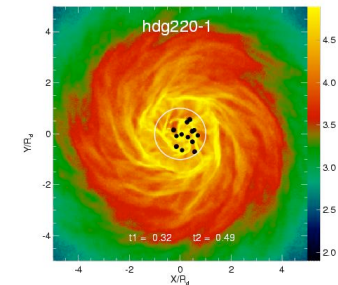
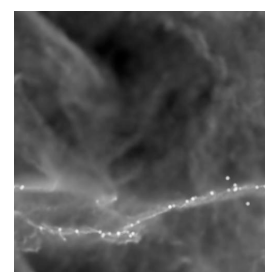
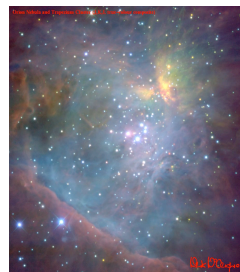
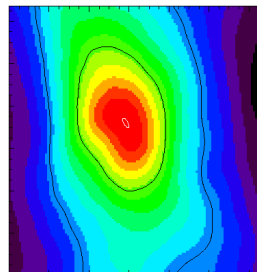
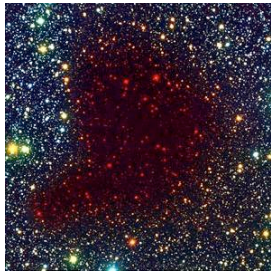




# ISM Turbulence



**Ralf Klessen**

Zentrum für Astronomie der Universität Heidelberg  
Institut für Theoretische Astrophysik

Deutsche  
Forschungsgemeinschaft  
**DFG**





# Agenda

- some phenomenology

- turbulence in the ISM

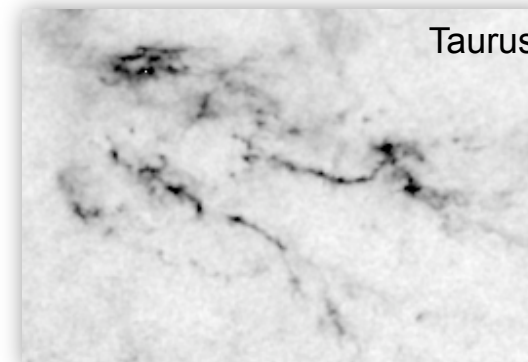
- a simple cartoon picture

- relation between ISM turbulence and star formation

- some basic thoughts

- accretion driven turbulence
- the stellar mass function
- formation first (strong) magnetic fields in the universe
- mixing in supersonic turbulence

- summary

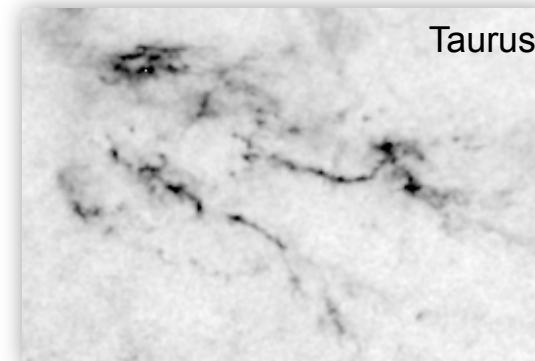




# Interstellar Matter: ISM

Abundances, scaled to 1.000.000 H atoms

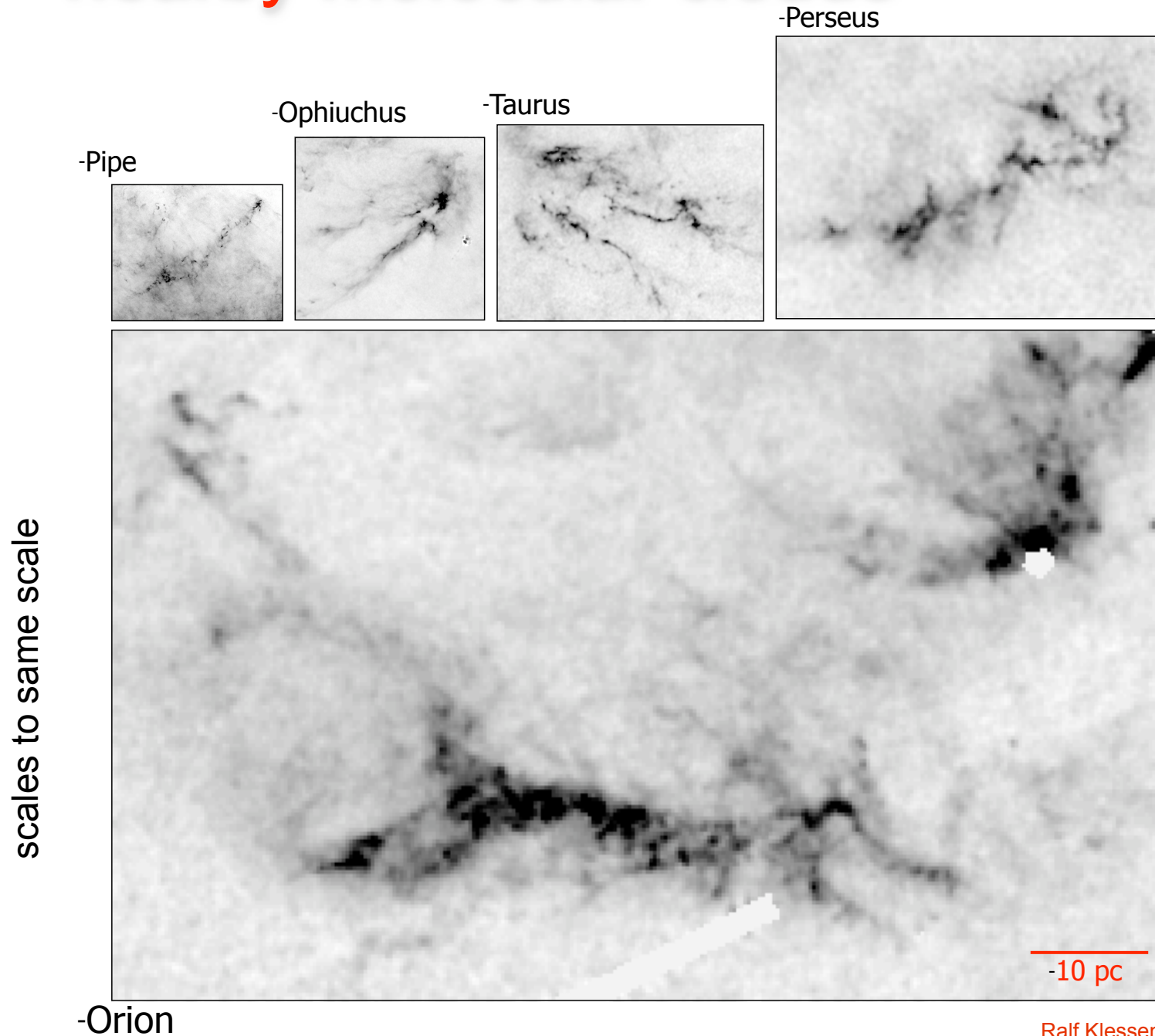
element	atomic number	abundance
hydrogen	H 1	1.000.000
deuterium	${}_1\text{H}^2$ 1	16
helium	He 2	68.000
carbon	C 6	420
nitrogen	N 7	90
oxygen	O 8	700
neon	Ne 10	100
sodium	Na 11	2
magnesium	Mg 12	40
aluminium	Al 13	3
silicium	Si 14	38
sulfur	S 16	20
calcium	Ca 20	2
iron	Fe 26	34
nickel	Ni 28	2



hydrogen is by far the most abundant element (more than 90% in number).



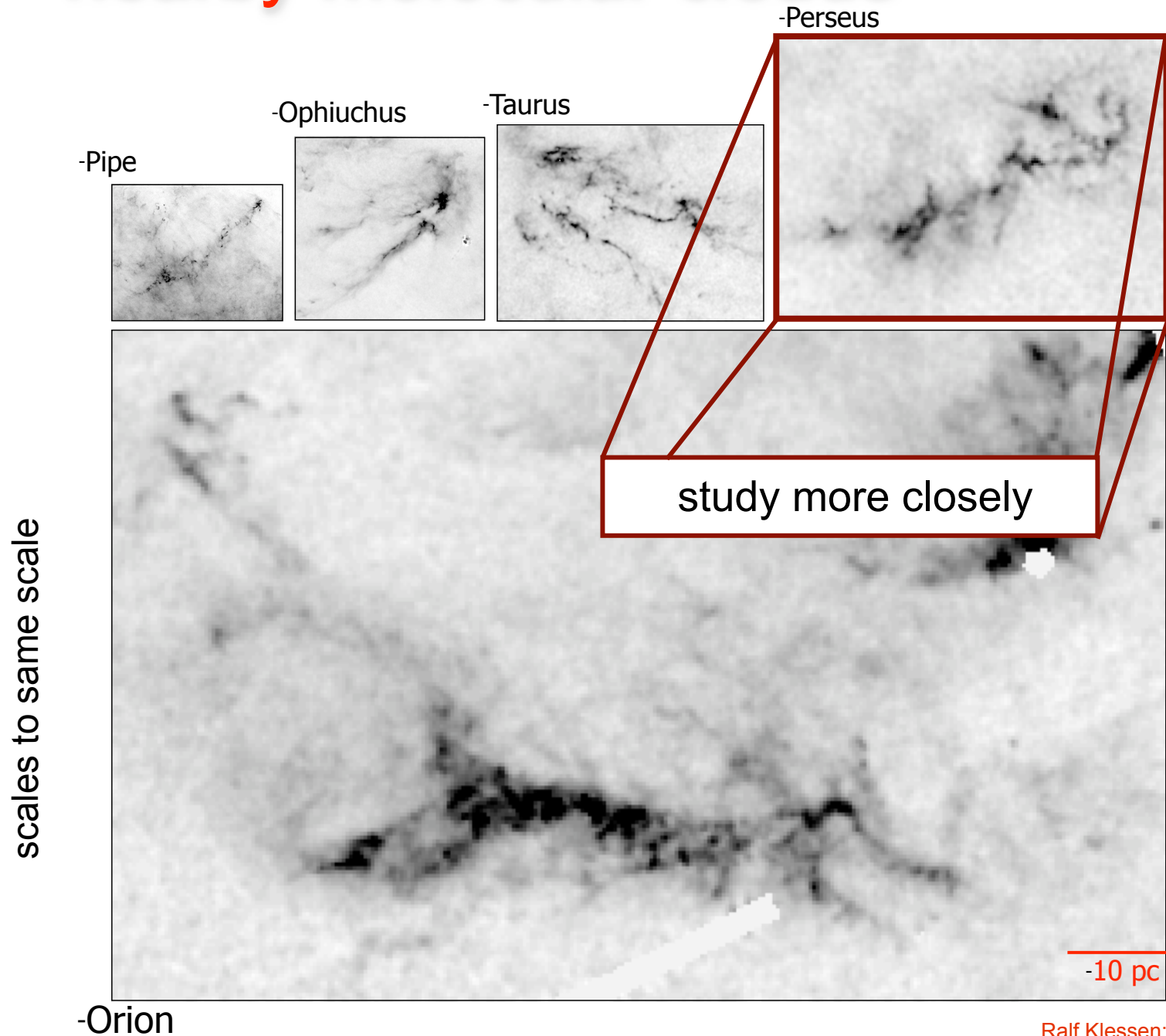
# nearby molecular clouds



(from A. Goodman)



# nearby molecular clouds



(from A. Goodman)

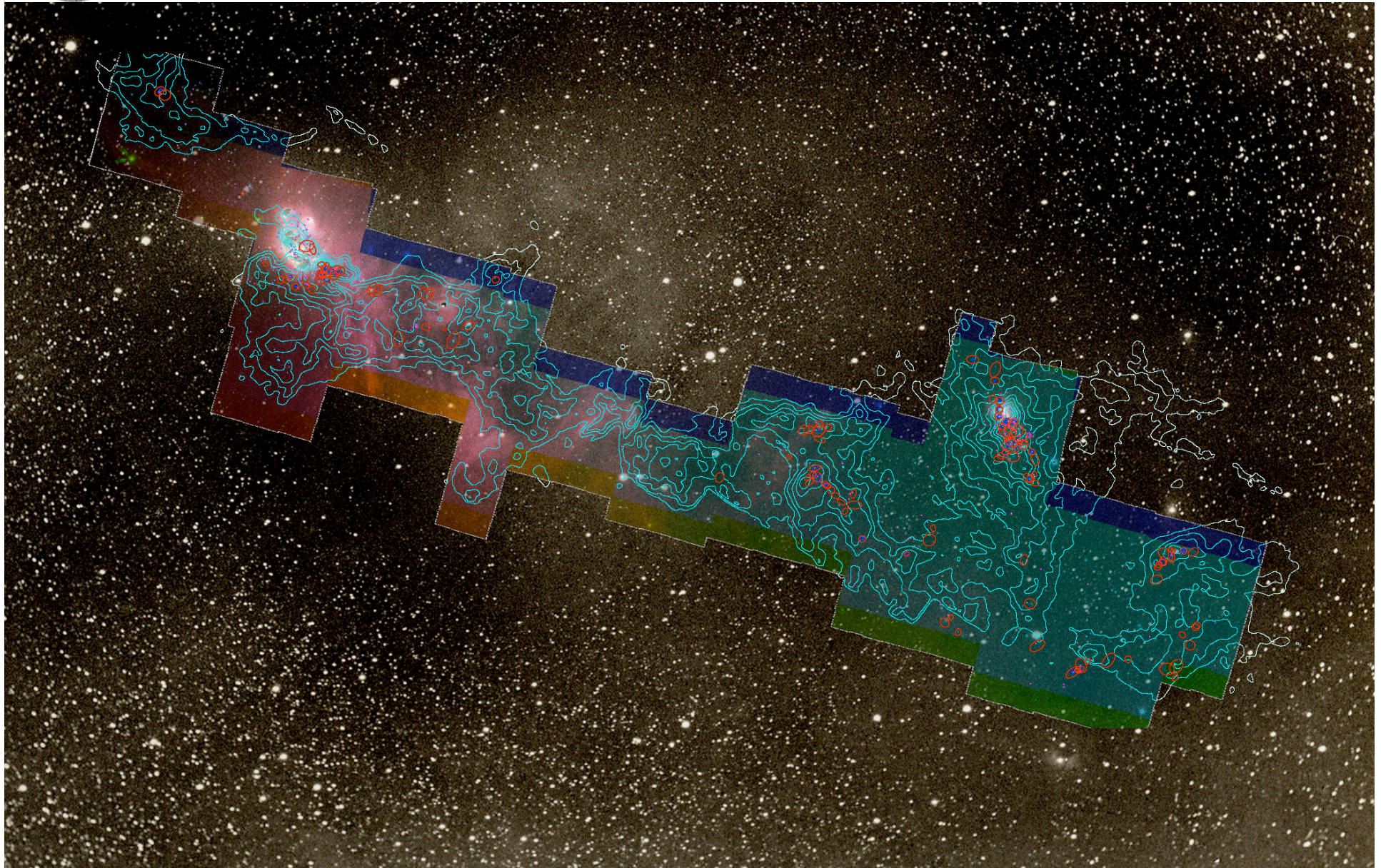
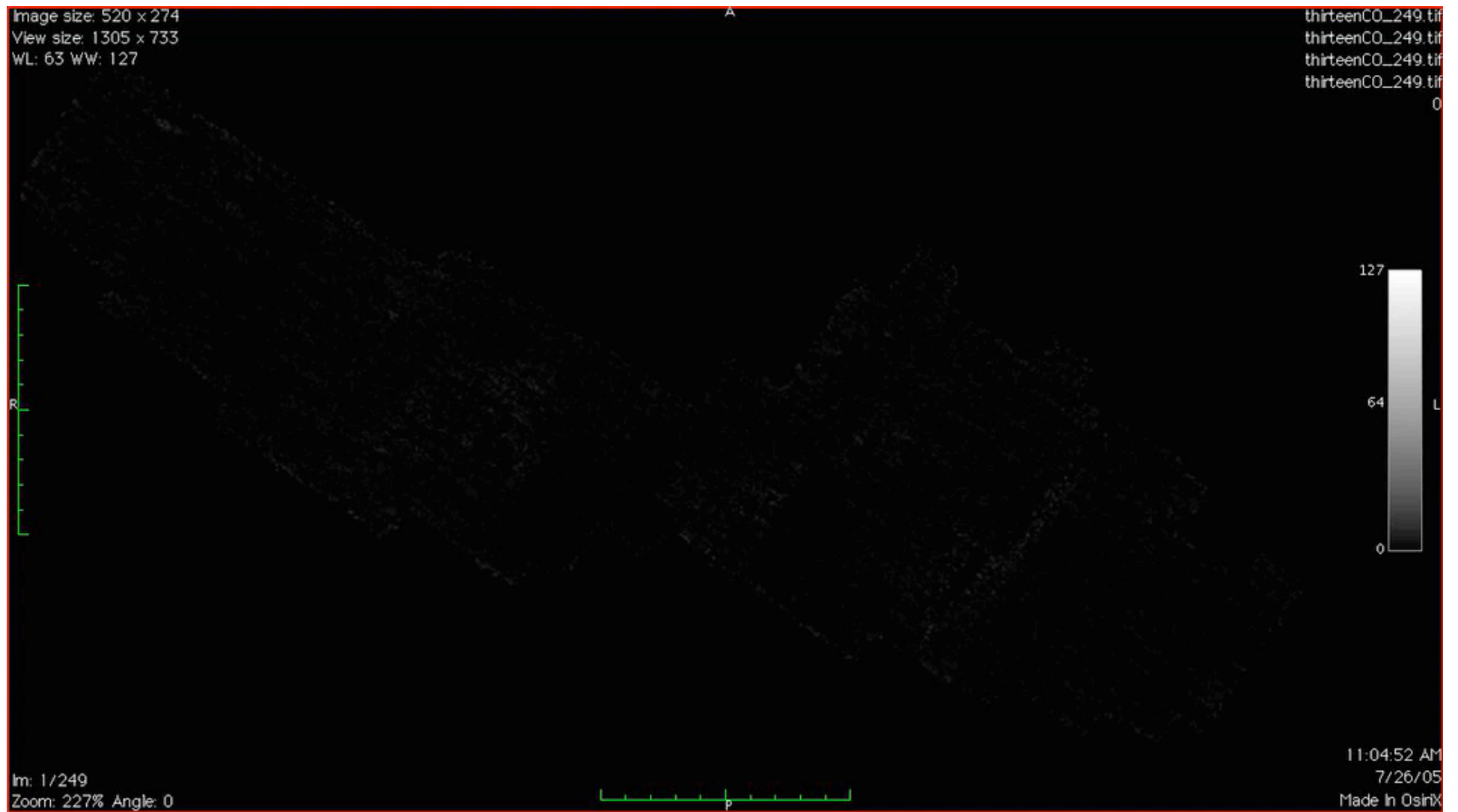


image from Alyssa Goodman: COMPLETE survey

Ralf Klessen: Santa Cruz, July 9.2010



velocity distribution in Perseus

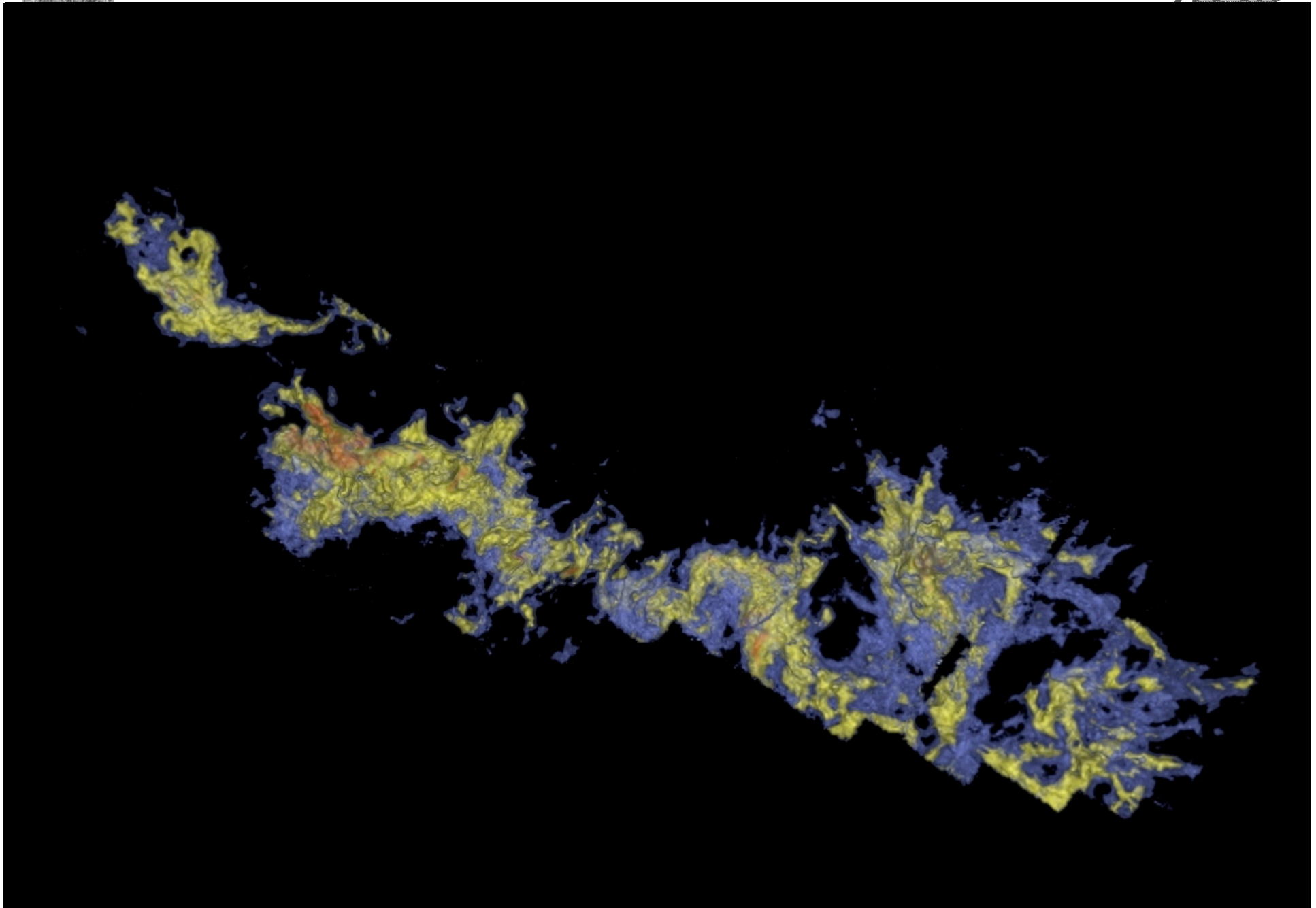
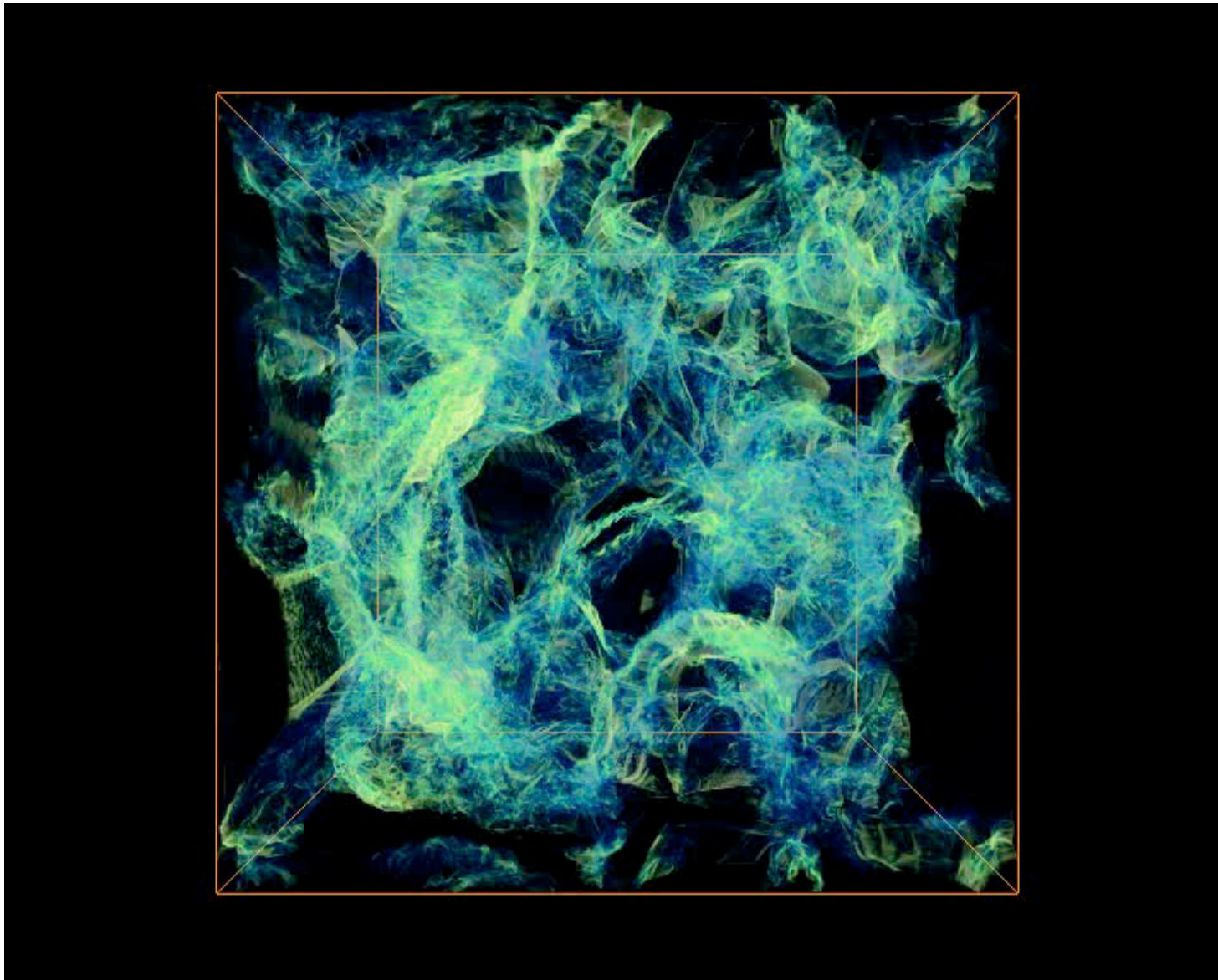


image from Alyssa Goodman: COMPLETE survey



(movie from Christoph Federrath)

