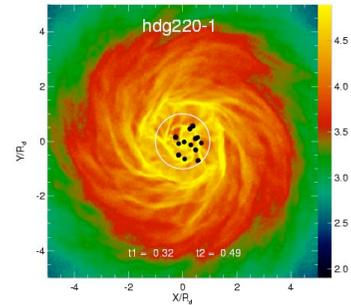
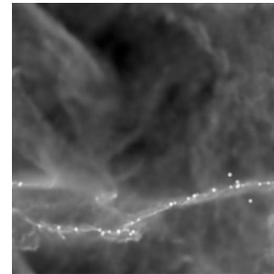
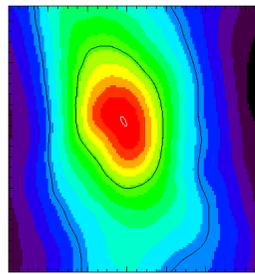
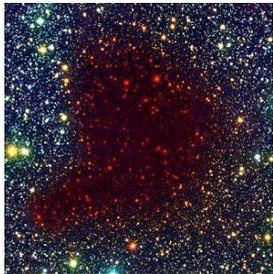


# Star formation today and in the early universe



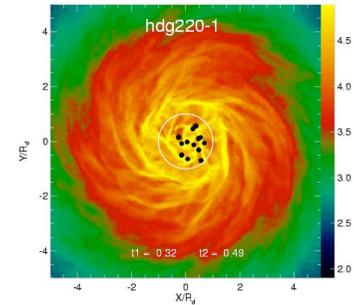
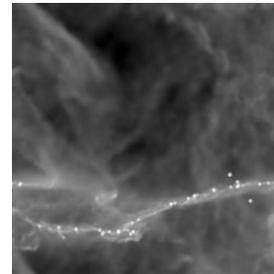
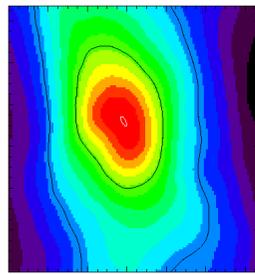
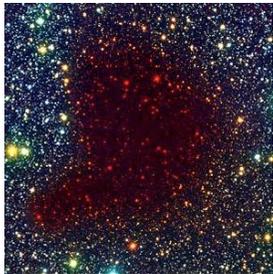
**Ralf Klessen**



Zentrum für Astronomie der Universität Heidelberg  
Institut für Theoretische Astrophysik



# What can we learn from present-day star formation about the first stars?



**Ralf Klessen**



Zentrum für Astronomie der Universität Heidelberg  
Institut für Theoretische Astrophysik



# thanks to ...



... people in the group in Heidelberg:

Robi Banerjee, Simon Glover, Rahul Shetty, Sharanya Sur, Daniel Seifried, Milica Milosavljevic, Florian Mandl, Christian Baczynski, Rowan Smith, Gustavo Dopcke, Jonathan Downing, Jayanta Dutta, Faviola Molina, Christoph Federrath, Erik Bertram, Lukas Konstandin, Paul Clark, Stefan Schmeja, Ingo Berentzen, Thomas Peters, Hsiang-Hsu Wang

... many collaborators abroad!



