in the inner region (dark red) than further away from the outflow axis (light red).



 $C^{18}O$ (1-0)









Protostellar collapse







Protostellar collapse







Protostellar collapse









What do we need to study ISM?



magneto-hydrodynamics

(multi-phase, non-ideal MHD, turbulence)

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