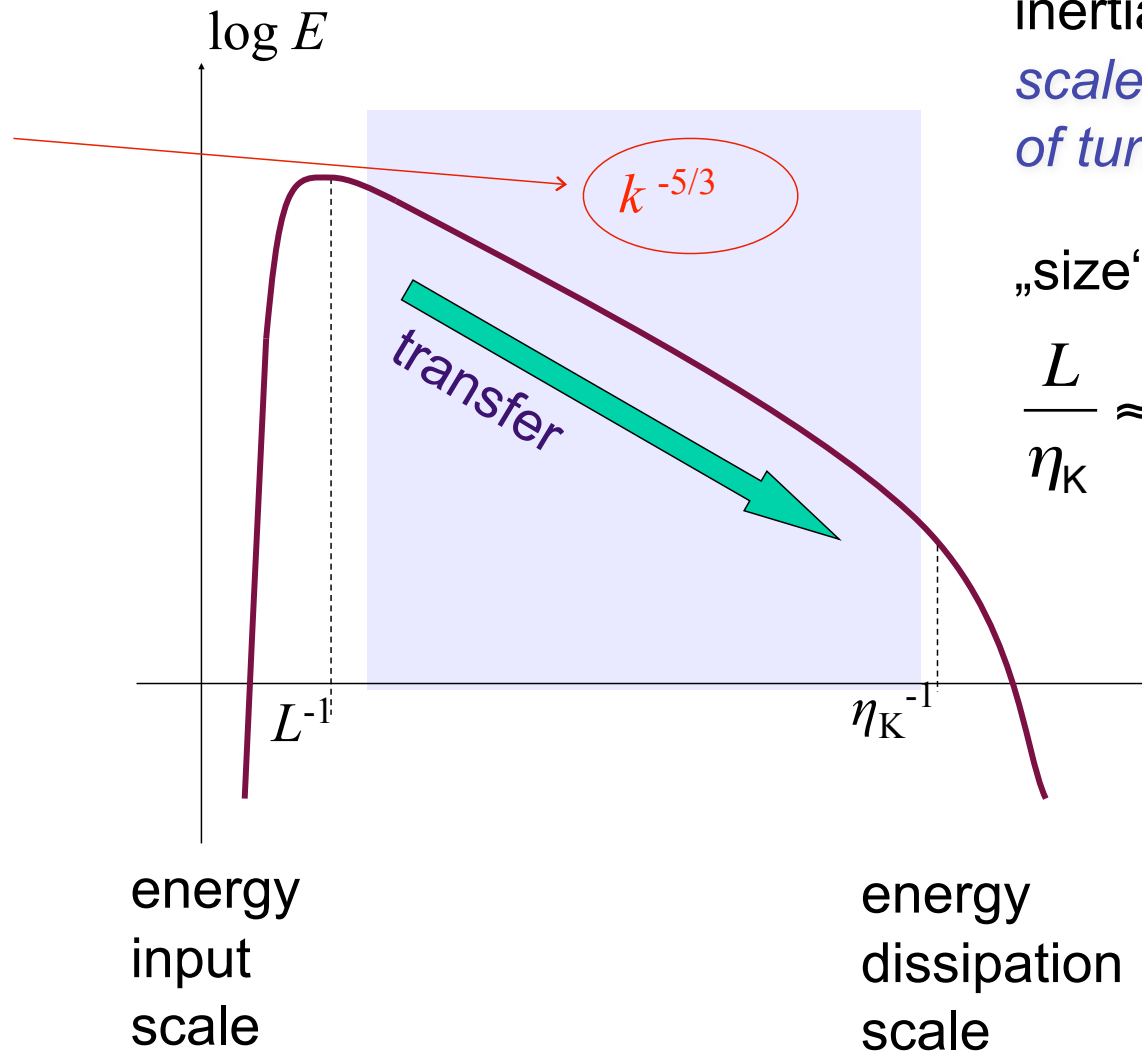


turbulence



# Turbulent cascade

Kolmogorov (1941) theory  
incompressible turbulence



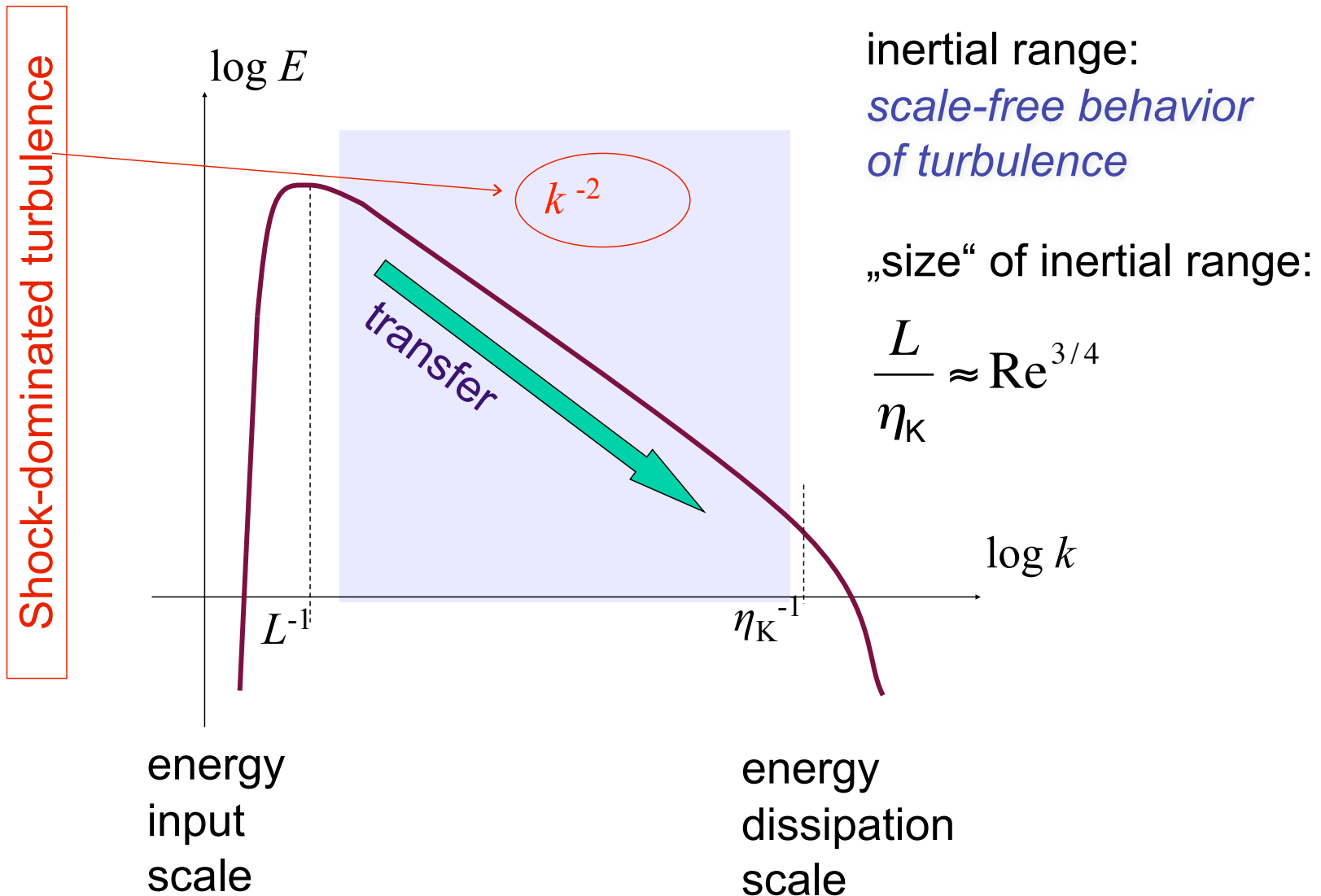
inertial range:  
*scale-free behavior  
of turbulence*

„size“ of inertial range:

$$\frac{L}{\eta_K} \approx \text{Re}^{3/4}$$

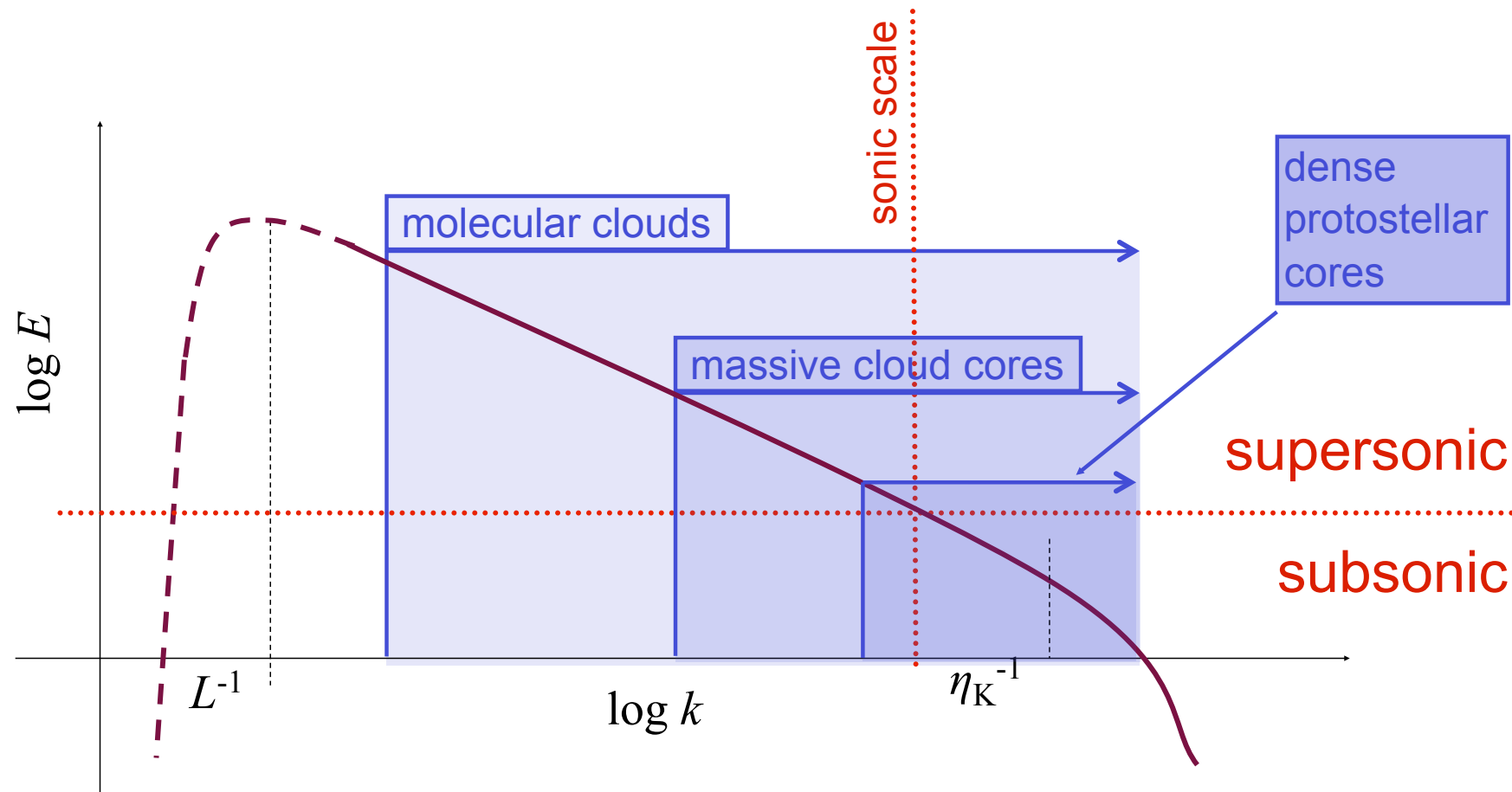


# Turbulent cascade





# Turbulent cascade in ISM



energy source & scale  
*NOT known*  
(supernovae, winds,  
spiral density waves?)

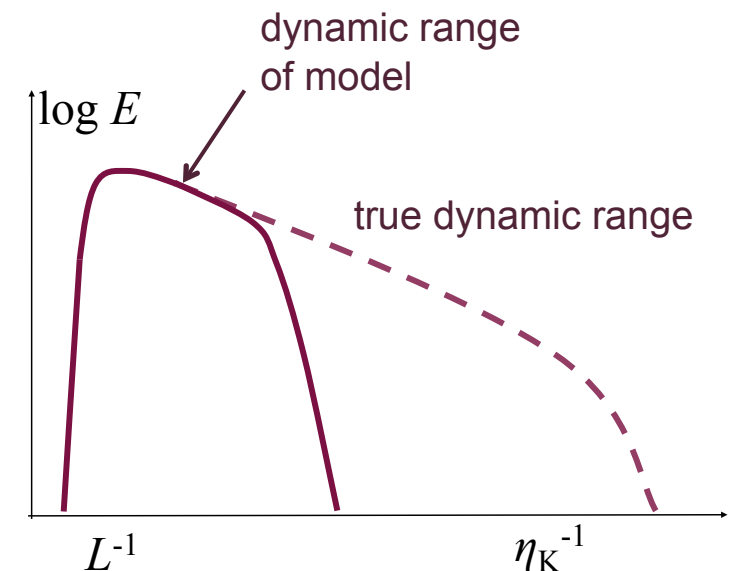
$$\sigma_{\text{rms}} \ll 1 \text{ km/s}$$
$$M_{\text{rms}} \leq 1$$
$$L \approx 0.1 \text{ pc}$$

dissipation scale not known  
(ambipolar diffusion,  
molecular diffusion?)

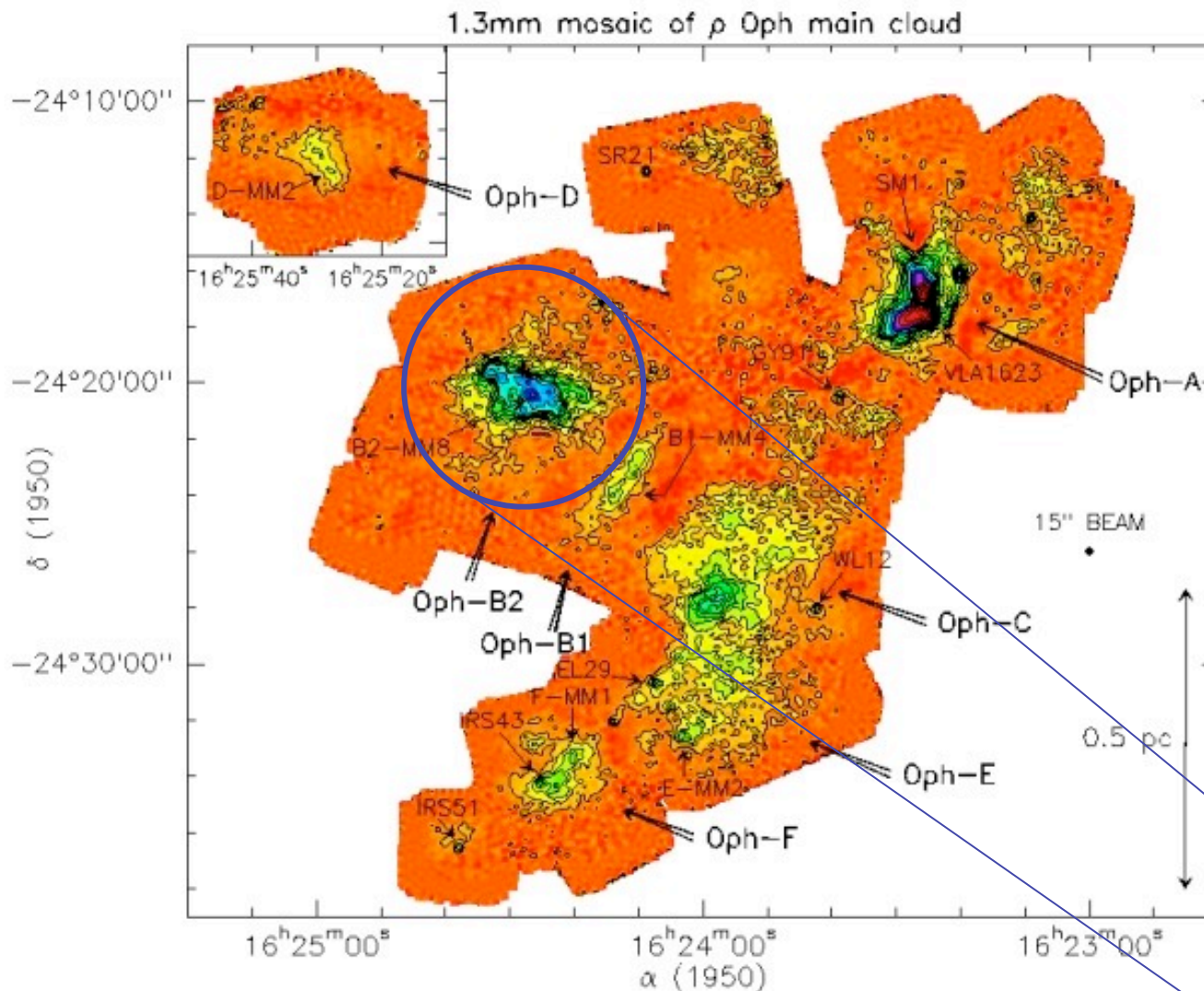


# Large-eddy simulations

- We use **LES** to model the large-scale dynamics
- Principal problem: only large scale flow properties
  - Reynolds number:  $Re = LV/\nu$  ( $Re_{nature} \gg Re_{model}$ )
  - dynamic range much smaller than true physical one
  - need **subgrid model** (in our case simple: only dissipation)
  - but what to do for more complex when processes on subgrid scale determine large-scale dynamics (chemical reactions, nuclear burning, etc)
- Turbulence is “space filling” --> difficulty for AMR (don't know what criterion to use for refinement)
- How **large** a Reynolds number do we need to catch basic dynamics right?



