Possible Lessons from Present-Day Star Formation



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- some phenomenology
- processes that influence present-day SF and their possible relevance for high-z SF:
 - turbulence
 - thermodynamics
 - magnetic fields
 - feedback

example: Orion



lets look at the Orion Nebula Cluster (ONC)

hydrogen H₂

(radio emission --

example: Orion



Orion molecular cloud

The Orion molecular cloud is the birthplace of several young embedded star clusters.

The Trapezium cluster is only visible in the IR and contains about 2000 newly born stars.



Trapezium cluster

example: Orion



NGC 602 in LMC



NGC 602 in the LMC: Hubble Heritage Image

end of formation phase: star cluster with HII region

stellar mass function







Orion, NGC 3603, 30 Doradus (Zinnecker & Yorke 2007, ARAA, 45, 481)

nearby molecular clouds









image from Alyssa Goodman: COMPLETE survey



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- some phenomenology
- Processes that influence present-day SF and their possible relevance for high-z SF:

turbulence plus cluster environment

- thermodynamics
- magnetic fields

example: model of Orion

"model" of Orion cloud: 15.000.000 SPH particles, 10⁴ M_{sun} in 10 pc, mass resolution 0,02 M_{sun}, forms ~2.500 "stars" (sink particles)

isothermal EOS, top bound, bottom unbound

has clustered as well as distributed "star" formation

efficiency varies from 1% to 20%

develops full IMF (distribution of sink particle masses)



(calculation by Ian Bonnell & Paul Clark)

example: model of Orion

"model" of Orion cloud: 15.000.000 SPH particles, 10⁴ M_{sun} in 10 pc, mass resolution 0,02 M_{sun}, forms ~2.500 "stars" (sink particles)

MASSIVE STARS

- form early in high-density gas clumps (cluster center)
- high accretion rates, maintained for a long time

LOW-MASS STARS

- form later as gas falls into potential well
- high relative velocities
- little subsequent accretion



(calculation by Ian Bonnell & Paul Clark)

in dense clusters protostellar interaction may be come important!



Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation (from Klessen & Burkert 2000, ApJS, 128, 287)

turbulence leads to fragmentation

this is true even for Z=0 (see talk by Paul Clark)



FIG. 1.— The column density images show the state of the POPIII.1 clouds after they have accreted 10 percent of their mass (100 M_{\odot}). The sink particles are denoted in the images by the white dots. In our lowest level of turbulence in which $\Delta \nu = 0.1c_s$, the velocities contain enough angular momentum to produce a disk around the sink particle. However once the strength of the initial velocities increased – to as little as 0.2 times the initial sound speed in the gas – the turbulence leads to fragmentation of the gas. At the point at which the simulations are shown here, the 0.1, 0.2, 0.4 and 0.8 c_s have produced 1, 6, 34 and 17 sink particles respectively.

FIG. 2.— Similar to Fig 1, the column density images show the state of the POPIII.2 clouds after they have accreted 10 percent of their mass (15 M_☉). Note that the scale in this figure differs from that in Fig 1. The clouds exhibit a different behaviour from their POPIII.1 counterparts. Although the clouds all form disks around their sink particles, due to the angular momentum contained in the initial turbulent motions, only the $\Delta v = 0.8c_s$ turbulent cloud has undergone fragmentation by this point in the evolution.

turbulence leads to fragmentation

Θ this is true even for Z=0

(see talk by Paul Clark)



FIG. 1.— The column density images show the state of the POPIII.1 clouds after they have accreted 10 percent of their mass (100 M_{\odot}). The sink particles are denoted in the images by the white dots. In our lowest level of urbulence in which $\Delta \nu = 0.1c_s$, the velocities contain enough angular momentum to produce a disk around the sink particle. However once the strength of the initial velocities is increased – to as little as 0.2 times the initial sound speed in the gas – the turbulence leads to fragmentation of the gas. At the point at which the simulations are shown here, the 0.1, 0.2, 0.4 and 0.8 c_s have produced 1, 6, 34 and 17 sink particles respectively.



turbulence leads to fragmentation

- this is true even for Z=0 (see talk by Paul Clark)
- full stellar mass range from brown dwarf regime onwards?
- it could be that Pop III.2 stars are more massive than Pop III.1
- key questions: which processes could prevent or weaken fragmentation?