Deutsche  
Forschungsgemeinschaft  
**DFG**

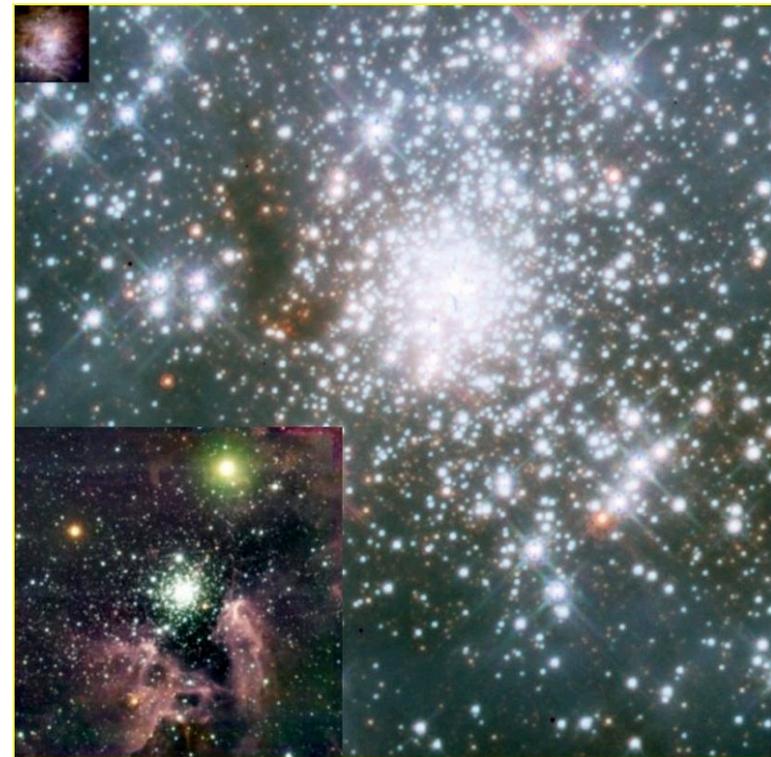
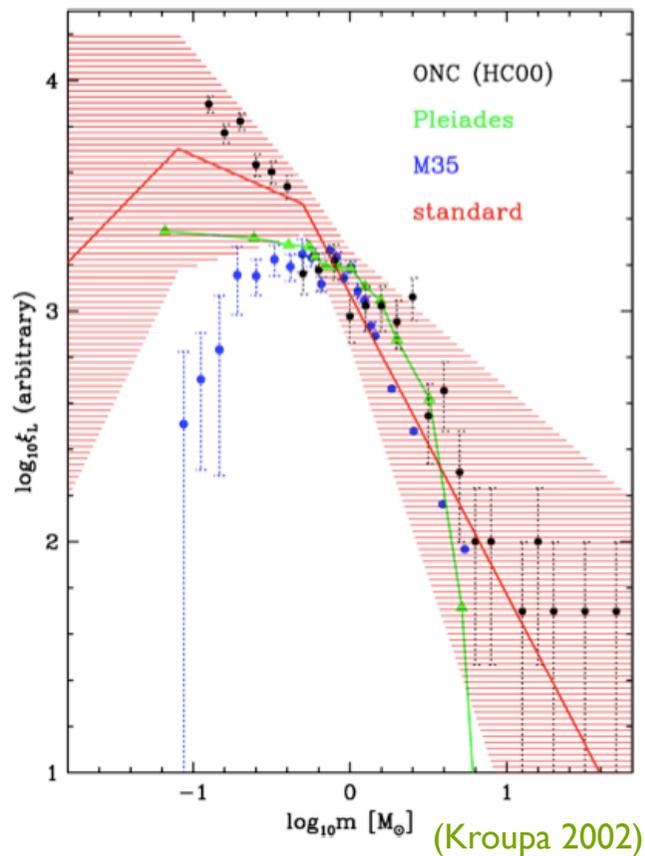




NGC 602 in the LMC: Hubble Heritage Image

# stellar mass function

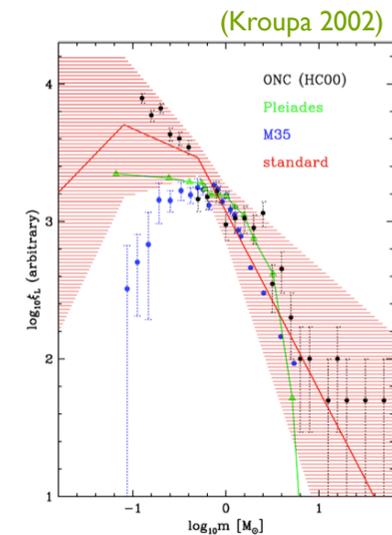
stars seem to follow a universal mass function at birth --> IMF



Orion, NGC 3603, 30 Doradus  
(Zinnecker & Yorke 2007)