# Density structure of MC's



molecular clouds  
are highly  
inhomogeneous

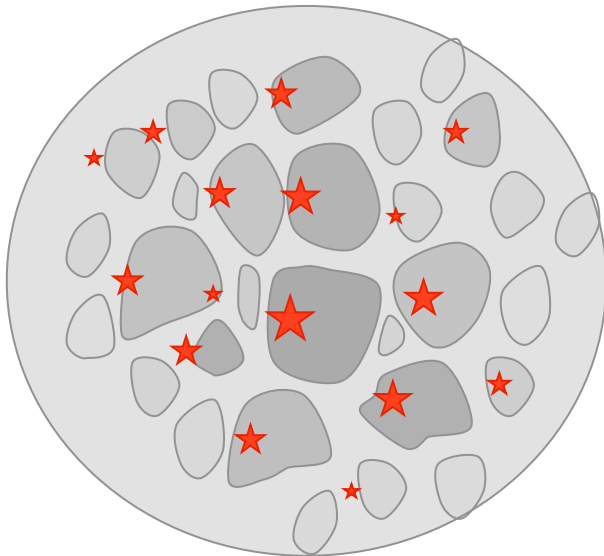
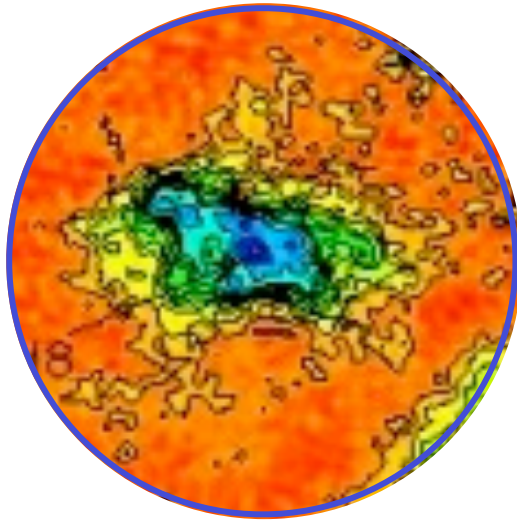
stars form in the  
densest and coldest  
parts of the cloud

$\rho$ -Ophiuchus cloud  
seen in dust  
emission

let's focus on  
a cloud core  
like this one

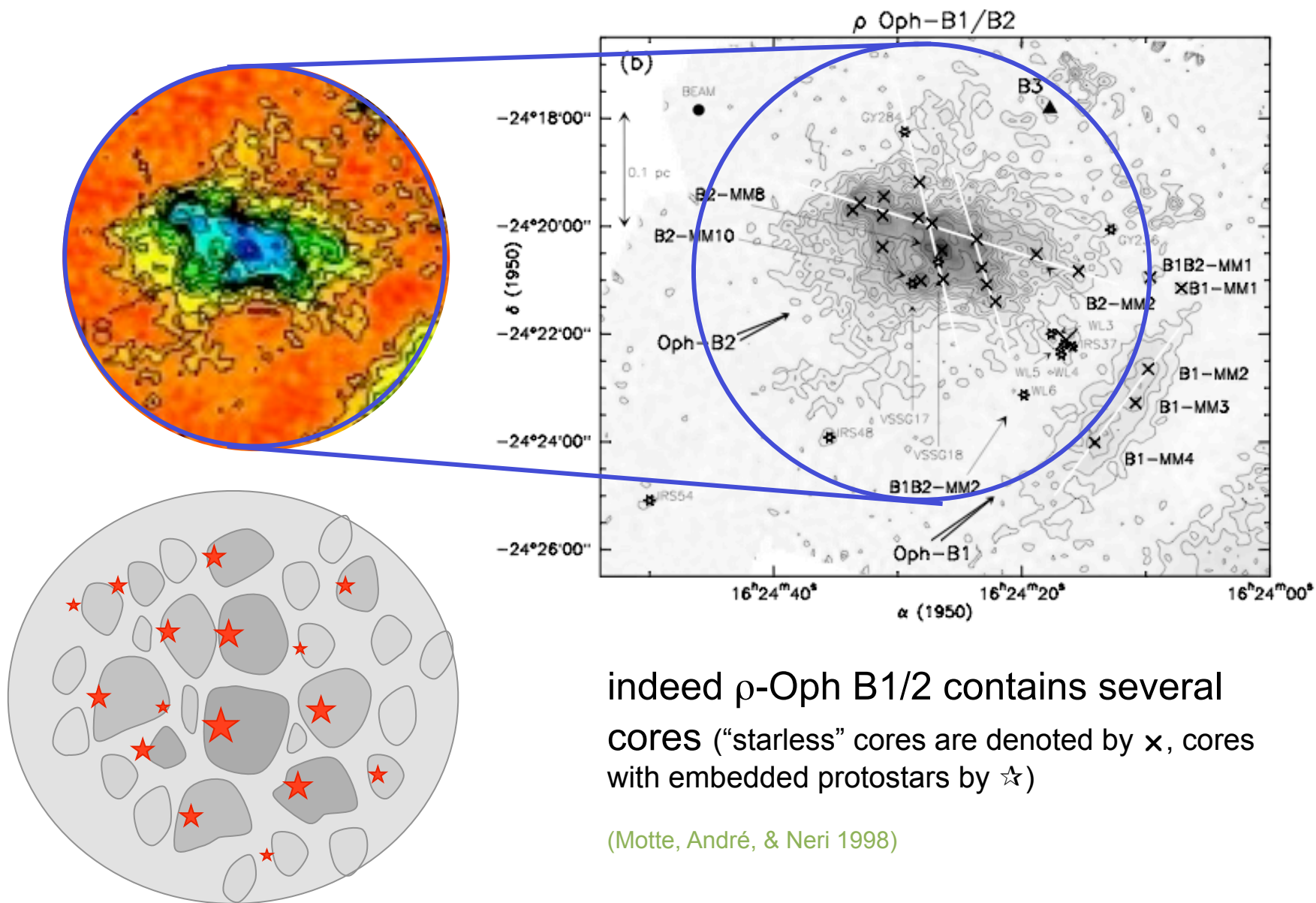
(Motte, André, & Neri 1998)

# Evolution of cloud cores



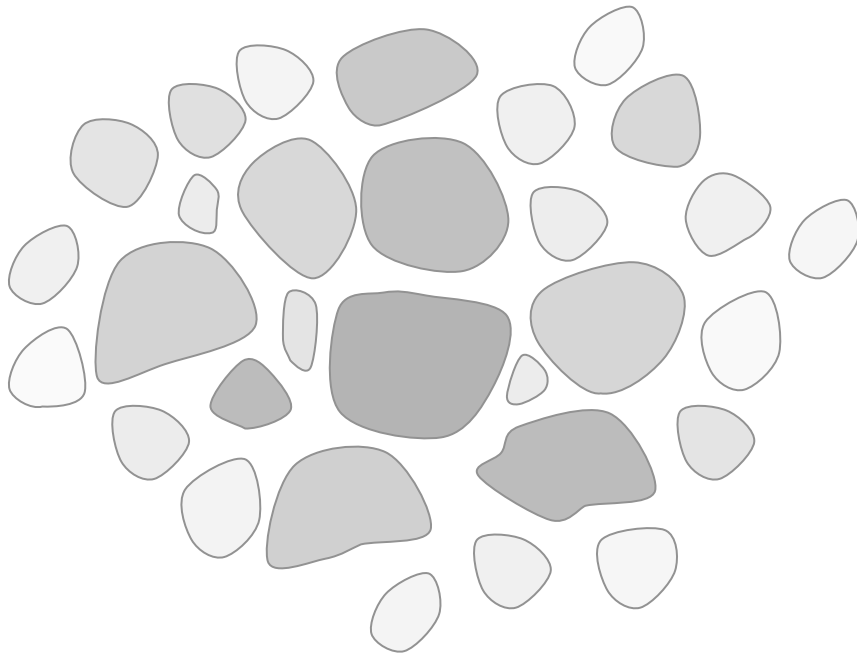
- How does this core evolve?  
Does it form one single massive star or cluster with mass distribution?
- Turbulent cascade „goes through“ cloud core  
--> NO *scale separation* possible  
--> NO *effective sound speed*
- Turbulence is supersonic!  
--> produces strong density contrasts:  
$$\delta\rho/\rho \approx M^2$$
  
--> with typical  $M \approx 10$  -->  $\delta\rho/\rho \approx 100!$
- many of the shock-generated fluctuations are Jeans unstable and go into collapse
- --> expectation: *core breaks up and forms a cluster of stars*

# Evolution of cloud cores



# Formation and evolution of cores

What happens to distribution of cloud cores?



Two extreme cases:

(1) turbulence dominates energy budget:

$$\alpha = E_{\text{kin}} / |E_{\text{pot}}| > 1$$

--> individual cores do *not* interact

--> *collapse of individual cores*  
dominates *stellar mass growth*

--> *loose cluster of low-mass stars*

(2) turbulence decays, i.e. gravity dominates:

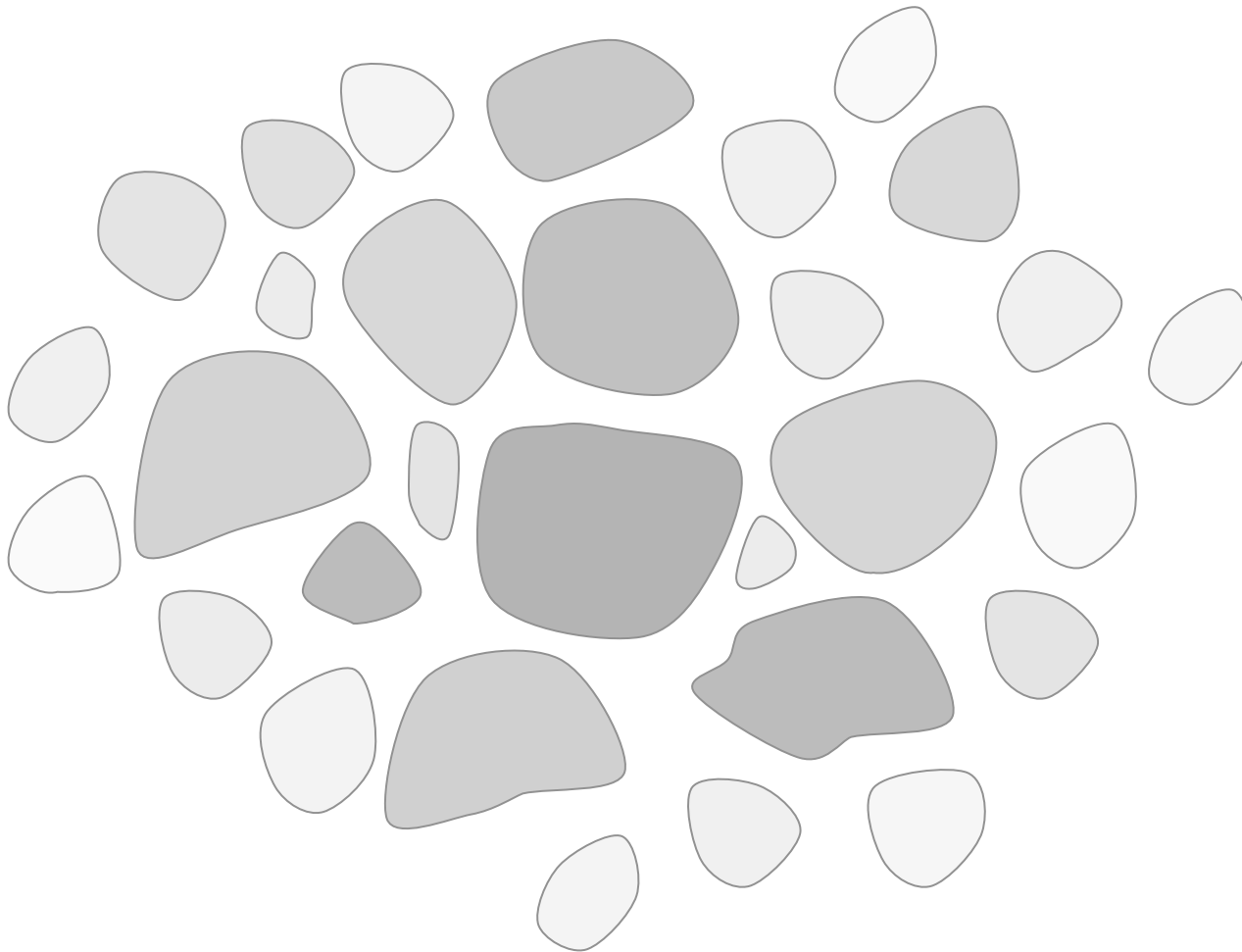
$$\alpha = E_{\text{kin}} / |E_{\text{pot}}| < 1$$

--> *global contraction*

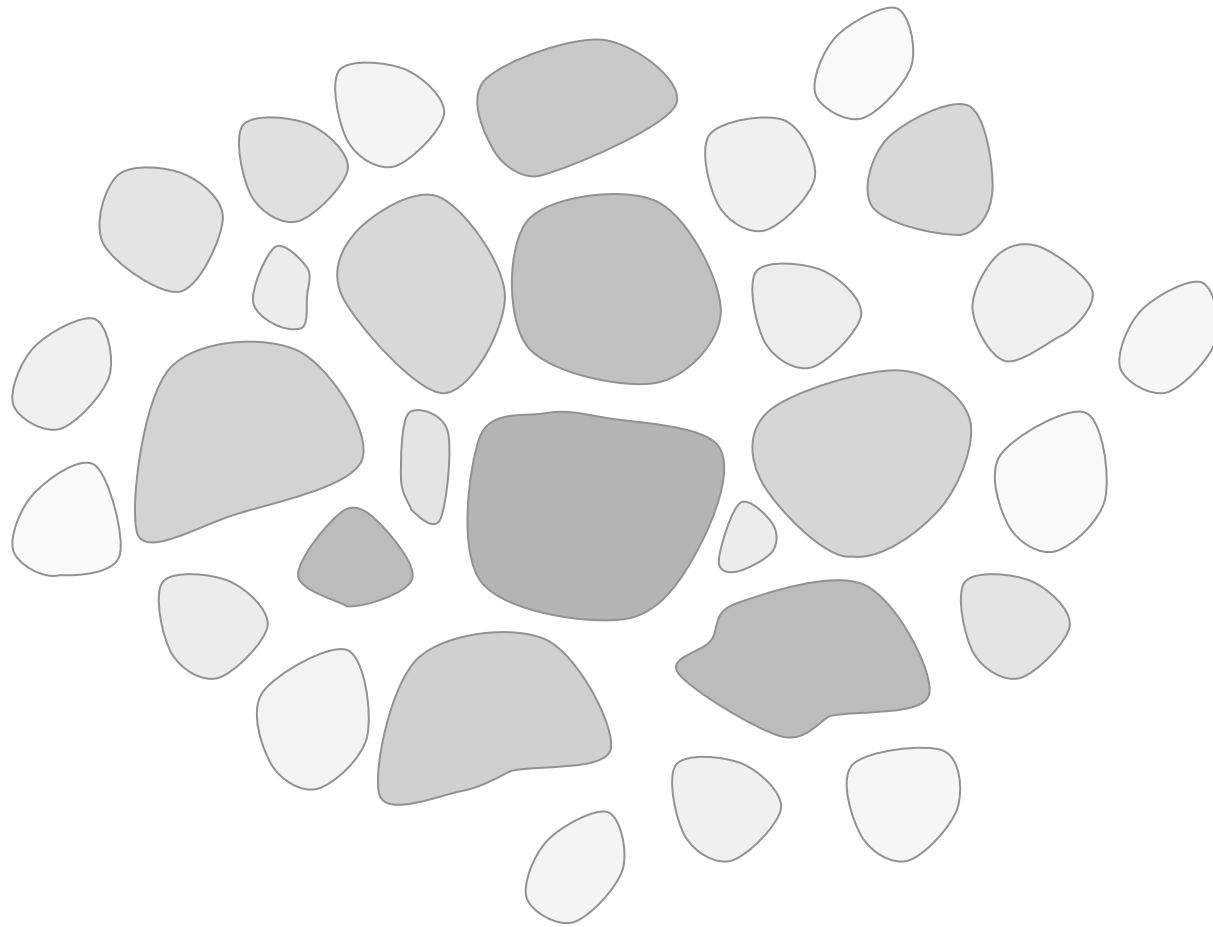
--> core do *interact* while collapsing

--> *competition* influences *mass growth*

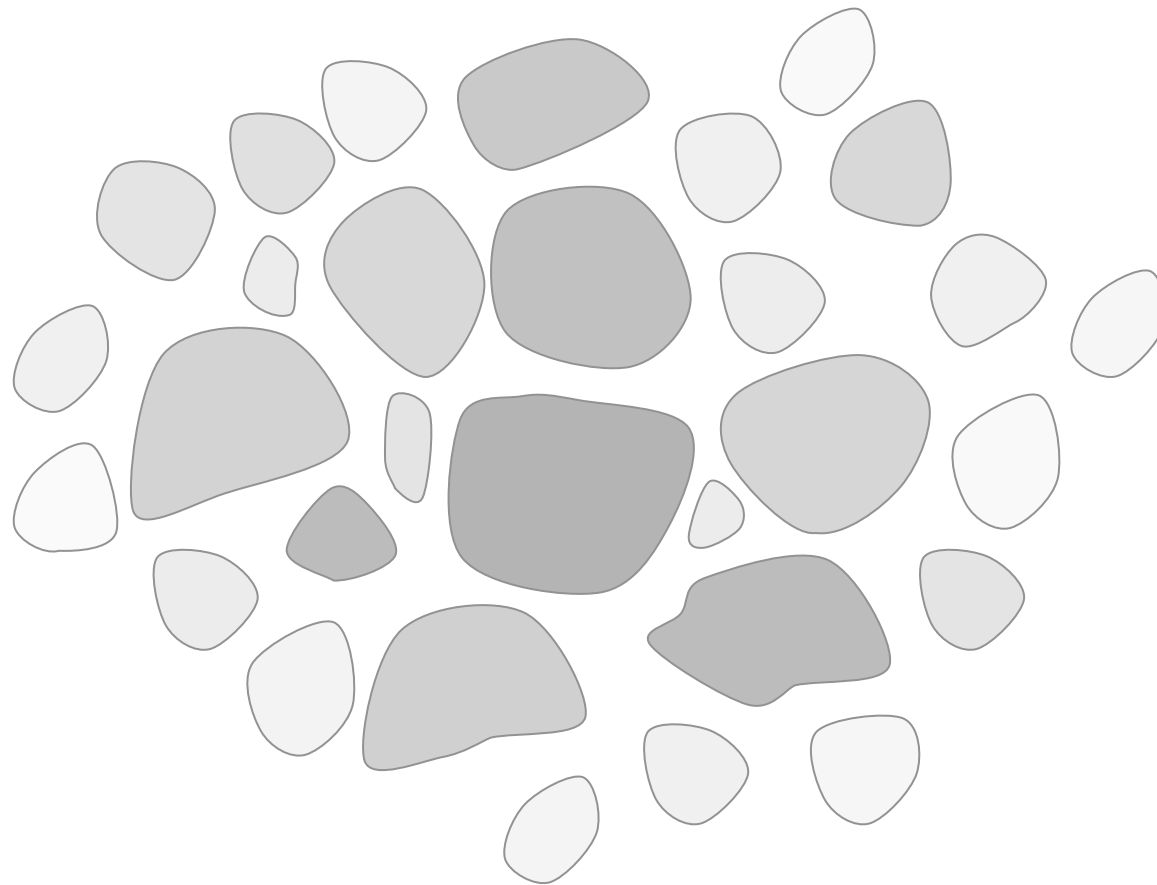
--> *dense cluster with high-mass stars*



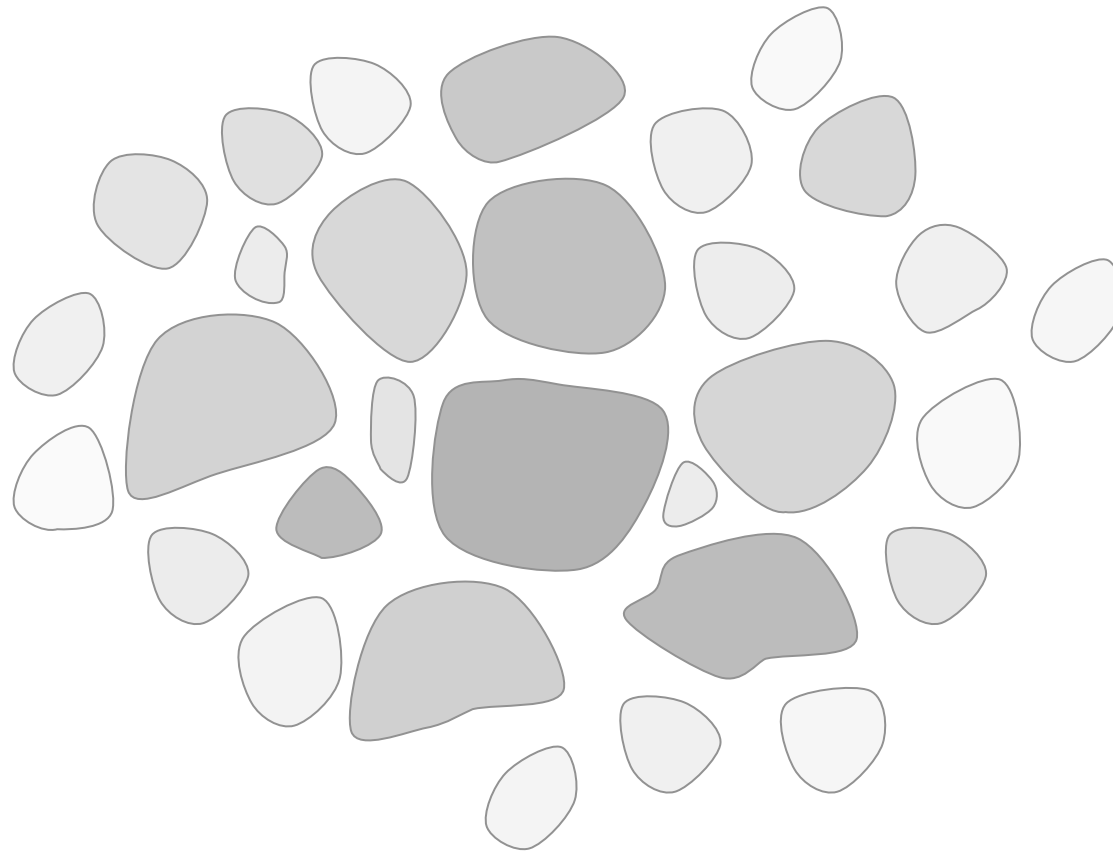
turbulence creates a hierarchy of clumps



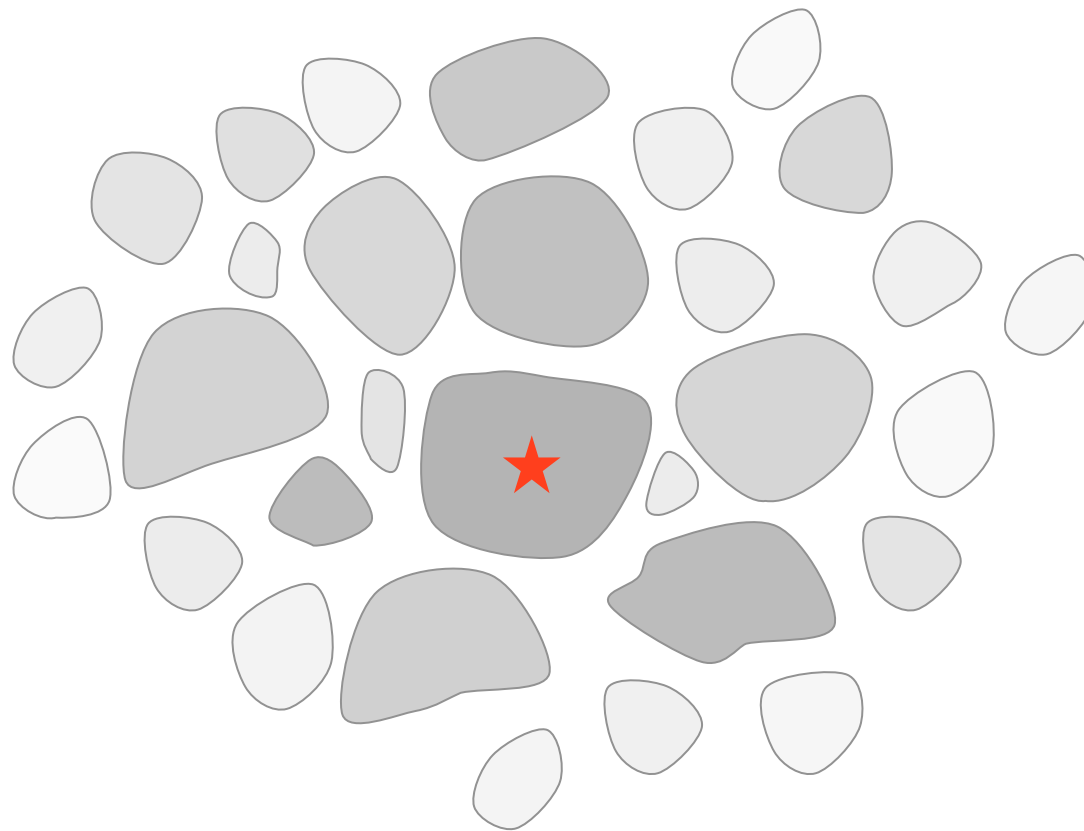
as turbulence decays locally, contraction sets in