- some phenomenology
- processes that influence present-day SF and their possible relevance for high-z SF:
 - ♀ turbulence

thermodynamics (balance heating / cooling)

- magnetic fields
- Ieedback

EOS as function of metallicity



⁽Omukai et al. 2005, ApJ, 626, 627)

present-day star formation



(Omukai et al. 2005, ApJ, 626, 627, Jappsen et al. 2005, A&A, 435, 611, Larson 2005, MNRAS, 359, 211)

present-day star formation





(Jappsen et al. 2005, A&A, 435, 611,)





⁽Jappsen et al. 2005, A&A, 435, 611,)



- combine scale free process \rightarrow POWER LAW BEHAVIOR
 - turbulence (Padoan & Nordlund 2002, Hennebelle & Chabrier 2008, 2009)
 - gravity in dense clusters (Bonnell & Bate 2006, Klessen 2001)
 - universality: dust-induced EOS kink insensitive to radiation field (Elmegreen et al. 2008)
- with highly stochastic processes → central limit theorem
 → GAUSSIAN DISTRIBUTION
 - basically mean thermal Jeans length (or feedback)
 - universality: insensitive to metallicity (Clark et al. 2010, submitted)

transition Pop III to Pop II.5



⁽Omukai et al. 2005, ApJ, 626, 627)

dependence on Z at low density

• at densities $n < 10^2$ cm⁻³ and metallicities $Z < 10^{-2}$ H_2 cooling dominates behavior.

fragmentation depends on *initial conditions*

- example 1: solid-body rotating top-hat initial conditions with dark matter fluctuations (a la Bromm et al. 1999) fragment no matter what metallicity you take (in regime n ≤ 10⁶ cm⁻³) → because unstable disk builds up (Jappsen et al. 2009a)
- Second Seco

transition Pop III to Pop II.5

- star formation will depend on degree of turbulence in protogalactic halo (see talk by Paul Clark)
- speculation: differences in stellar mass function?
- speculation:
 - Iow-mass halos → low level of turbulence → relatively massive stars?



⁽Greif et al. 2008, MNRAS, 387, 1021)

Inigh-mass halos (atomic cooling halos) → high degree of turbulence → wider mass spectrum, peak at lowermasses?



see also talk by Thomas Greif

turbulence developing in an atomic cooling halo

(Greif et al. 2008, MNRAS, 387, 1021, see also Wise & Abel 2007)



turbulence developing in an atomic cooling halo (Greif et al. 2008, MNRAS, 387, 1021)

see also talk by *Thomas Greif*

transition: Pop III to Pop II.5



(Omukai et al. 2005, ApJ, 626, 627, also Schneider et al. 2006, MNRAS, 369, 1437)

dust induced fragmentation at Z=10⁻⁵

t = t_{SF} - 67 yr



 $t = t_{SF} - 20 yr$



 $t = t_{SF}$



 $t = t_{SF} + 53 \text{ yr}$



 $t = t_{SF} + 233 \text{ yr}$



 $t = t_{SF} + 420 \text{ yr}$



(Clark et al. 2008, ApJ, 672, 757)

dust induced fragmentation at Z=10⁻⁵



dense cluster of low-mass protostars builds up:

- mass spectrum peaks below 1 M_{sun}
- cluster VERY dense $n_{stars} = 2.5 \times 10^9 \, pc^3$
- fragmentation at density $n_{gas} = 10^{12} - 10^{13} \text{ cm}^{-3}$

(Clark et al. 2008, ApJ, 672, 757, see also Machida et al. 2009, 399, 1255)

binary fragmentation $(Z=10^{-4})$



dust induced fragmentation at Z=10⁻⁵



dust induced fragmentation at Z=10⁻⁵



(plot from Salvadori et al. 2006, data from Frebel et al. 2005)

dense cluster of low-mass protostars builds up:

 mass spectrum peaks below I M_{sun}
 cluster VERY dense n_{stars} = 2.5 x 10⁹ pc⁻³

- predictions:



⁽Clark et al. 2008, ApJ, 672, 757)

caveats / questions

how good is EOS approach?

fime to reach chemical + thermal equilibrium shorter than dynamical time?

how does EOS depend on dynamics? (e.g. ID collapse with large-gradient approx. versus complex 3D turbulent flows)

how important is heating from stars?