# stellar masses

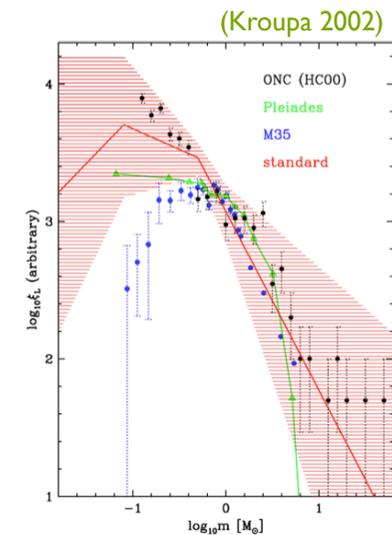
- distribution of stellar masses depends on
  - turbulent initial conditions
    - > mass spectrum of prestellar cloud cores
  - collapse and interaction of prestellar cores
    - > accretion and  $N$ -body effects
  - thermodynamic properties of gas
    - > balance between heating and cooling
    - > EOS (determines which cores go into collapse)
  - (proto) stellar feedback terminates star formation
    - ionizing radiation, bipolar outflows, winds, SN



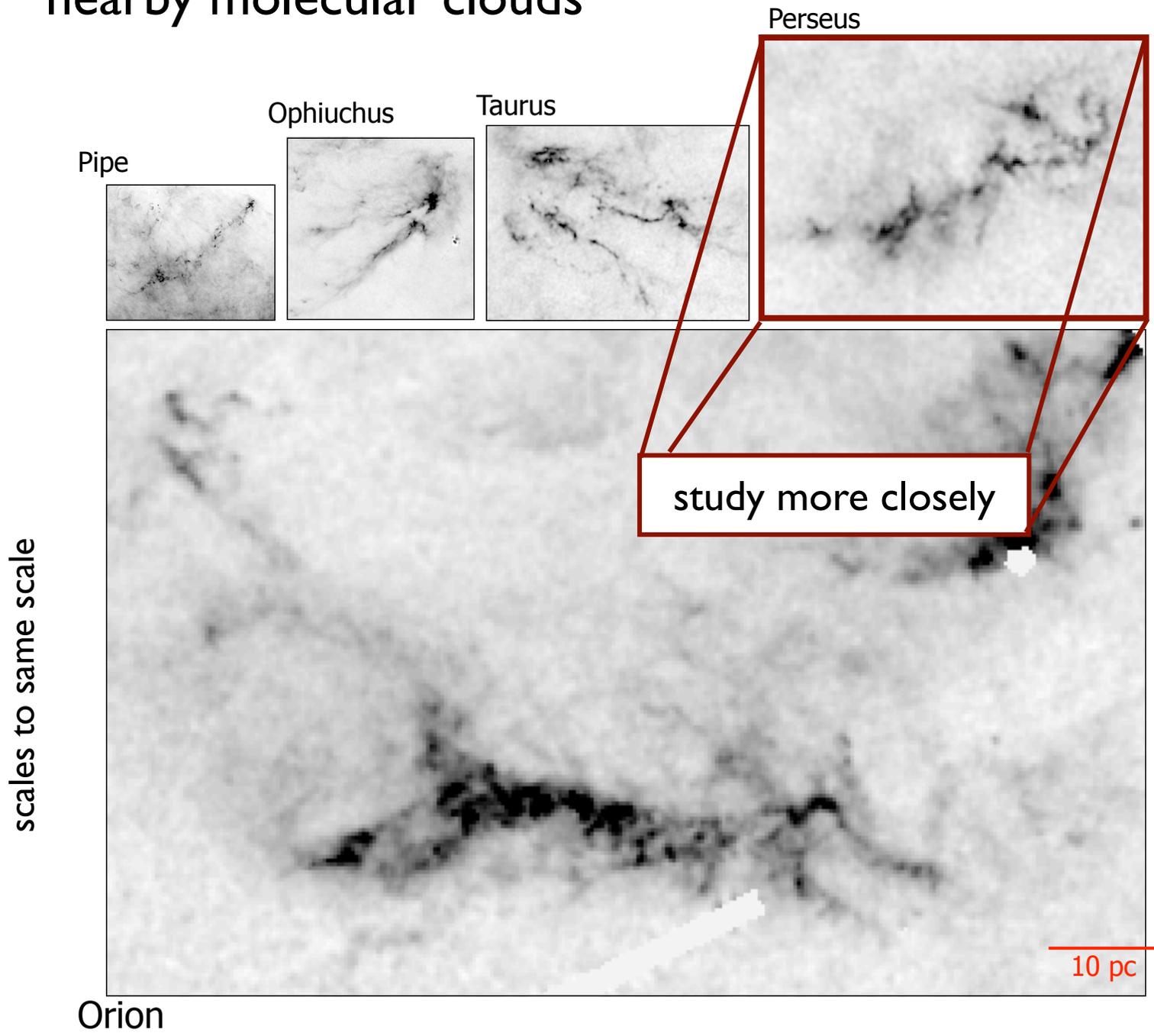
# stellar masses

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# nearby molecular clouds



scales to same scale

(from A. Goodman)