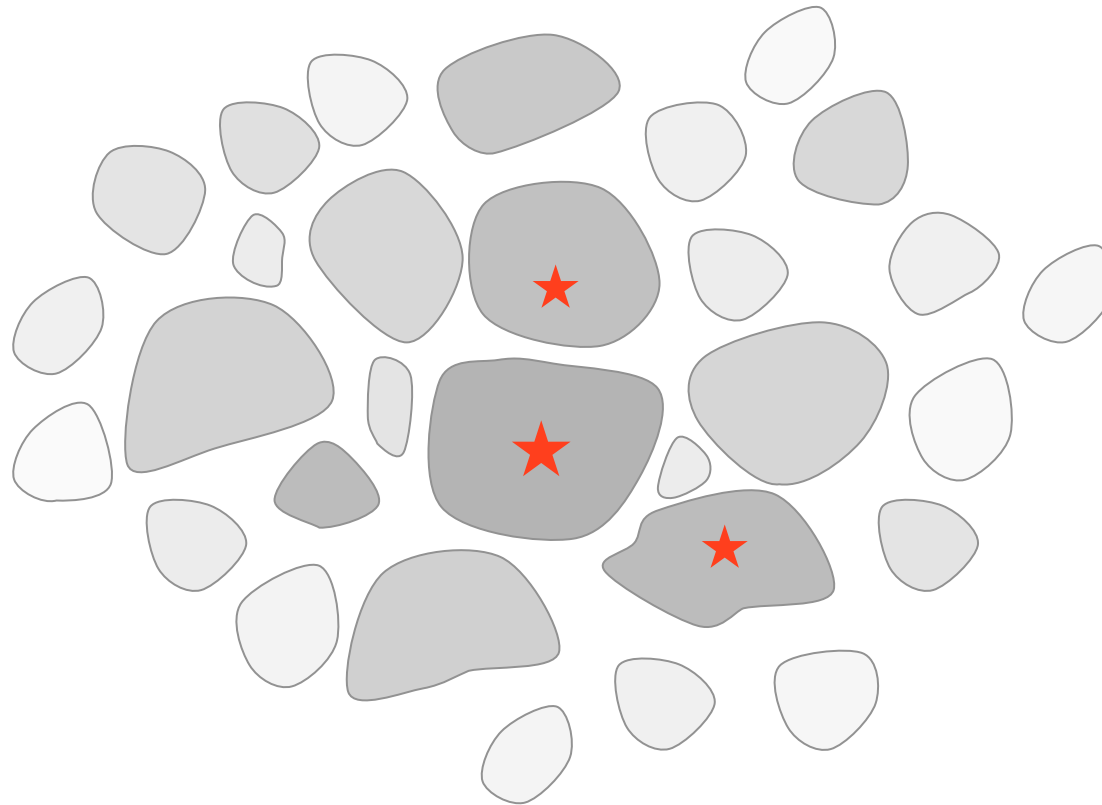
as turbulence decays locally, contraction sets in



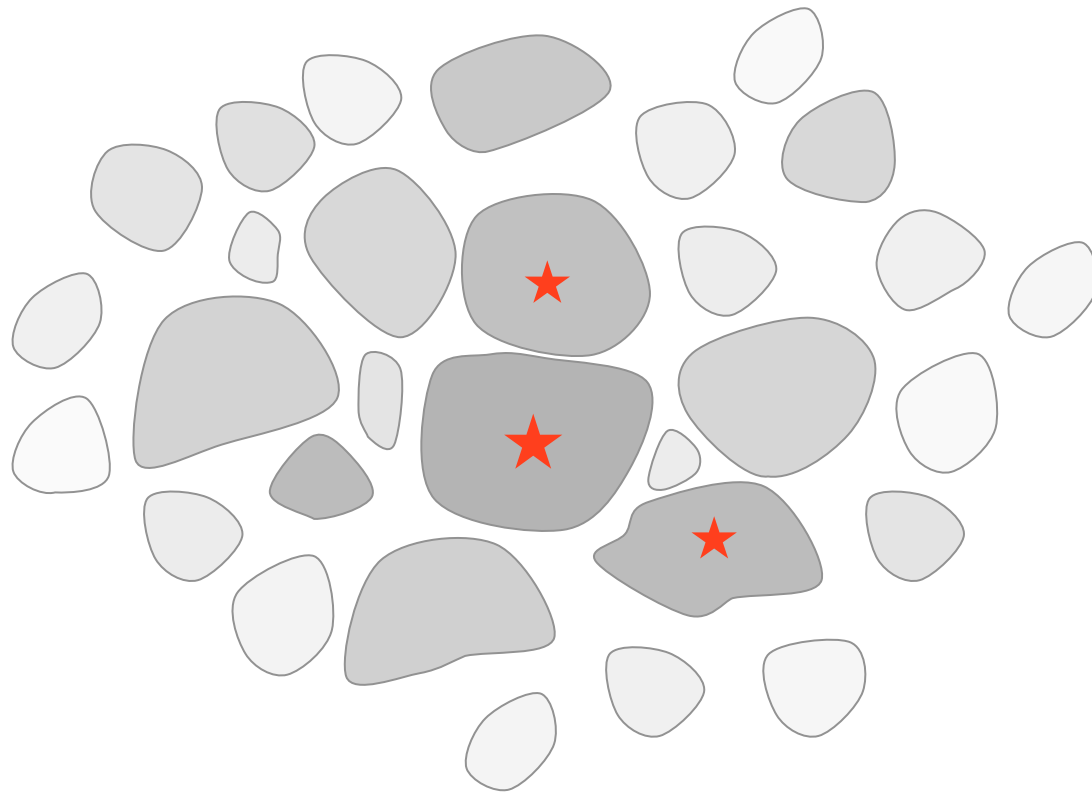
while region contracts, individual clumps collapse to form stars



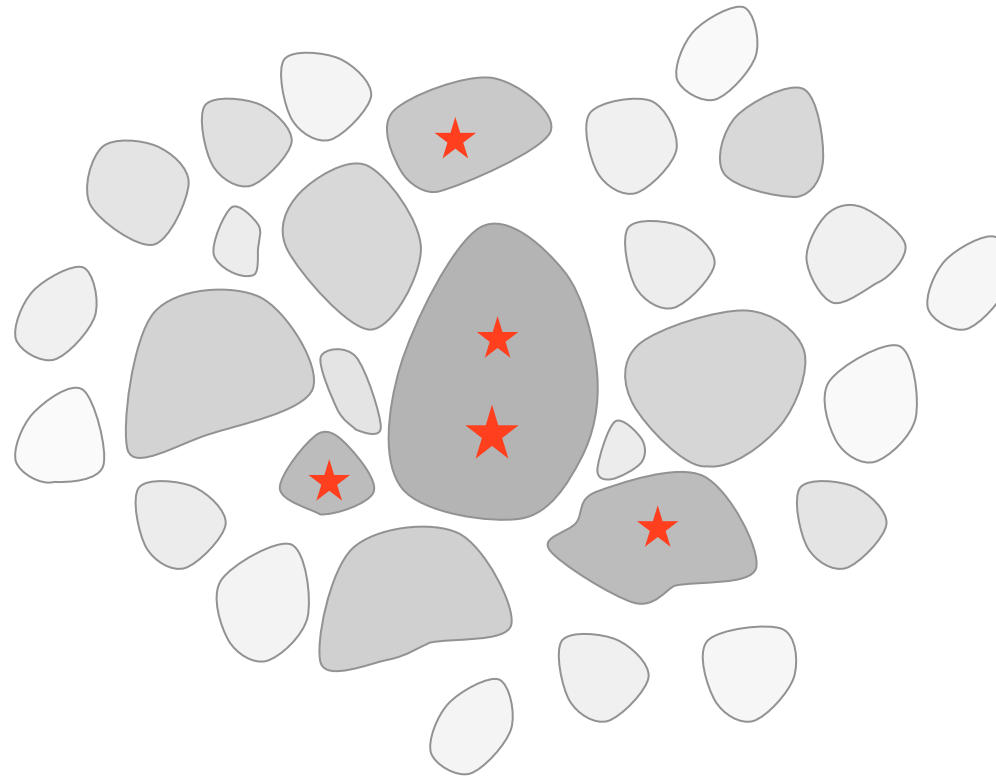
while region contracts, individual clumps collapse to form stars



individual clumps collapse to form stars

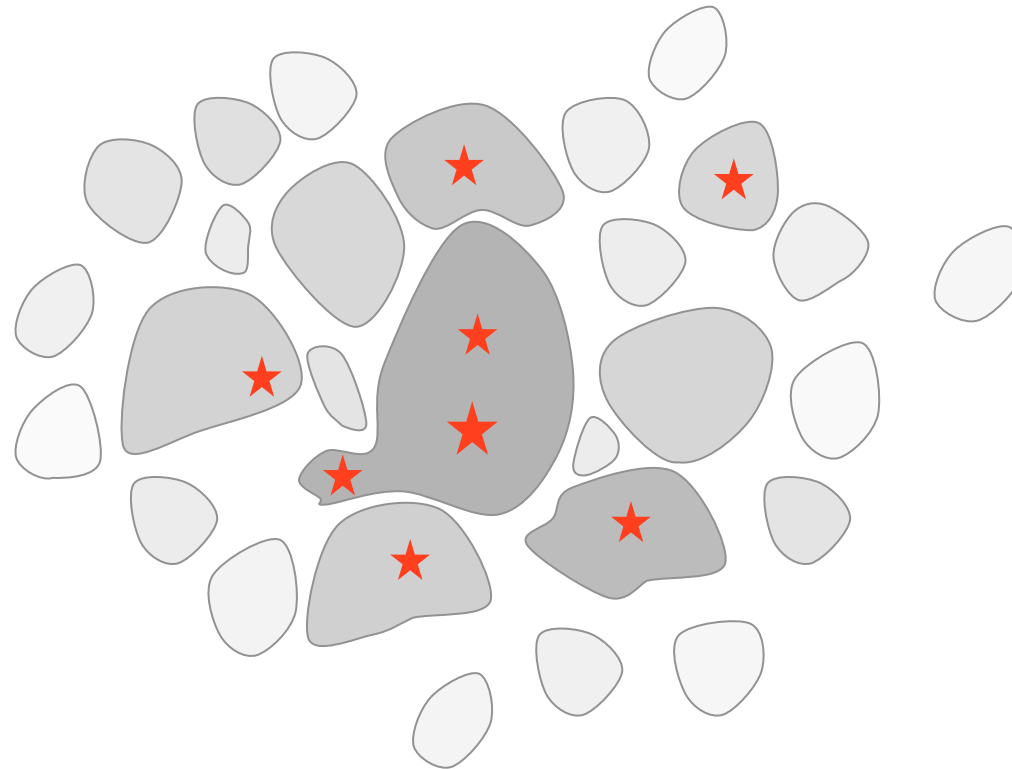


individual clumps collapse to form stars

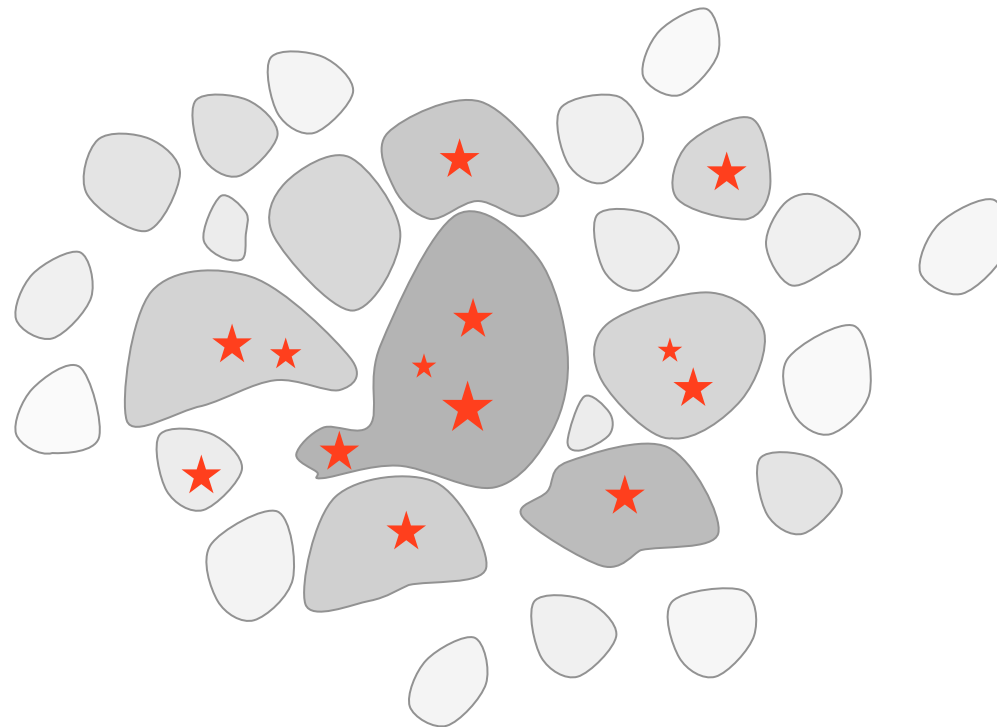


$$\alpha = E_{\text{kin}} / |E_{\text{pot}}| < 1$$

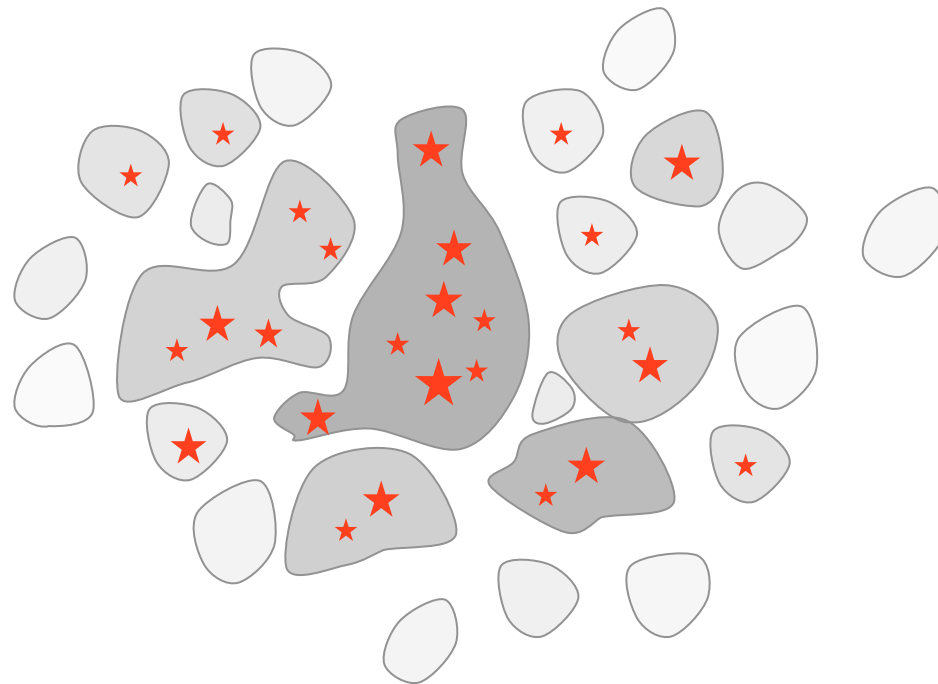
in *dense clusters*, clumps may merge while collapsing  
--> then contain multiple protostars



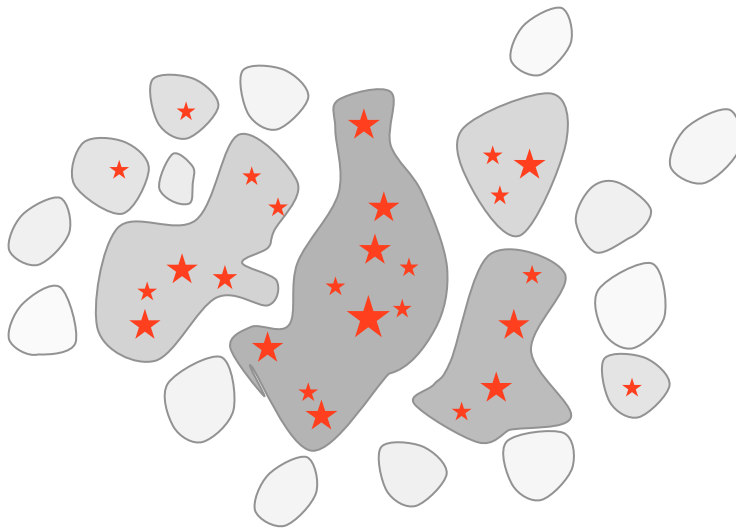
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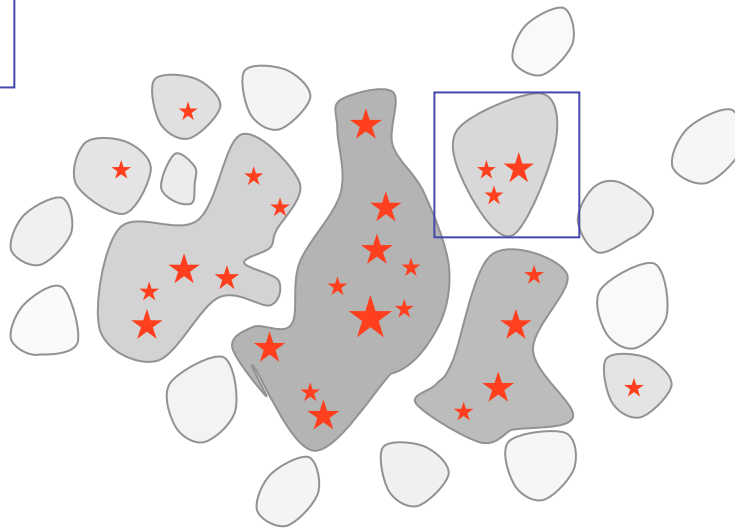
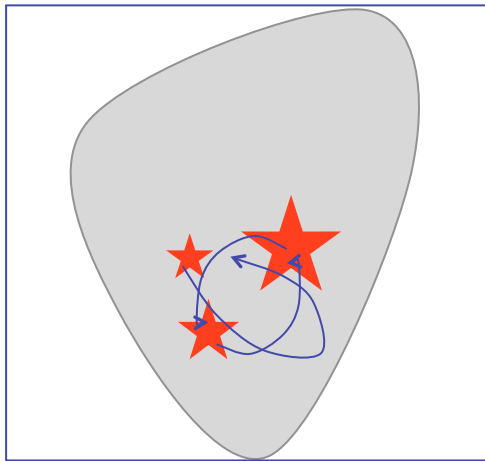
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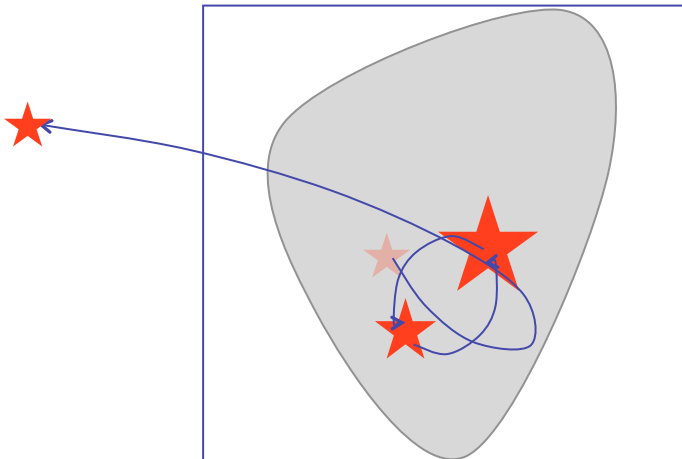
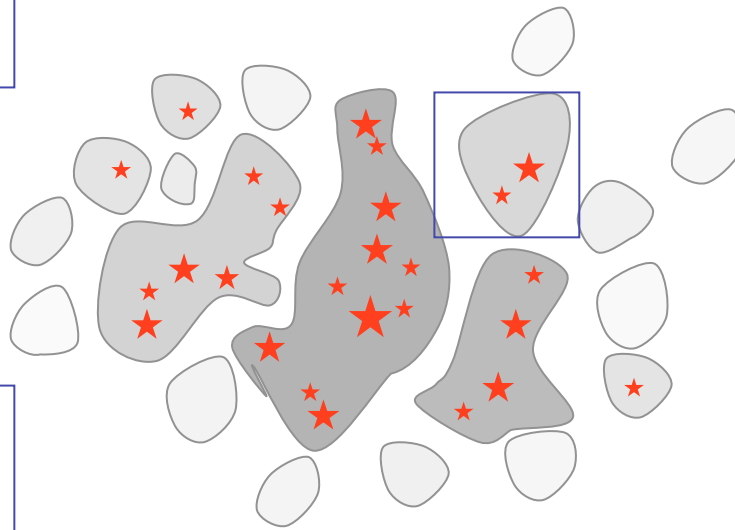
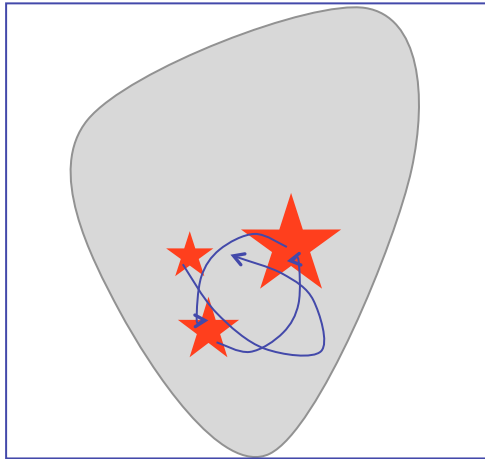
in *dense clusters*, competitive mass growth  
becomes important



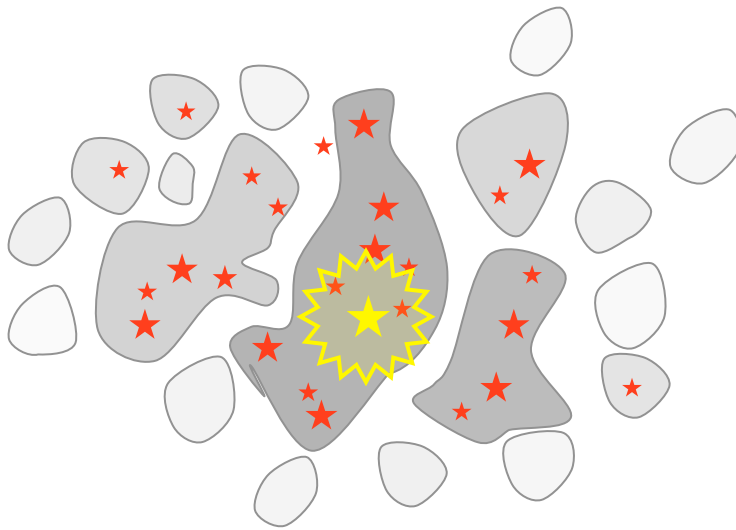
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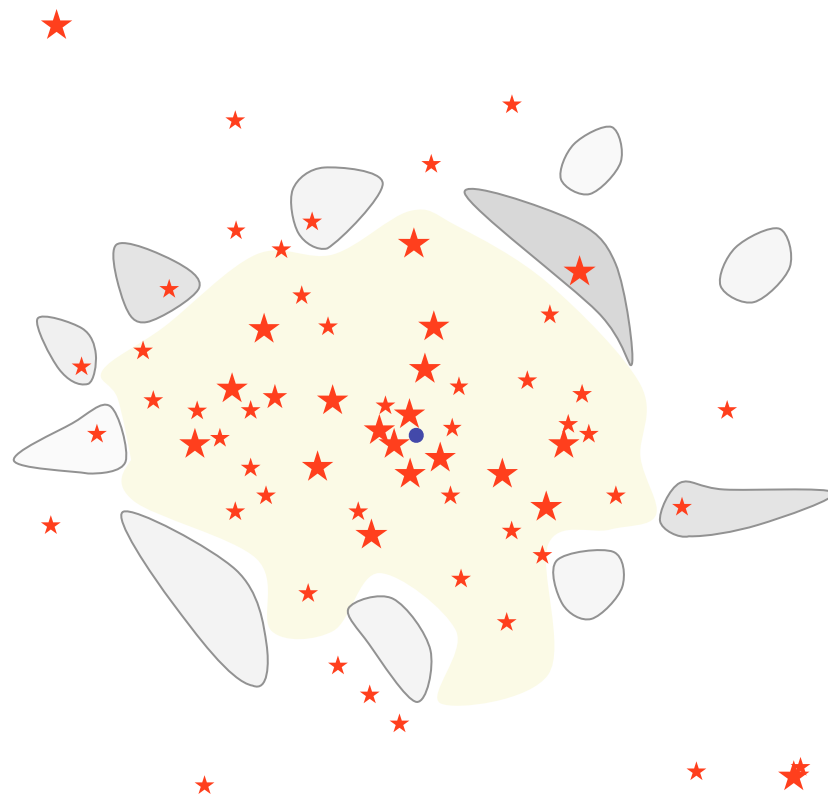
in *dense clusters*,  $N$ -body effects influence mass growth



low-mass objects may  
become ejected --> accretion stops



feedback terminates star formation



result: *star cluster*, possibly with HII region



NGC 602 in the LMC: Hubble Heritage Image

result: *star cluster* with HII region

predictions



## some details

- *what drives turbulence?*
  - accretion driven turbulence on ALL scales  
galaxies, molecular clouds, protostellar disks
- *stellar initial mass function (IMF)*
  - focus on effects of thermodynamics
- *formation of molecular cloud in the turbulent ISM*
  - combining MHD with time-dependent chemistry
- *the first (strong) B-fields in the universe*
  - the turbulent dynamo in action
- *diffusion in supersonic turbulence*
  - mixing length approach