- accretion luminosity may heat gas and reduce degree of cloud fragmentation (cluster formation vs. high-mass SF)
- how can we model that best?
 - full radiation transfer vs. approximate schemes

effects of accretion heating

In how important is heating from stars?

- accretion luminosity may heat gas and reduce degree of cloud fragmentation (cluster formation vs. high-mass SF)
- HOWEVER: the effect is NOT large (see poster by Rowan Smith)



fragmentation of Pop III disk (Z=0) without accretion heating --> fragmentation radius 10.4 AU



fragmentation of Pop III disk (Z=0) with accretion heating --> fragmentation radius 18.9 AU

Evolution of the protostellar disc

No stellar feedback





metall-free star formation



⁽Omukai et al. 2005, ApJ, 626, 627)

metall free star formation

- ♀ first disks are expected to fragment!
 - halos have large angular momentum --> disk forms around first protostar --> roughly isothermal disk are known to be unstable (see talk by Paul Clark) (see also Turk, Abel. & O'Shea 2009, Science, 325,601, Stacy, Greif, & Bromm, 2010, MNRAS, in press, Clark et al. in preparation)
 - also turbulence may lead to fragmentation! (see talk by Paul Clark)
- It do first all first stars form in small clusters? what is the IMF?



- some phenomenology
- processes that influence present-day SF and their possible relevance for high-z SF:
 - turbulence
 - thermodynamics

magnetic fields

🥯 feedback

effects of magnetic fields

magnetic fields can

- suppress disk fragmentation (Ziegler 2005, A&A, 435, 385, Hennebelle & Fromang 2008, A&A, 477, 9, Hennebelle & Teyssier 2008, A&A, 477, 25)
- Grive jets and outflows (Machida et al. 2006, ApJ, 647, LI)
- Induce additional turbulence (MRI / dynamo) (Balbus & Hawley 1998, RMP, 70, I, Brandenburg & Subramanian 2005, Phy. Rep. 417, I)
- maybe present even in Z=0 gas (either as primordial fields or generated by dynamo action) see talk by Dominik Schleicher on Thursday
- ${}^{\scriptsize extsf{@}}$ need to be taken into account ${}^{\scriptsize extsf{!!}}$

effects of magnetic fields



weak field: binary formation

sufficient field strength: single star

(protostellar collapse with m=2 perturbation and B-field: Ziegler 2005, A&A, 435, 385, for further discussion, see Hennebelle & Teyssier 2008, A&A, 477, 25



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feedback (focus on ionizing feedback)

single star: HII region and outflow



disk edge on

disk plane

(Peters et al. 2010, ApJ, 711, 1017)

multiple protostars: dynamics of HII regions



disk edge on

disk plane

(Peters et al. 2010, ApJ, 711, 1017)

accretion history



- ionizing radiation cannot stop accretion
- however, fragmentation of disk can stop mass growth of the central star

--> fragmentation induced starvation (Peters et al. 2010, ApJ, 711, 1017)

feedback in Pop III star formation

- we expect feedback during PopIII protostellar collapse
 - NOT to stop fragmentation (but possibly reduce number of fragments)
 - NOT to prevent mass growth (if at all *fragmentation induced starvation* stops further mass growth)
- we expect the effects of magnetic fields to be potentially more important.

summary

- - turbulence
 - thermodynamics
 - magnetic fields
 - feedback

to influence Pop III/II star formation.



- masses of Pop III stars still uncertain (expect surprises from new generation of high-resolution calculations that go beyond first collapse)
- disks unstable: Pop III stars should be binaries or part of small cluster
- effects of feedback less important than in present-day SF