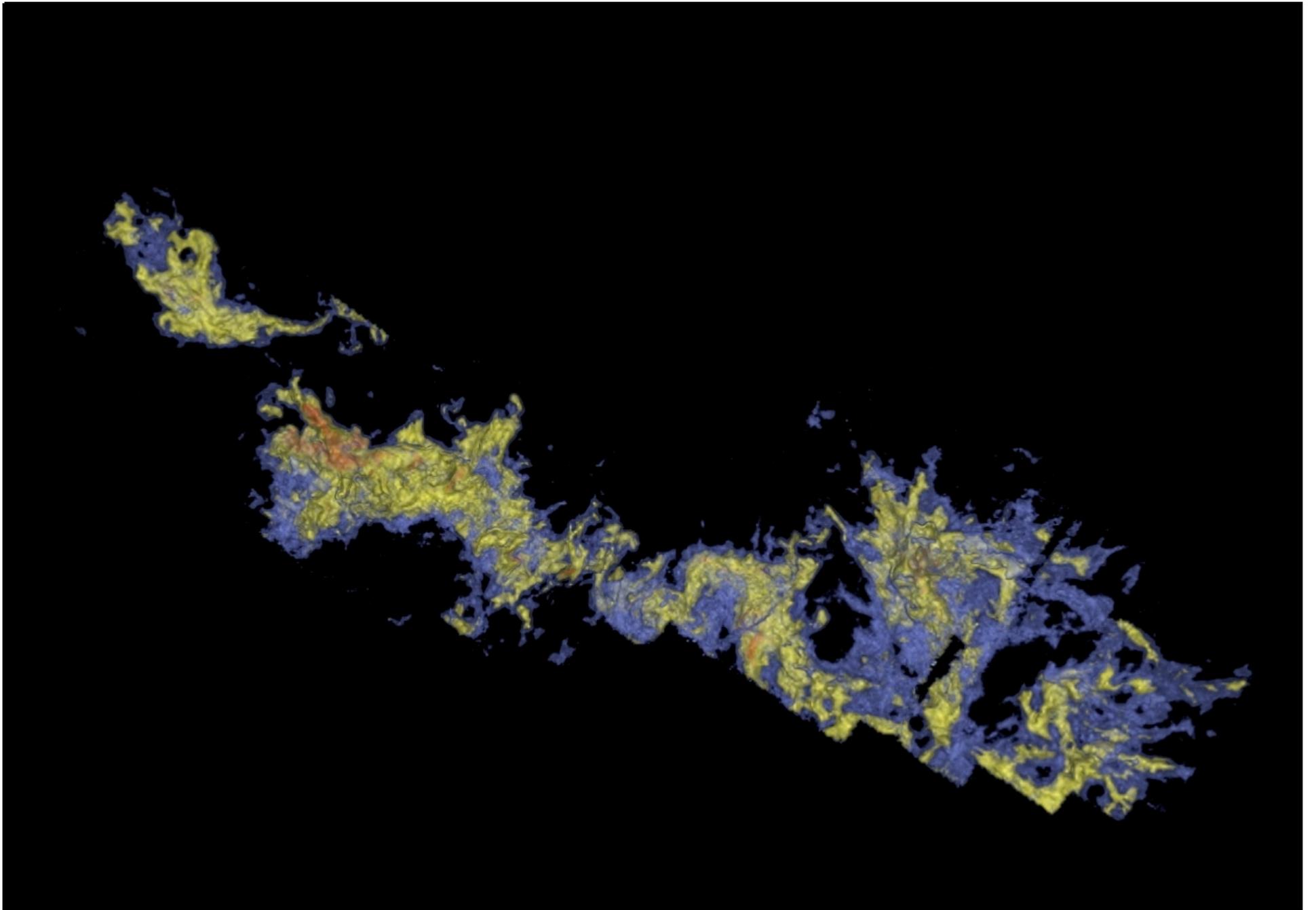