# accretion driven turbulence

## ● thesis:

- astrophysical objects *form* by *accretion* of ambient material
- the *kinetic energy* associated with this process is a key agent *driving internal turbulence*.
- this works on *ALL* scales:
  - galaxies
  - molecular clouds
  - protostellar accretion disks

# concept

- turbulence decays on a crossing time

$$\tau_d \approx \frac{L_d}{\sigma}$$

- energy decay rate

$$\dot{E}_{\text{decay}} \approx \frac{E}{\tau_d} = -\frac{1}{2} \frac{M \sigma^3}{L_d}$$

- kinetic energy of infalling material

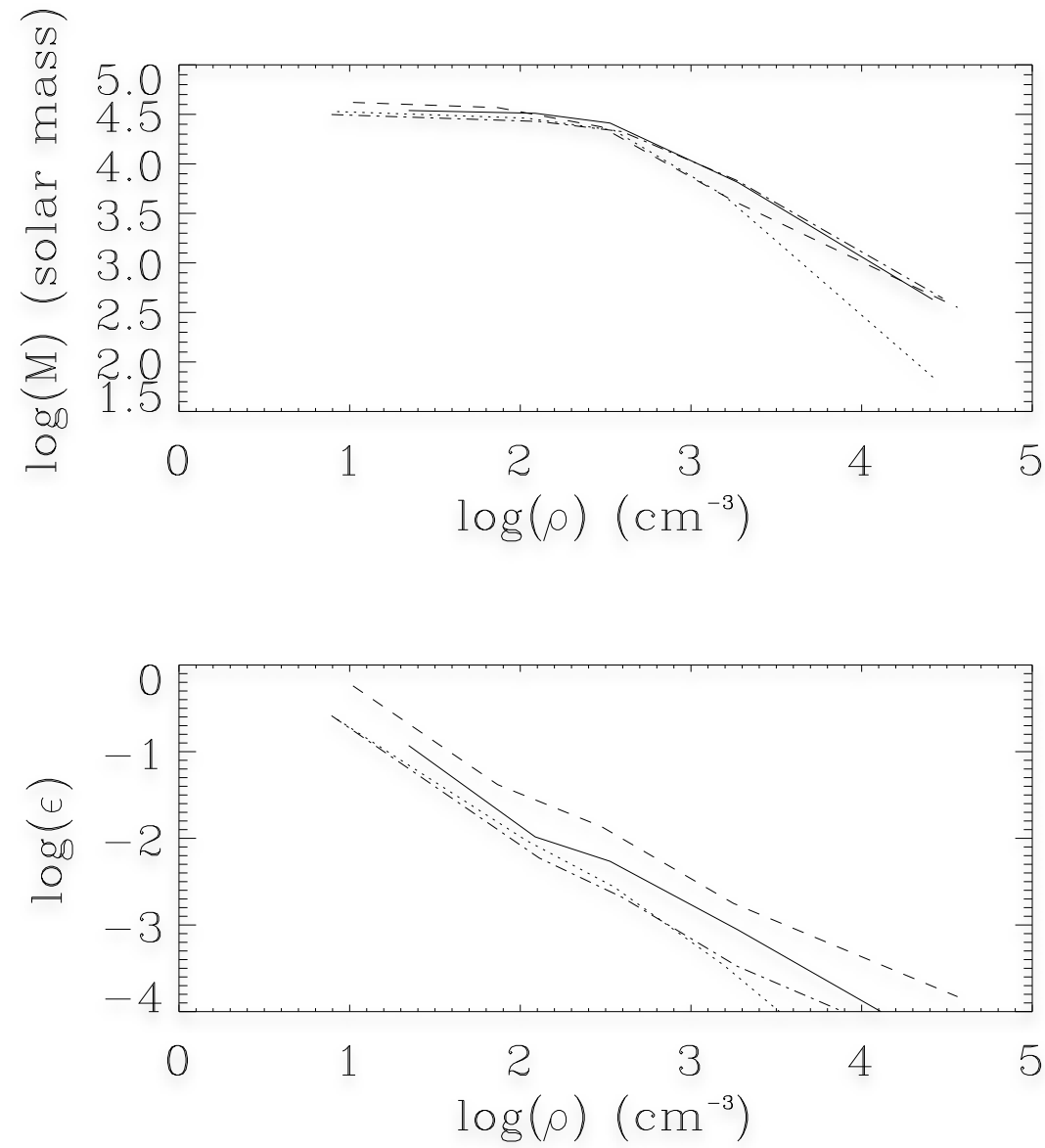
$$\dot{E}_{\text{in}} = \frac{1}{2} \dot{M}_{\text{in}} v_{\text{in}}^2$$

- can both values match, modulo some efficiency?

$$\epsilon = \left| \frac{\dot{E}_{\text{decay}}}{\dot{E}_{\text{in}}} \right|$$



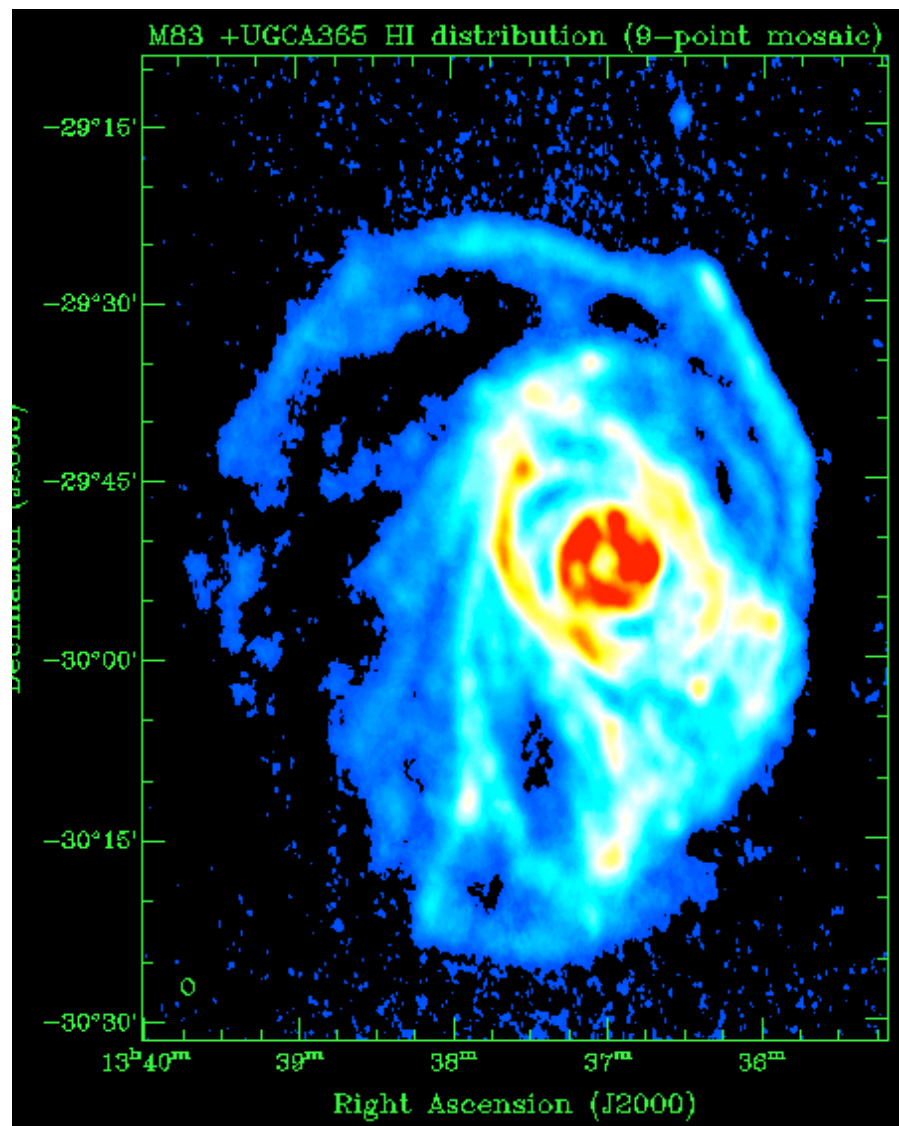
## some estimates from convergent flow studies



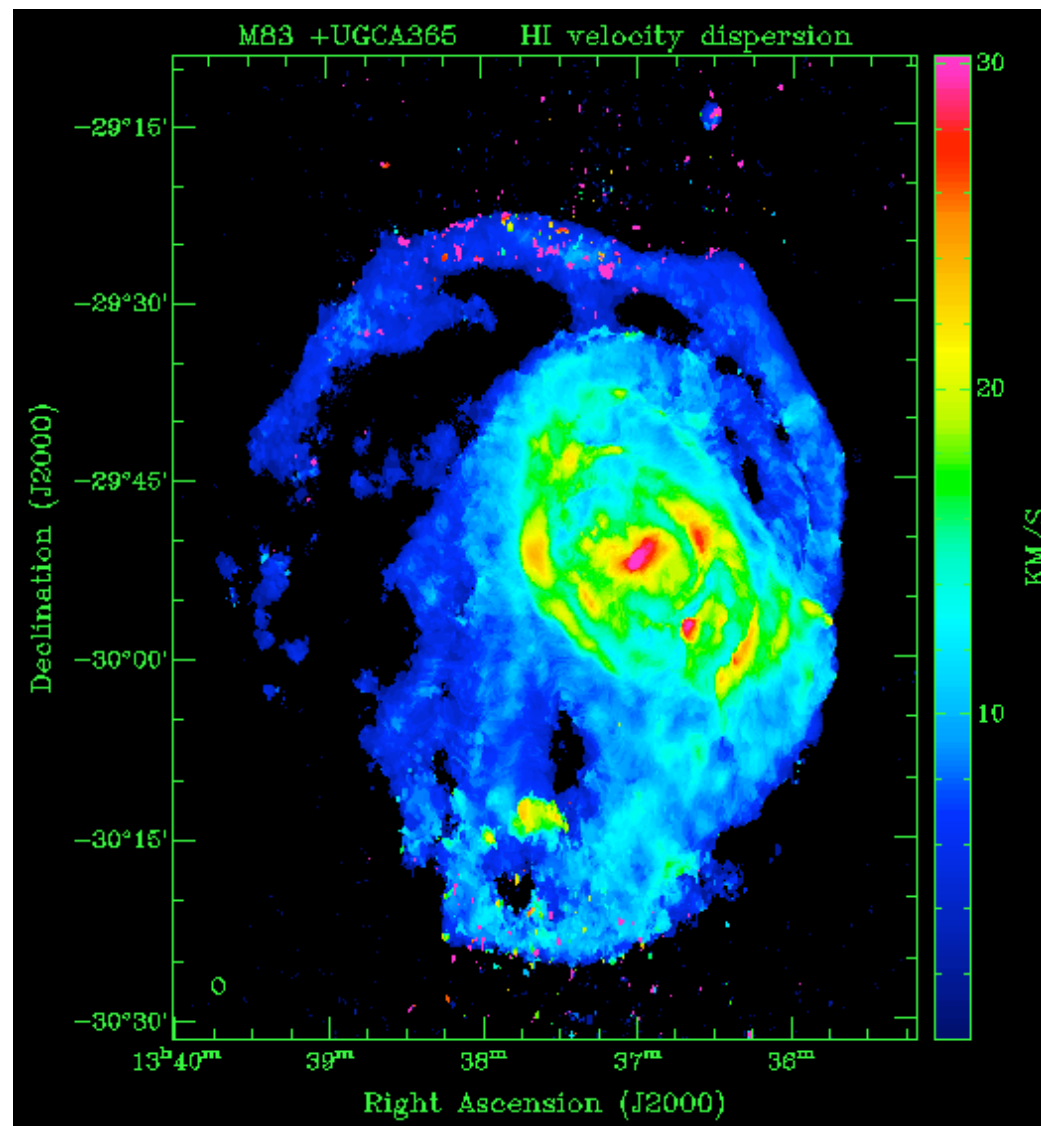
# application to galaxies

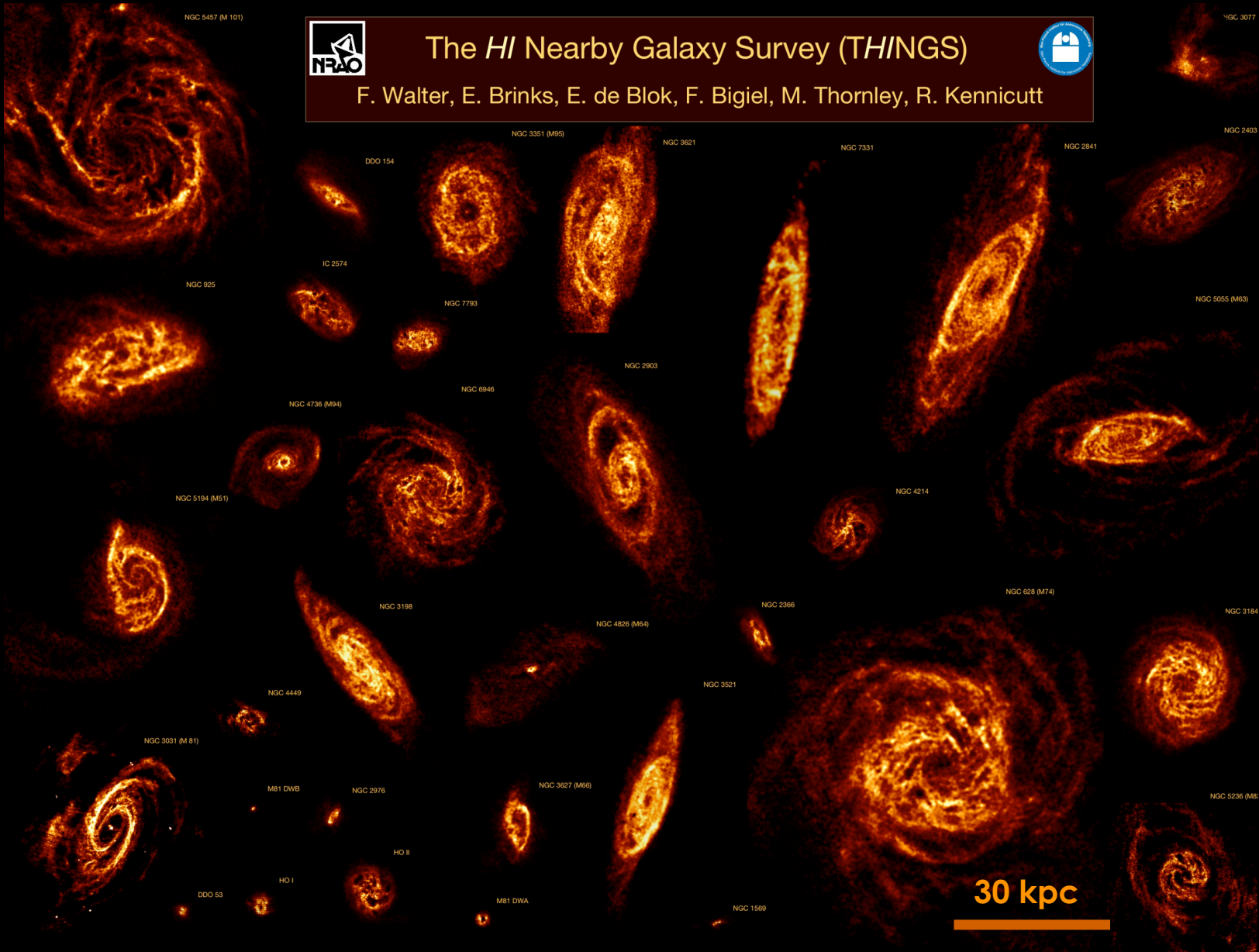
- underlying assumption
  - galaxy is in steady state
    - > accretion rate equals star formation rate
  - what is the required efficiency for the method to work?
- study Milky Way and 11 THINGS
  - excellent observational data in HI:  
velocity dispersion, column density, rotation curve

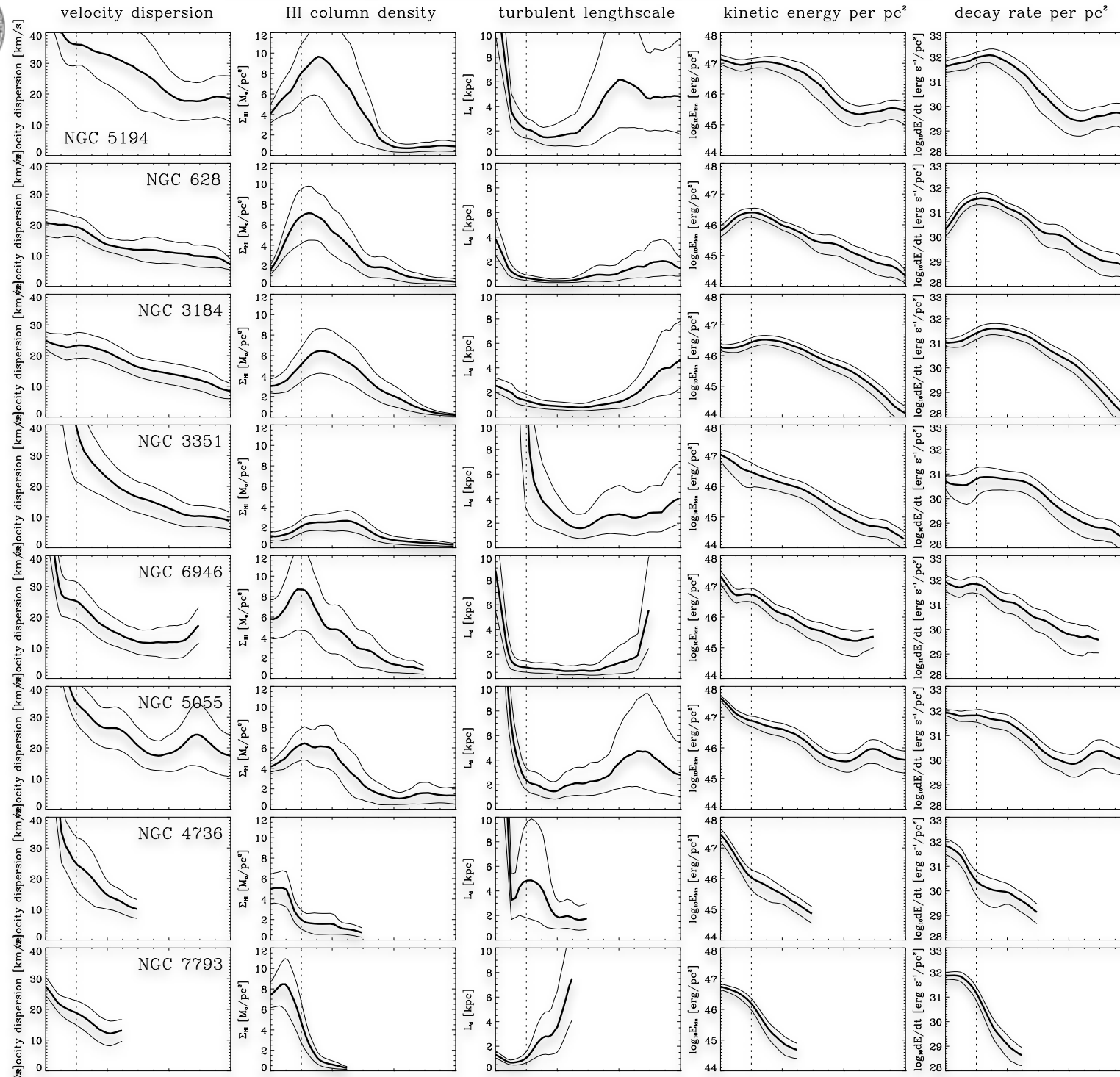
M83 HI column



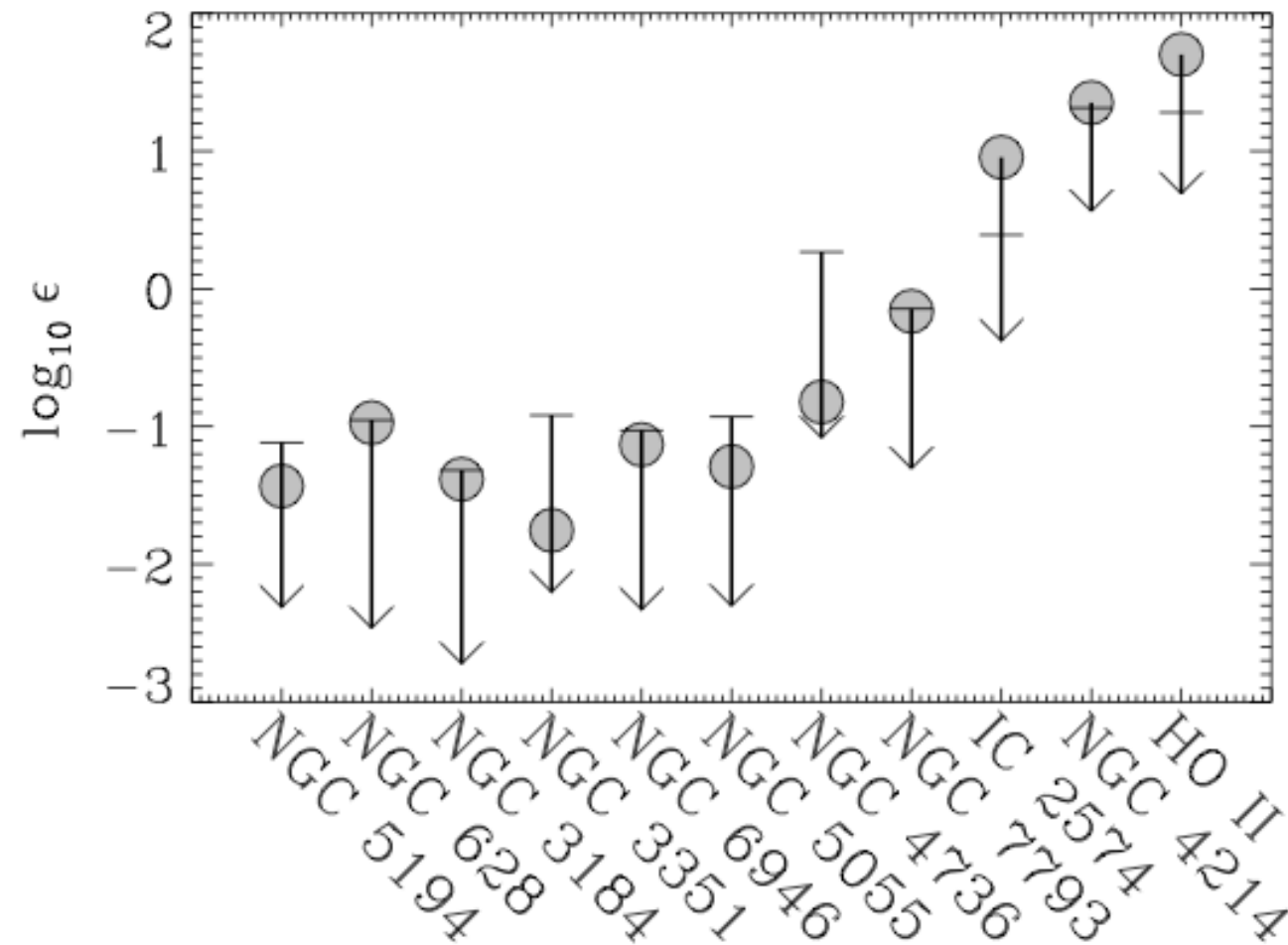
M83 HI velocity dispersion







# 11 THINGS galaxies



# galactic disks

- method works for Milky Way type galaxies:
  - required efficiencies are  $\sim 1\%$  only!
- relevant for outer disks (extended HI disks)
  - there are not other sources of turbulence (certainly not stellar sources, maybe MRI)
- works well for molecular clouds
  - example clouds in the LMC (Fukui et al.)
- potentially interesting for TTS
  - model reproduces  $dM/dt - M$  relation (e.g Natta et al. 2006, Muzerolle et al. 2005, Muhanty et al. 2005, Calvet et al. 2004, etc.)

IMF



# IMF

- distribution of stellar masses depends on
  - turbulent initial conditions
    - > mass spectrum of prestellar cloud cores
  - collapse and interaction of prestellar cores
    - > competitive accretion and  $N$ -body effects
  - thermodynamic properties of gas
    - > balance between heating and cooling
    - > EOS (determines which cores go into collapse)
  - (proto) stellar feedback terminates star formation
    - ionizing radiation, bipolar outflows, winds, SN

(e.g. Larson 2003, Prog. Rep. Phys.; Mac Low & Klessen, 2004, Rev. Mod. Phys, 76, 125 - 194)



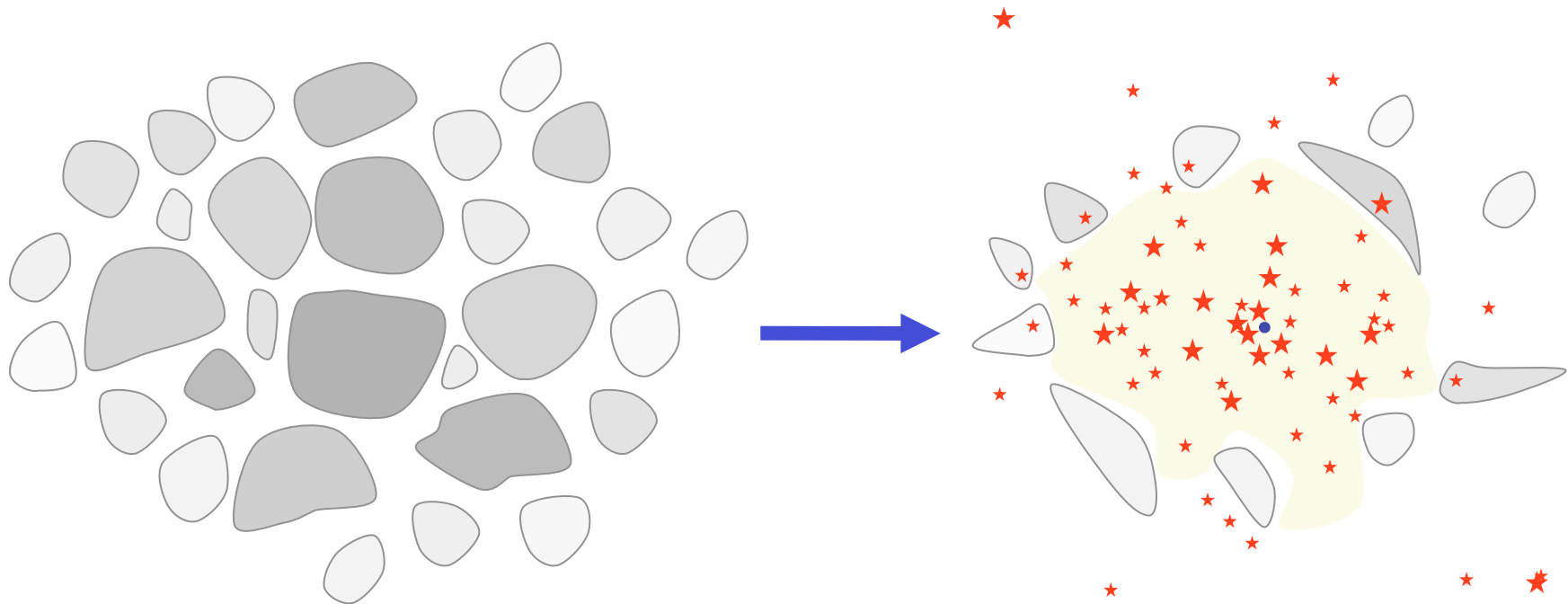
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(e.g. Larson 2003, Prog. Rep. Phys.; Mac Low & Klessen, 2004, Rev. Mod. Phys, 76, 125 - 194)

# Star cluster formation

Most stars form in clusters → *star formation = cluster formation*



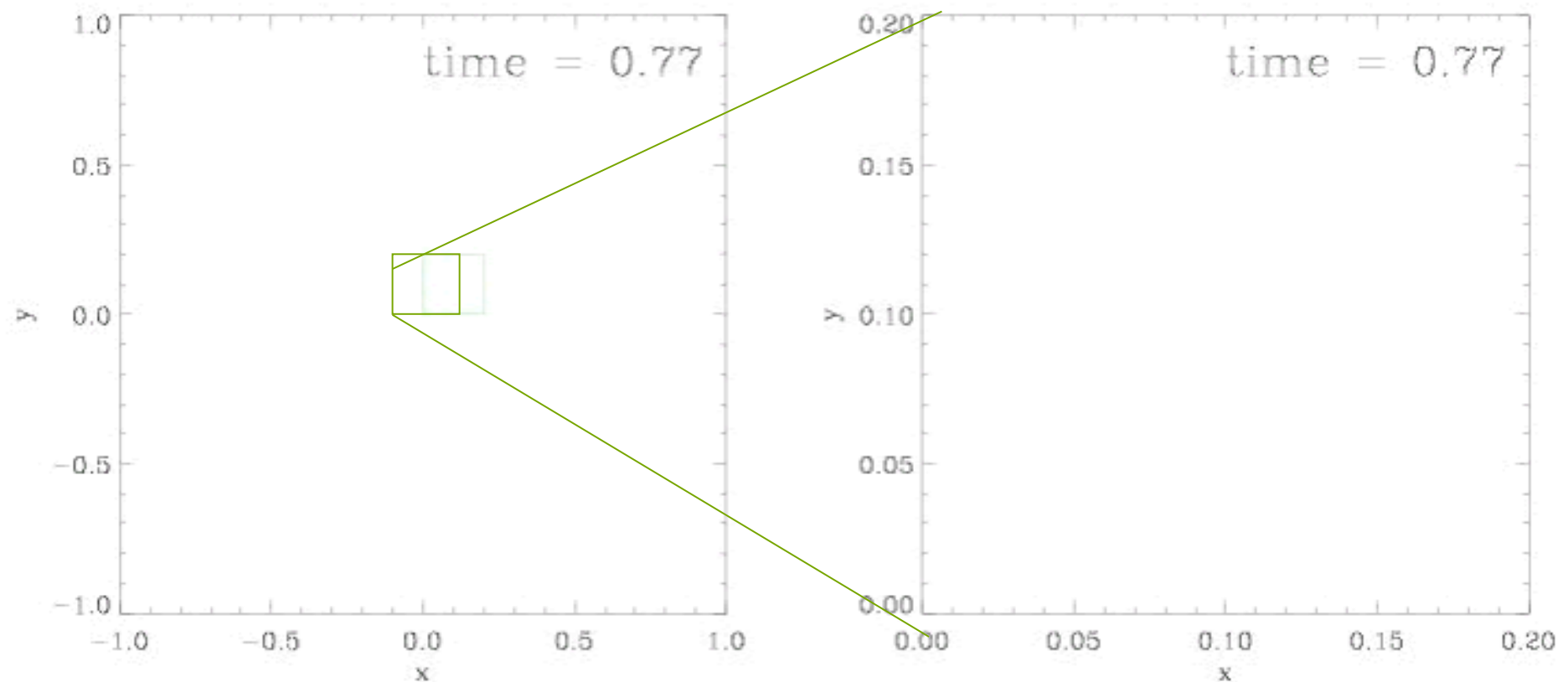
How to get from **cloud cores** to **star clusters**?

How do the stars **acquire mass**?

(e.g. Larson 2003, Prog. Rep. Phys.; Mac Low & Klessen, 2004, Rev. Mod. Phys, 76, 125 - 194)

# Dynamics of nascent star cluster

in dense clusters protostellar interaction may become important!

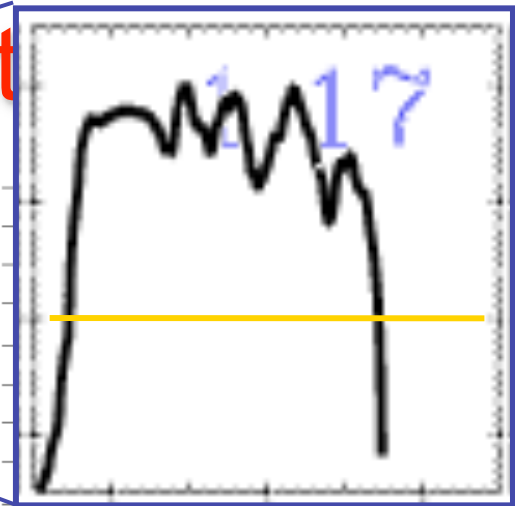
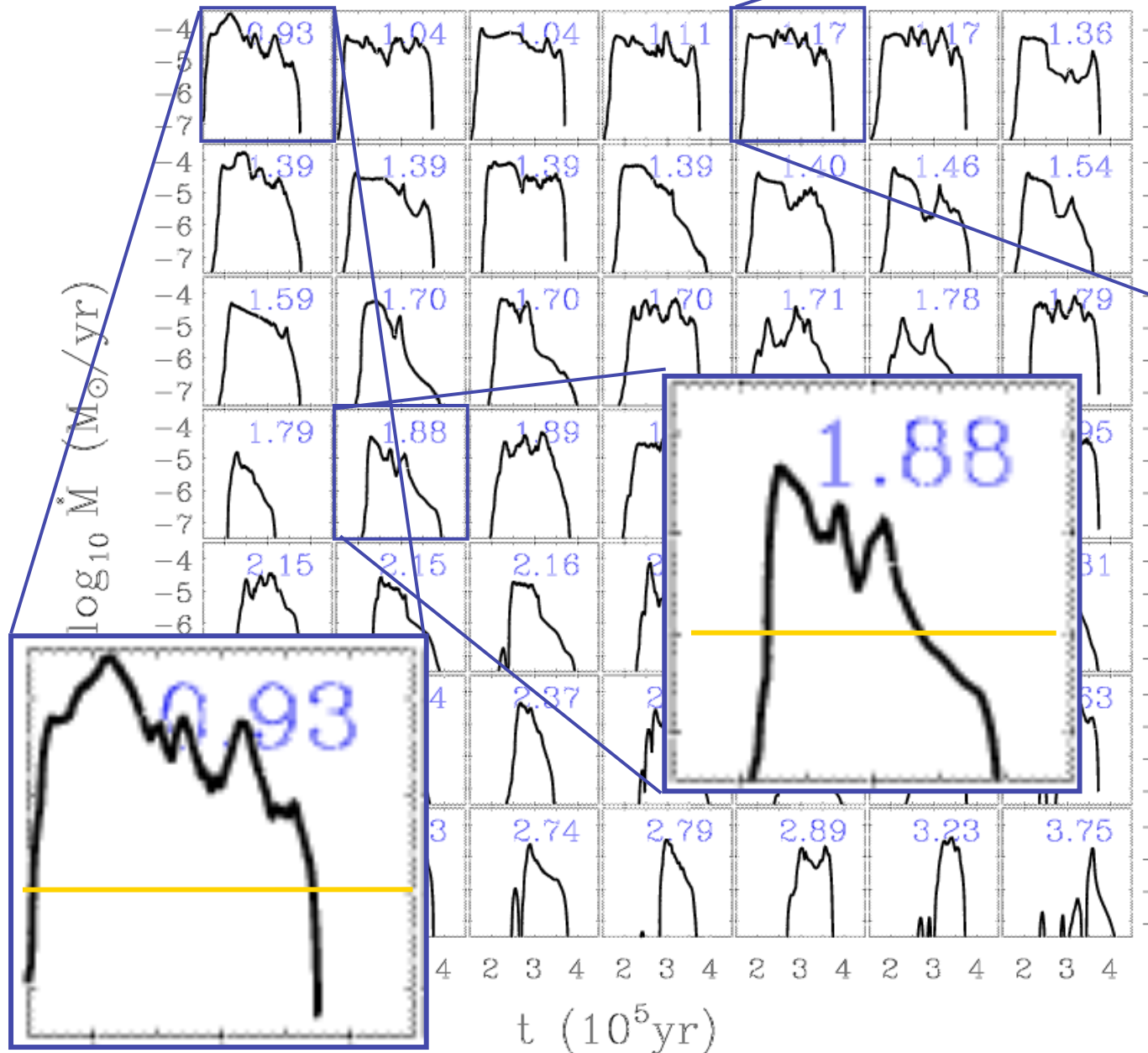


Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation

(from Klessen & Burkert 2000, ApJS, 128, 287)

Ralf Klessen: PKU/KIAA, 17.03.2008

# Accretion rates in clust



Mass accretion rates *vary with time* and are strongly *influenced* by the *cluster environment*.

(Klessen 2001, ApJ, 550, L77;  
also Schmeja & Klessen,  
2004, A&A, 419, 405)



# IMF

- distribution of stellar masses depends on
  - turbulent initial conditions
    - > mass spectrum of prestellar cloud cores
  - collapse and interaction of prestellar cores
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  - *thermodynamic properties of gas*
    - > *balance between heating and cooling*
    - > *EOS (determines which cores go into collapse)*
  - (proto) stellar feedback terminates star formation
    - ionizing radiation, bipolar outflows, winds, SN



# dependency on EOS

- degree of fragmentation depends on *EOS*!

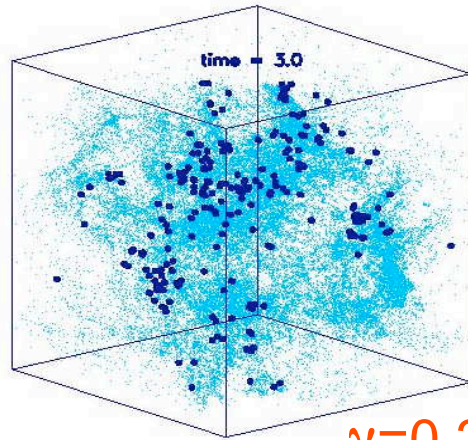
polytropic EOS:  $p \propto \rho^\gamma$

- $\gamma < 1$ : dense cluster of low-mass stars
- $\gamma > 1$ : isolated high-mass stars

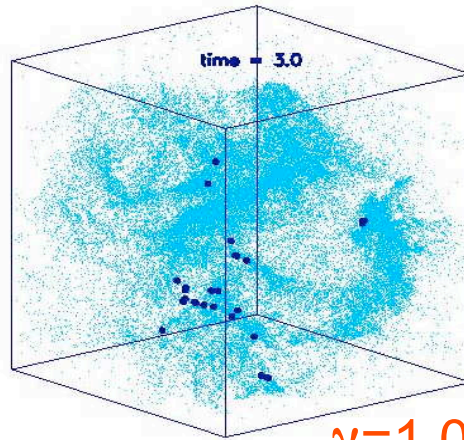
(see Li, Klessen, & Mac Low 2003, ApJ, 592, 975; also Kawachi & Hanawa 1998, Larson 2003)



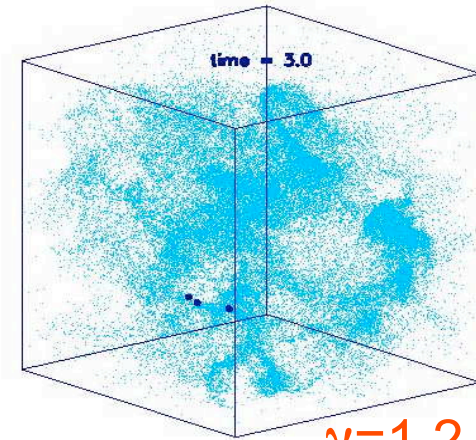
# dependency on EOS



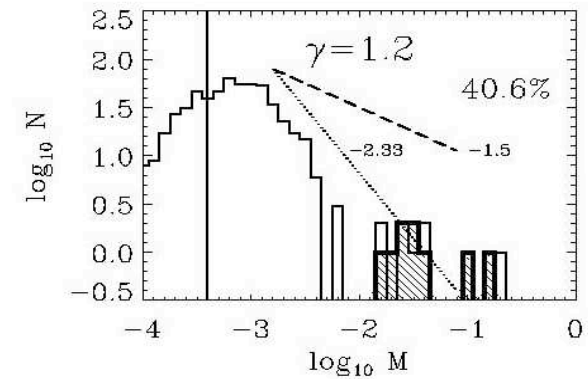
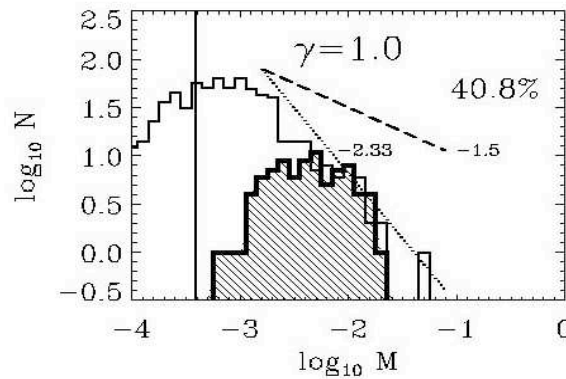
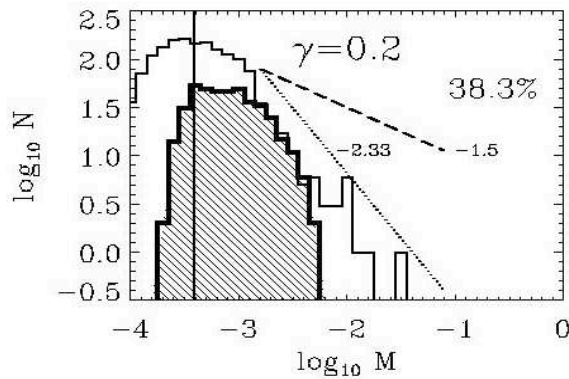
$\gamma=0.2$



$\gamma=1.0$



$\gamma=1.2$



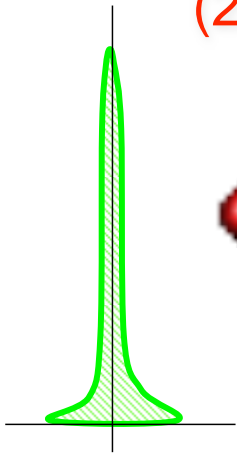
for  $\gamma < 1$  fragmentation is enhanced  $\rightarrow$  *cluster of low-mass stars*  
for  $\gamma > 1$  it is suppressed  $\rightarrow$  formation of *isolated massive stars*



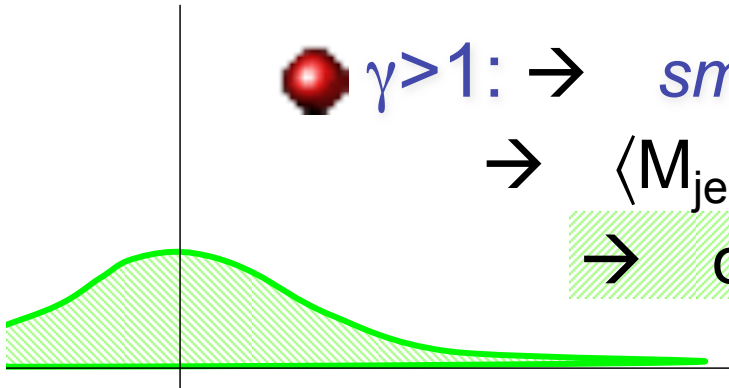
# how does that work?

(1)  $p \propto \rho^\gamma \rightarrow \rho \propto p^{1/\gamma}$

(2)  $M_{\text{jeans}} \propto \gamma^{3/2} \rho^{(3\gamma-4)/2}$



- $\gamma < 1$ :  $\rightarrow$  *large* density excursion for given pressure  
 $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  becomes small  
 $\rightarrow$  number of fluctuations with  $M > M_{\text{jeans}}$  is large

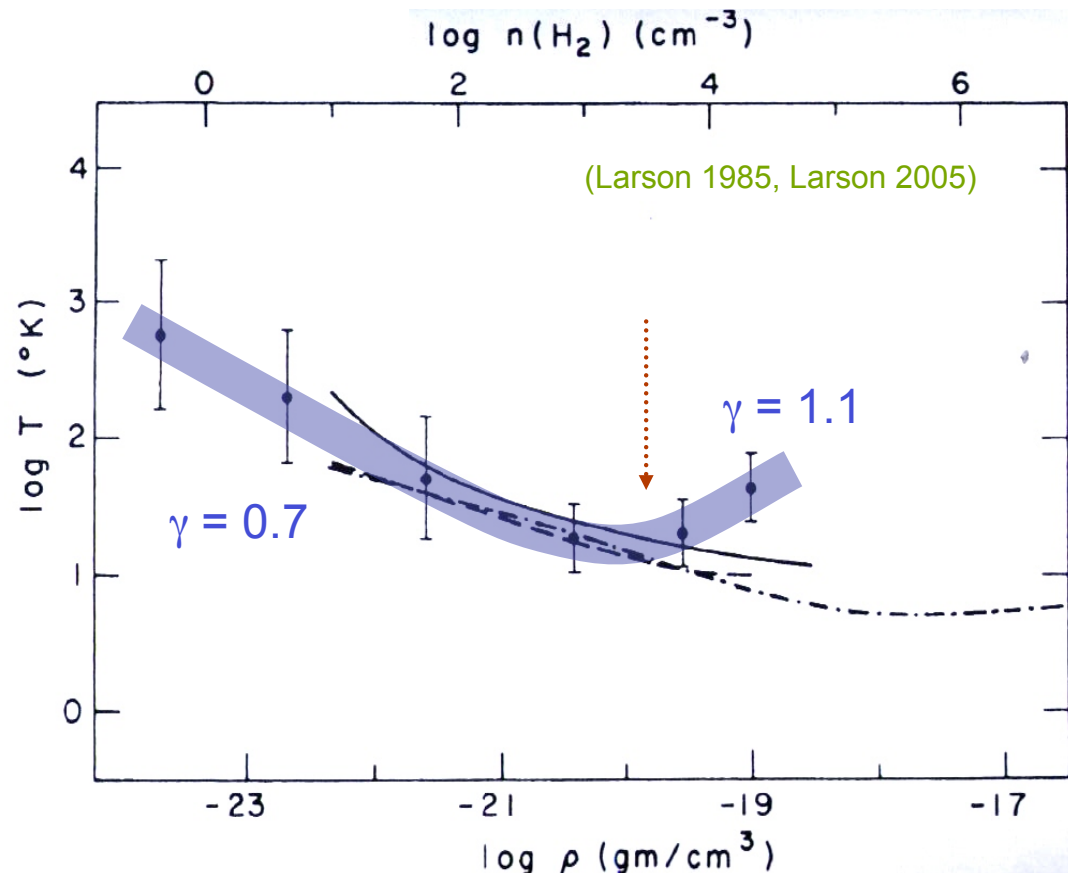


- $\gamma > 1$ :  $\rightarrow$  *small* density excursion for given pressure  
 $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  is large  
 $\rightarrow$  only few and massive clumps exceed  $M_{\text{jeans}}$