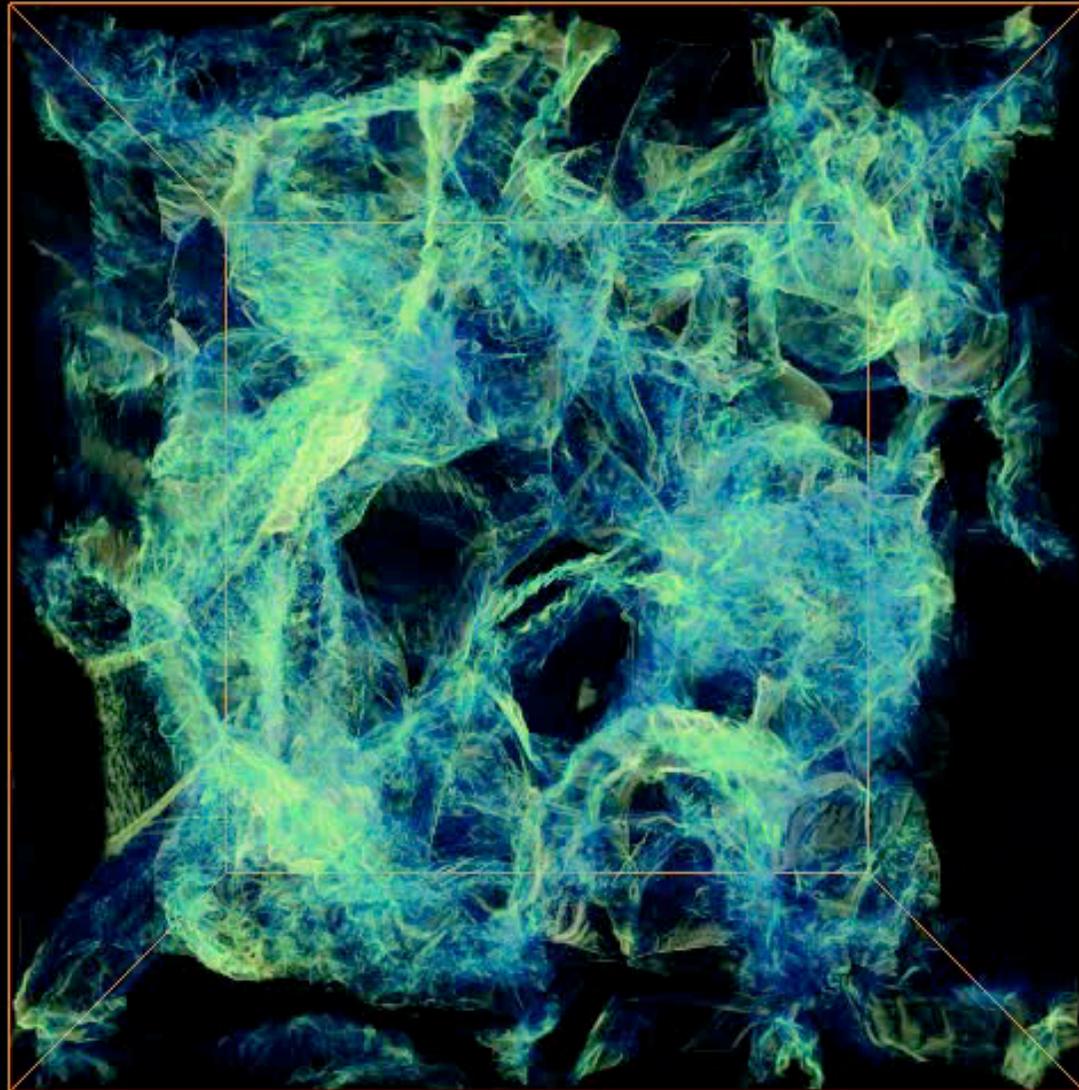