# EOS for solar neighborhood

below  $10^{-18} \text{ gcm}^{-3}$ :  $\rho$   $\uparrow$   $\rightarrow$   $\downarrow$   
above  $10^{-18} \text{ gcm}^{-3}$ :  $\rho$   $\uparrow$   $\rightarrow$   $\uparrow$



$$P \propto \rho^\gamma$$

$$P \propto \rho T$$

$$\rightarrow \gamma = 1 + d\ln T / d\ln \rho$$

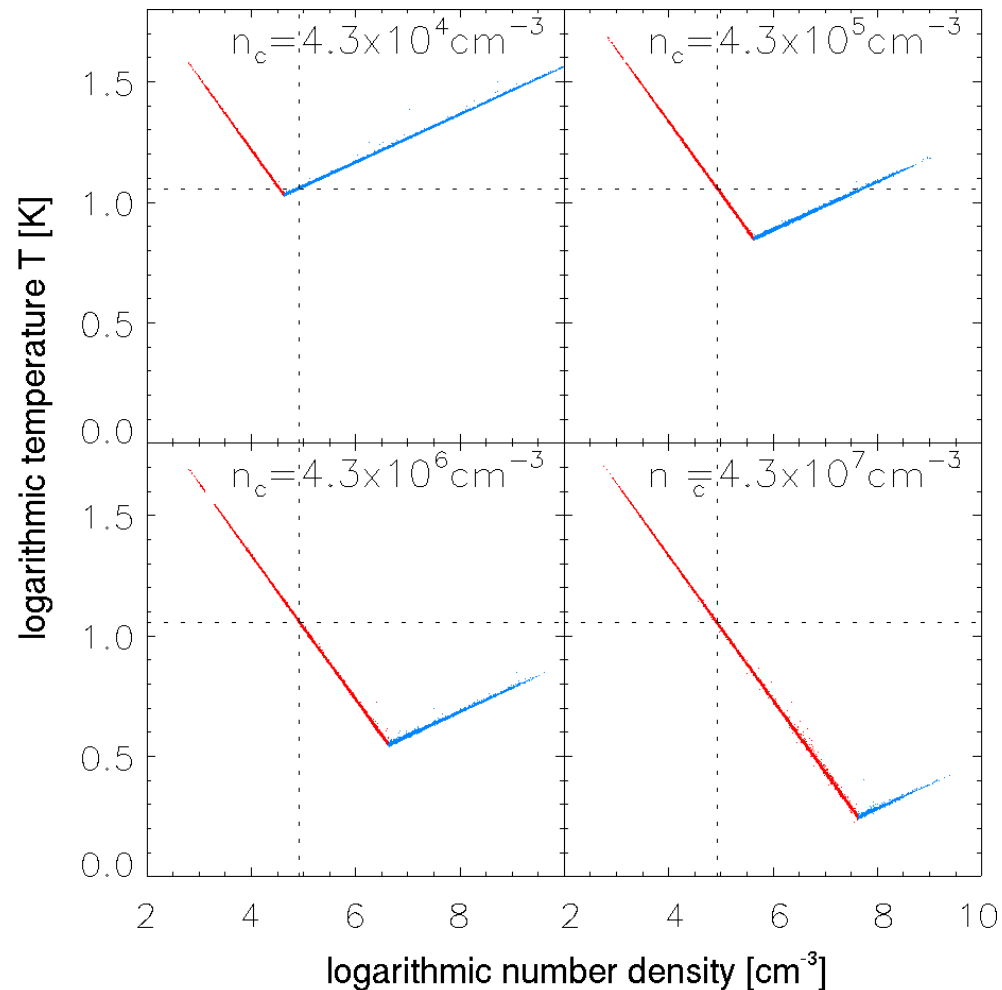


# IMF from simple piece-wise polytropic EOS

$$\gamma_1 = 0.7$$

$$\gamma_2 = 1.1$$

$$T \sim \rho^{\gamma-1}$$

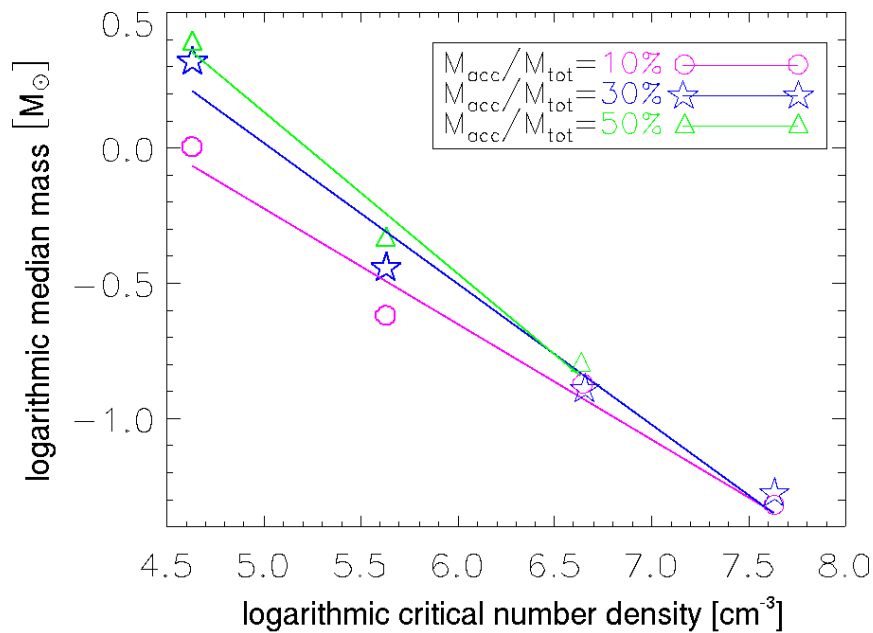


(Jappsen et al. 2005)

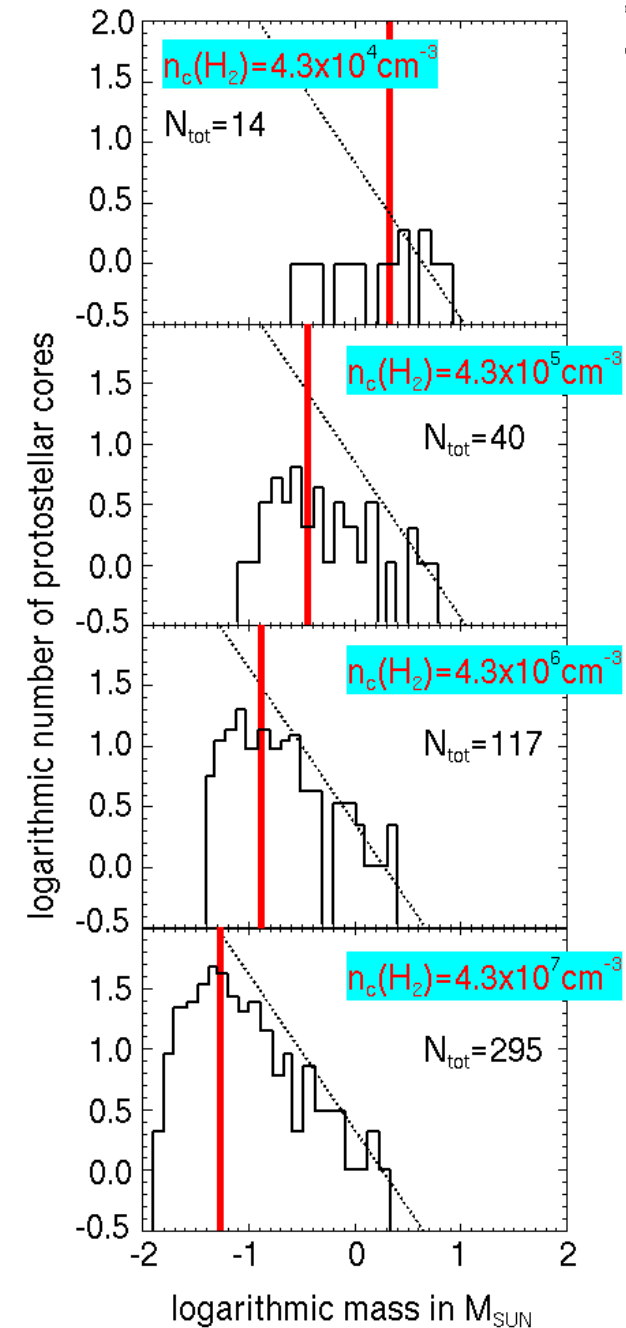


# IMF from simple piece-wise EOS

critical density  $\uparrow$   $\longrightarrow$  median mass

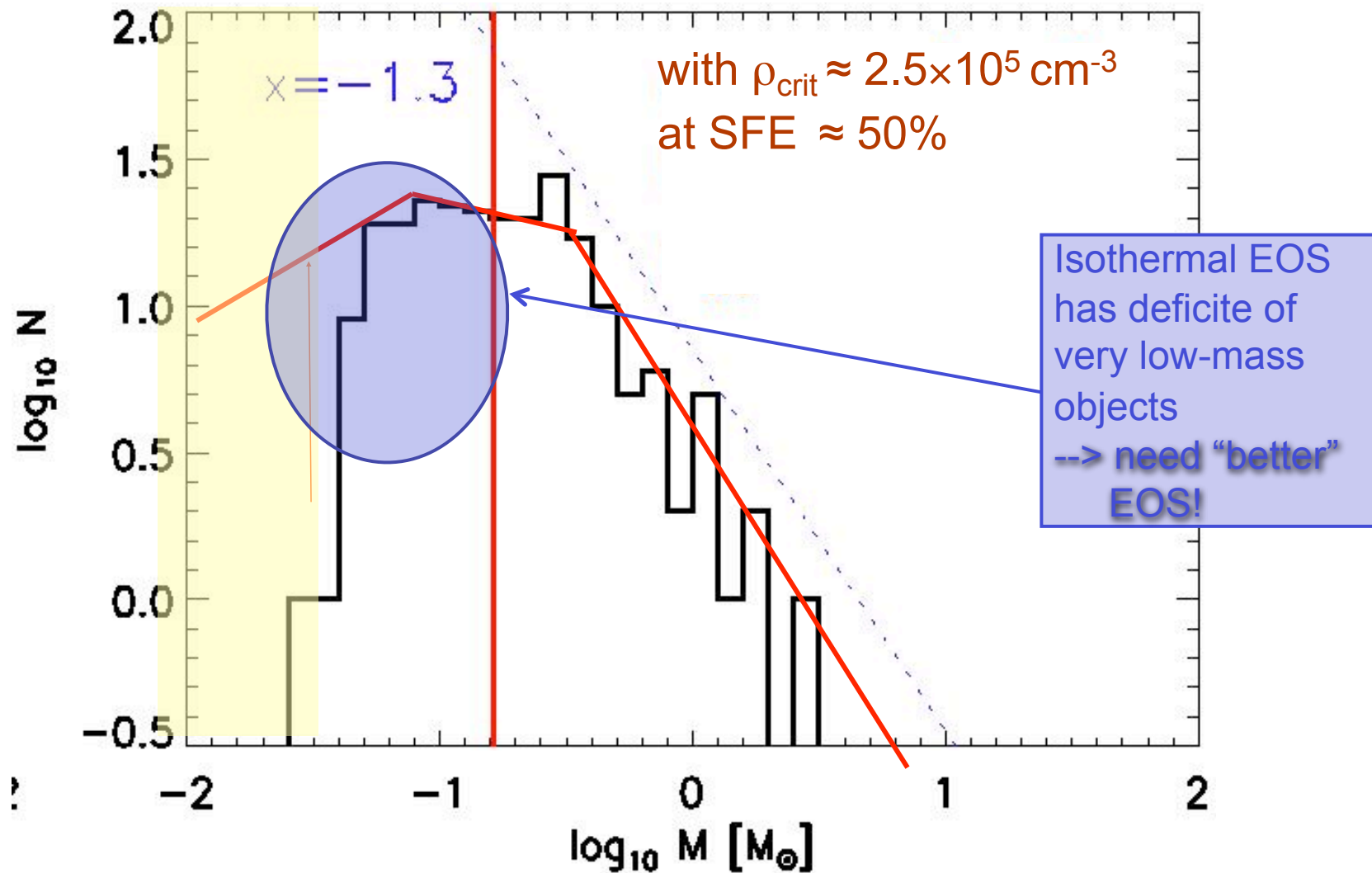


(Jappsen et al. 2005)





# IMF in nearby molecular clouds





# B-fields



# how to make strong B-fields

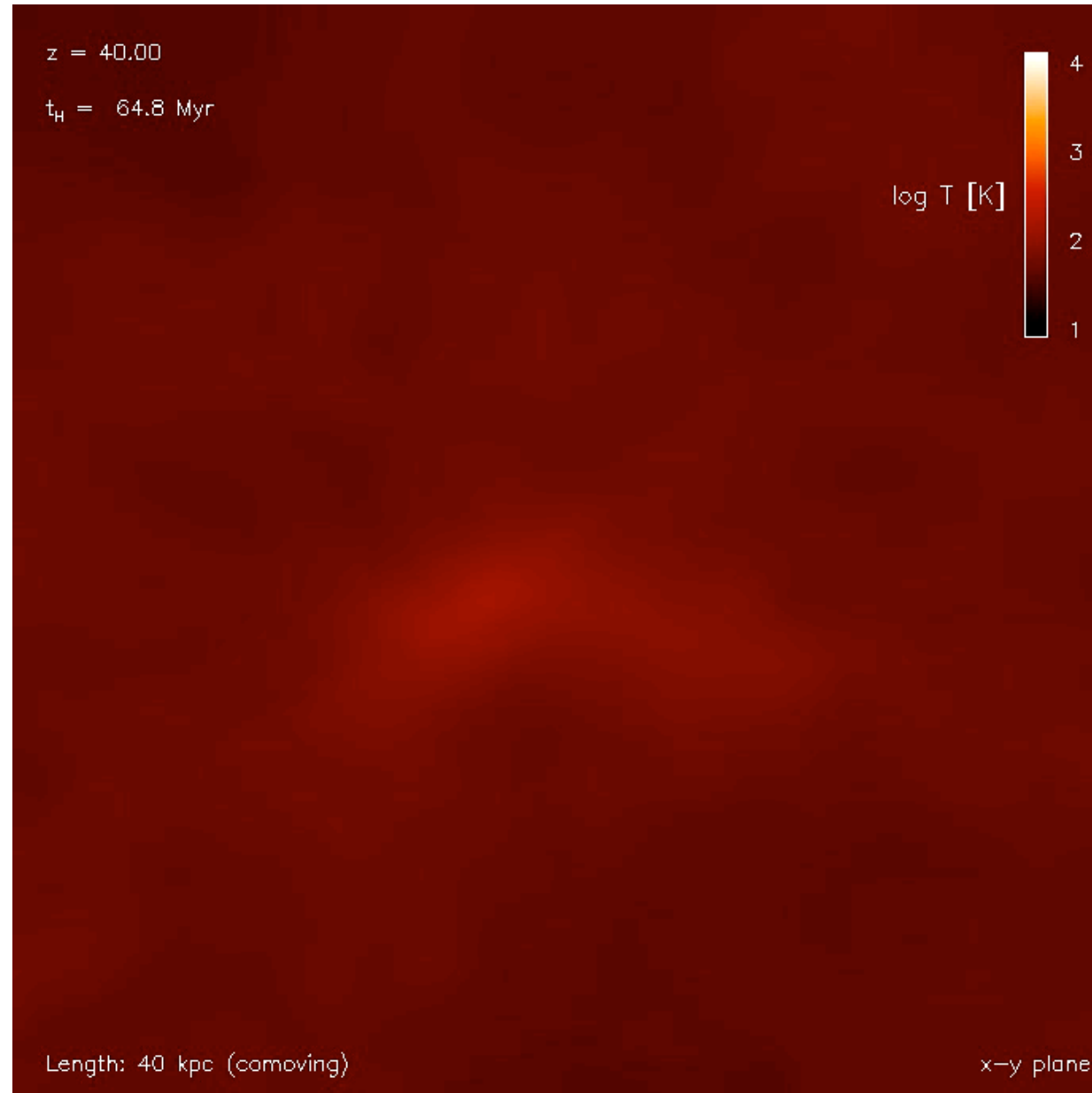
- we know the universe is magnetized (now)
- knowledge about B-fields in the high-redshift universe is extremely uncertain
  - inflation / QCD phase transition / Biermann battery / Weibel instability
- they are thought to be extremely small
- however, THIS MAY BE WRONG!



# small-scale turbulent dynamo

- *idea*: the small-scale turbulent dynamo can generate strong magnetic fields from very small seed fields
- *approach*: model collapse of primordial gas ---> formation of the first stars in low-mass halo at redshift  $z \sim 20$
- *method*: solve ideal MHD equations with very high resolution
  - grid-based AMR code FLASH
  - resolution up to  $128^3$  cells per Jeans volume (effective resolution  $65536^3$  cells)

# turbulence developing in an atomic cooling halo

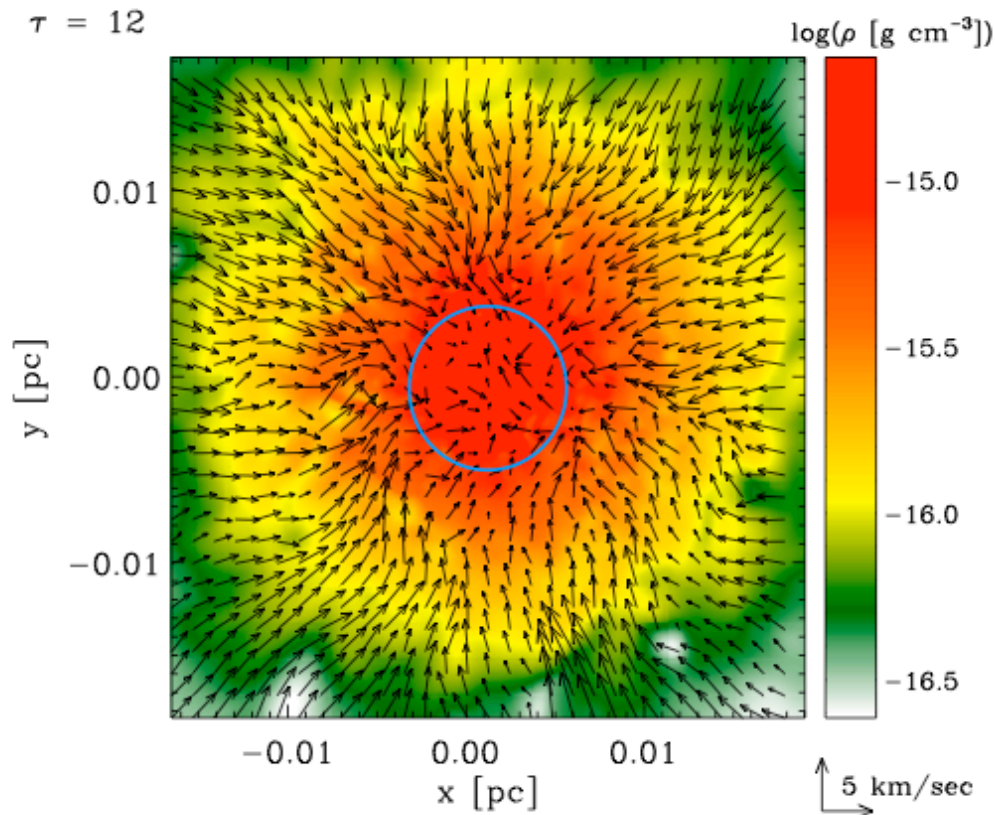


see also talk  
by *Thomas Greif*

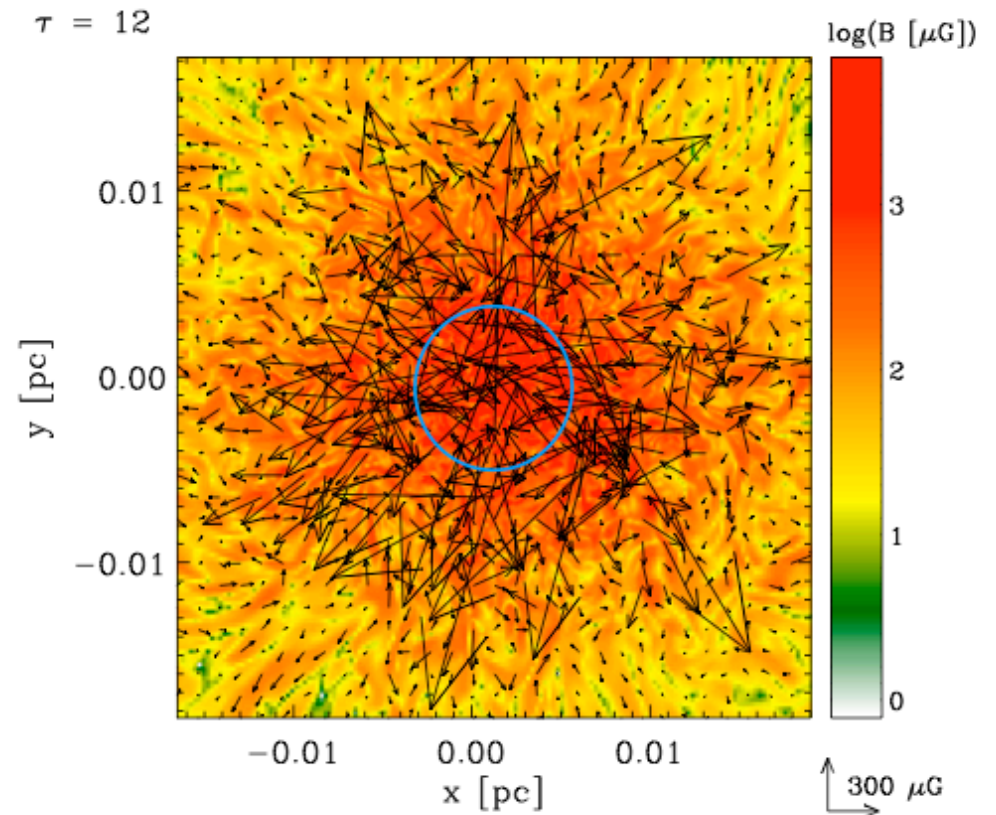
(Greif et al. 2008, MNRAS, 387, 1021, see also Wise & Abel 2007)



# small-scale turbulent dynamo



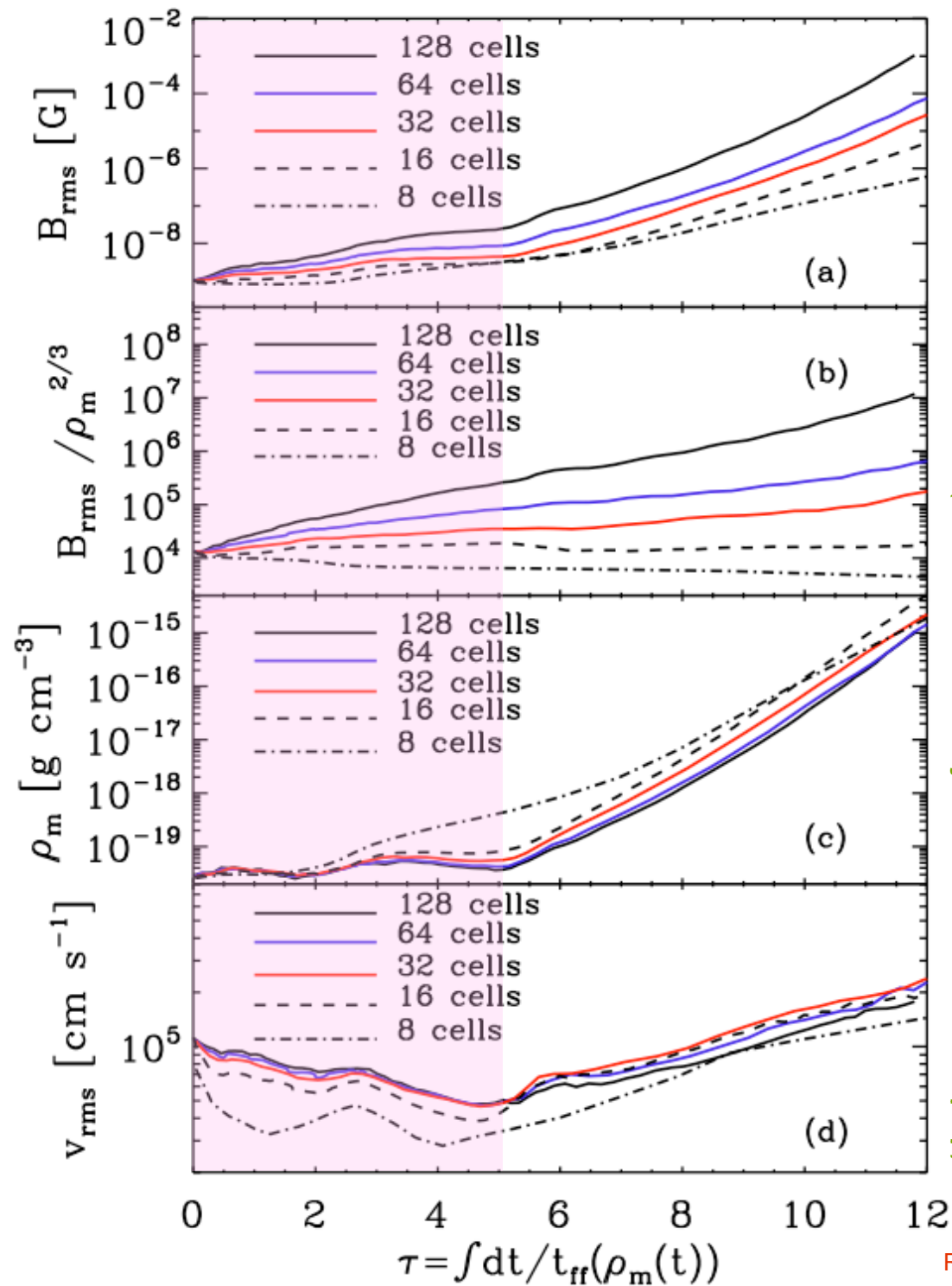
density, xy-velocity



B-field amplitude, xy-component



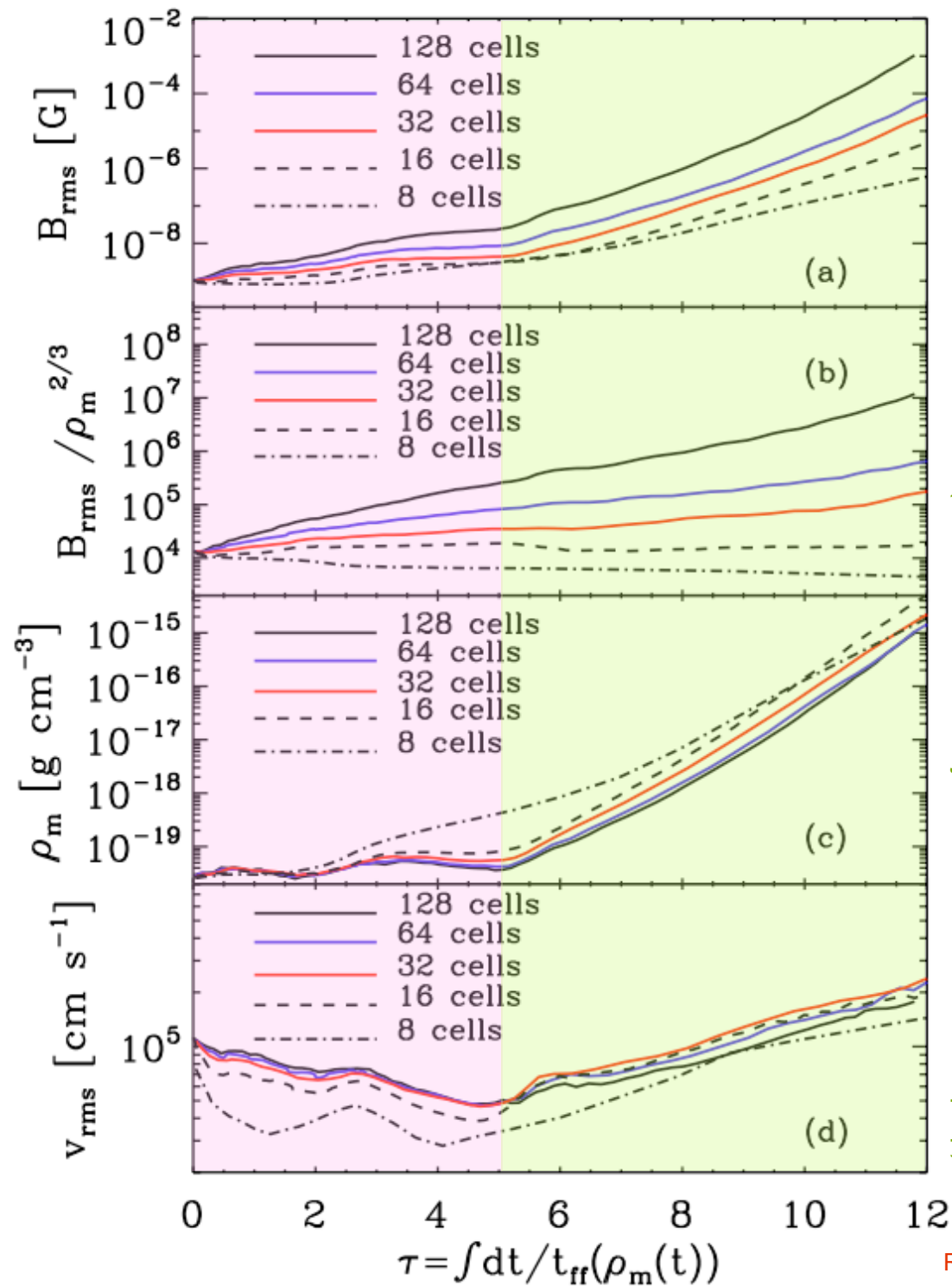
# decay of turbulence



(Sur, Schleicher, Banerjee, Federrath, Klessen, in prep)



# decay of turbulence

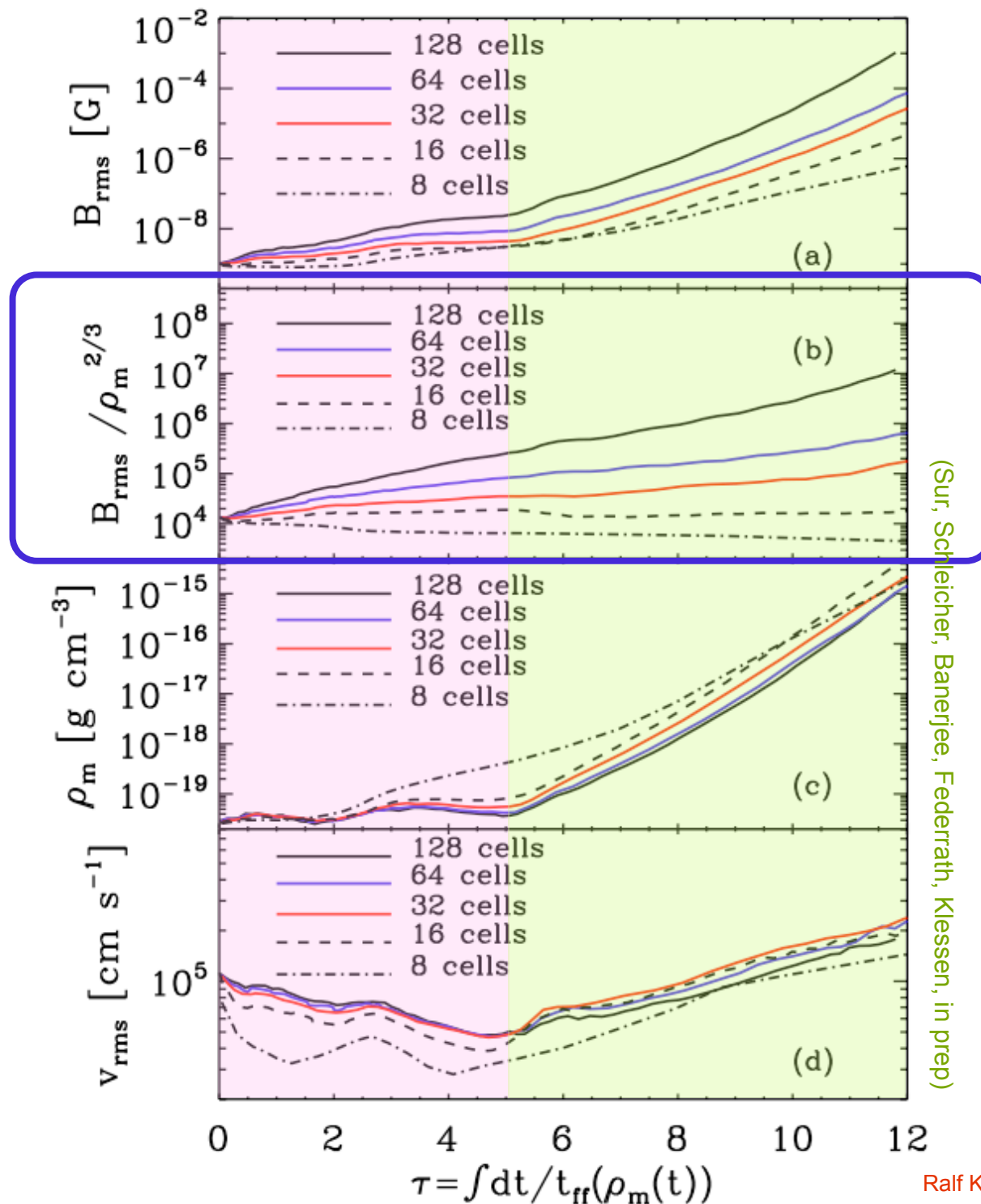


(Sur, Schleicher, Banerjee, Federrath, Klessen, in prep)

gravitational collapse,  
accretion driven turbulence



# decay of turbulence



gravitational collapse,  
accretion driven turbulence

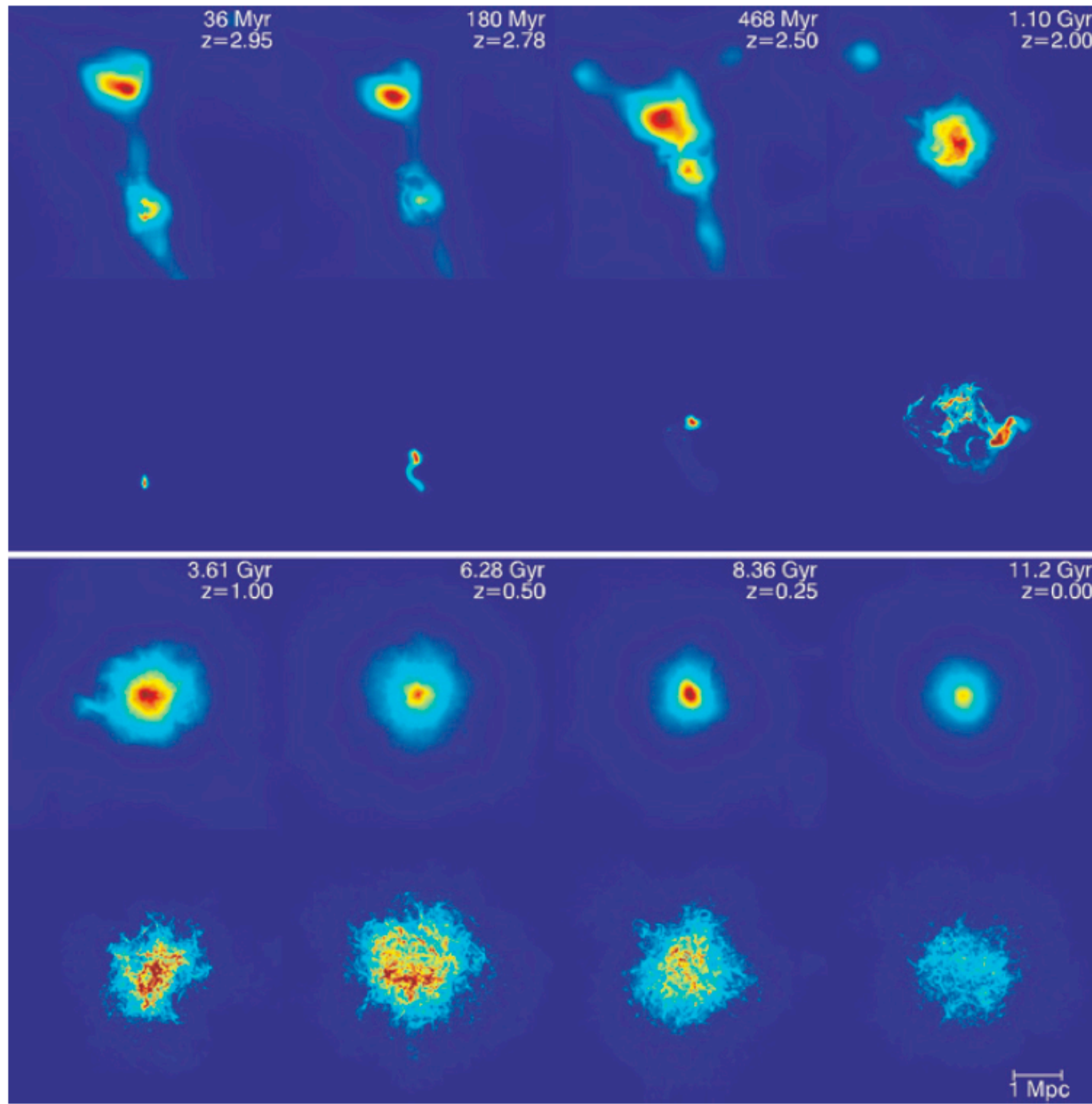


# small-scale turbulent dynamo

- small-scale turbulent dynamo can generate dynamically significant B-fields in the during first star formation!
- expected saturation level  $\sim 10\% - 20\%$  of equipartition value (Subramanian & Brandenburg 2005)
- needs very high resolution to be seen:

cells per Jeans length	max. refinement	max. effective resolution	min.cell size [cm]	mean growth rate $\Gamma$
8	10	4096	$1.17 \times 10^{16}$	—
16	11	8192	$5.86 \times 10^{15}$	0.021
32	12	16384	$2.93 \times 10^{15}$	0.108
64	13	32768	$1.46 \times 10^{15}$	0.188
128	14	65536	$7.32 \times 10^{14}$	0.25

# similar processes on scales of galaxy clusters



density

magnetic  
energy  
density

density

magnetic  
energy  
density



diffusion

# Turbulent diffusion I

- Observations of young star clusters exhibit an enormous degree of chemical homogeneity  
(e.g. in the Pleiades: Wilden et al. 2002)
- Star-forming gas must be well mixed.
- How does this constrain models of interstellar turbulence?
- → Study mixing in supersonic compressible turbulence....

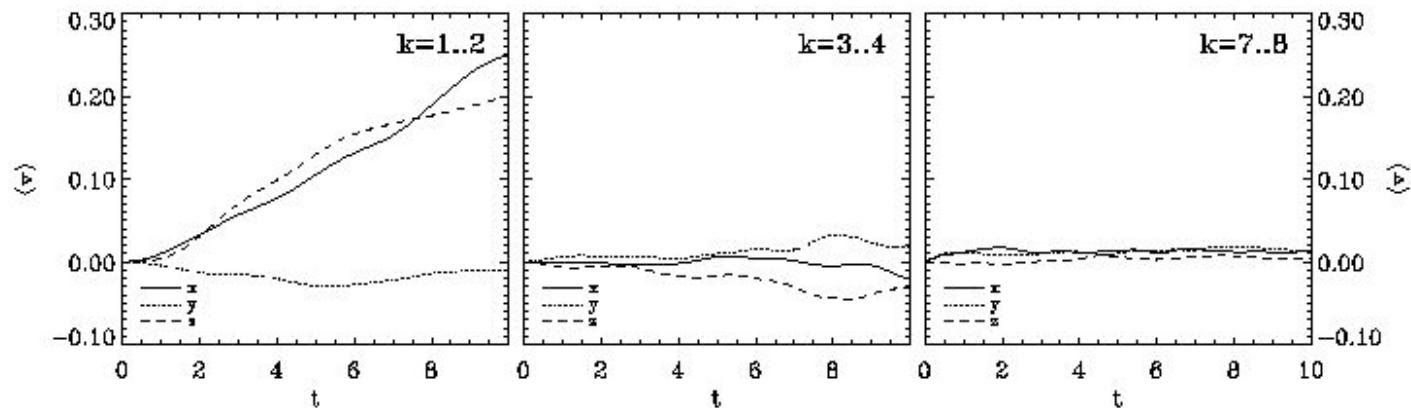
$$\xi_r^2(t-t') = \langle [\vec{r}_i(t) - \vec{r}_i(t')]^2 \rangle_i$$

mean particle displacement  
for a given time interval

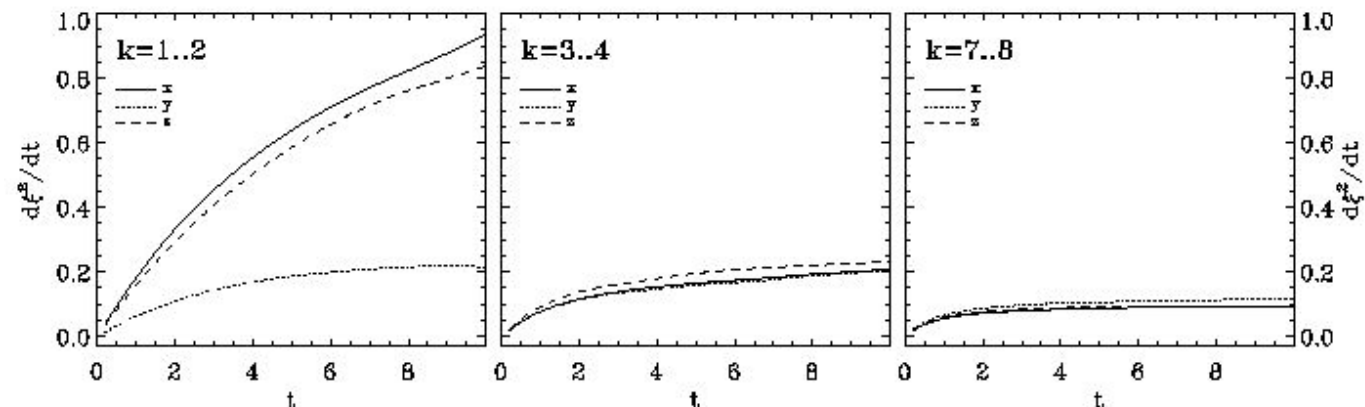
$$\frac{\partial n}{\partial t} = D \nabla^2 n \quad D(t) = \frac{d\xi_r^2(t)}{dt} = 2 \langle \vec{r}_i(t) \cdot \vec{v}_i(t) \rangle_i$$

# Turbulent diffusion II

- Large-scale turbulence associated with bulk motion.



- Super-diffusive behavior.



# Turbulent diffusion III

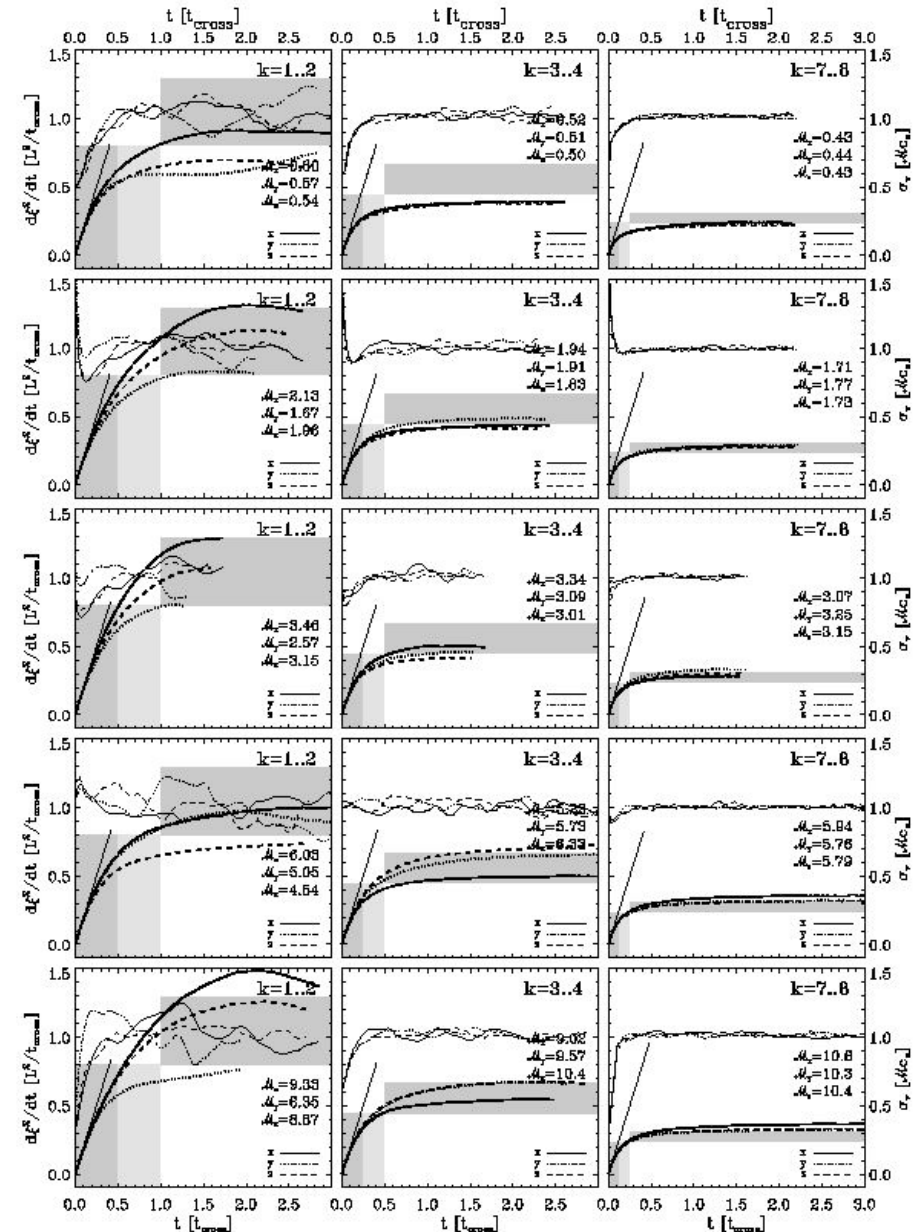
- Mean-motion corrected diffusion

- Simple mixing-length approach:

- $D(t) \approx v_{\text{rms}}^2 t \quad t < \tau$

- $D(t) \approx v_{\text{rms}}^2 \tau$   
 $= v_{\text{rms}} \ell \quad t > \tau$

- With  $v_{\text{rms}}$  = rms velocity  
and  $\ell = L/k =$  shock sep.





# Summary



# summary

- modeling ISM turbulence & star formation is feasible
  - more physics is needed (chemistry, radiation field, realistic sources of turbulence, etc.)
- caveat: all LES of ISM turbulence model “crude oil” not interstellar gas
- examples
  - stellar mass spectrum
  - molecular cloud formation
  - generation of the first dynamically significant magnetic fields in the universe
  - effective diffusivity in isotropic supersonic turbulence



thanks