image from Alyssa Goodman: COMPLETE survey



Schmidt et al. (2009, A&A, 494, 127)



# example: model of Orion cloud

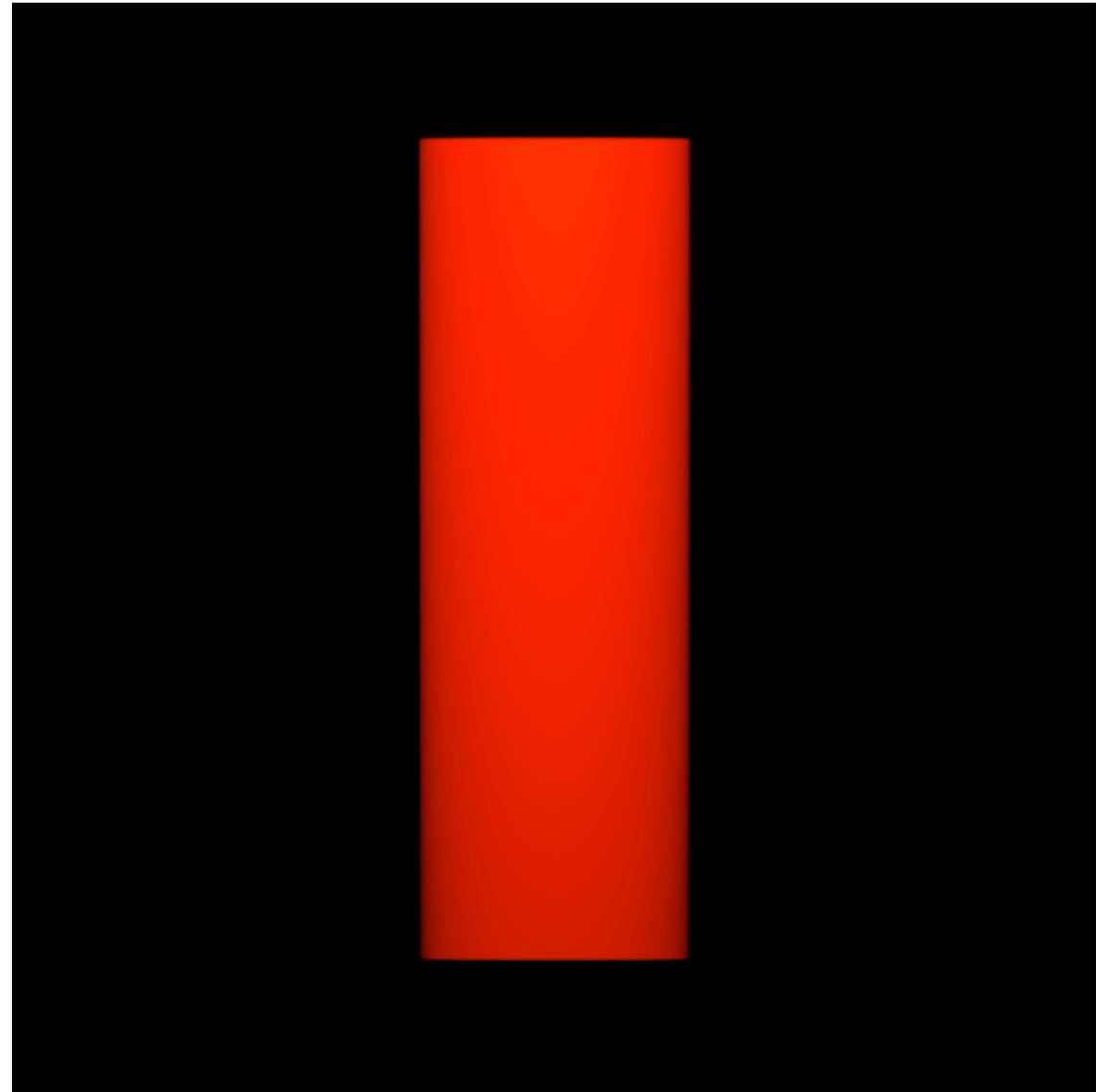
„model“ of Orion cloud:  
15.000.000 SPH particles,  
 $10^4 M_{\text{sun}}$  in 10 pc, mass resolution  
 $0,02 M_{\text{sun}}$ , forms  $\sim 2.500$   
„stars“ (sink particles)

isothermal EOS, top bound, bottom  
unbound

has clustered as well as distributed  
„star“ formation

efficiency varies from 1% to 20%

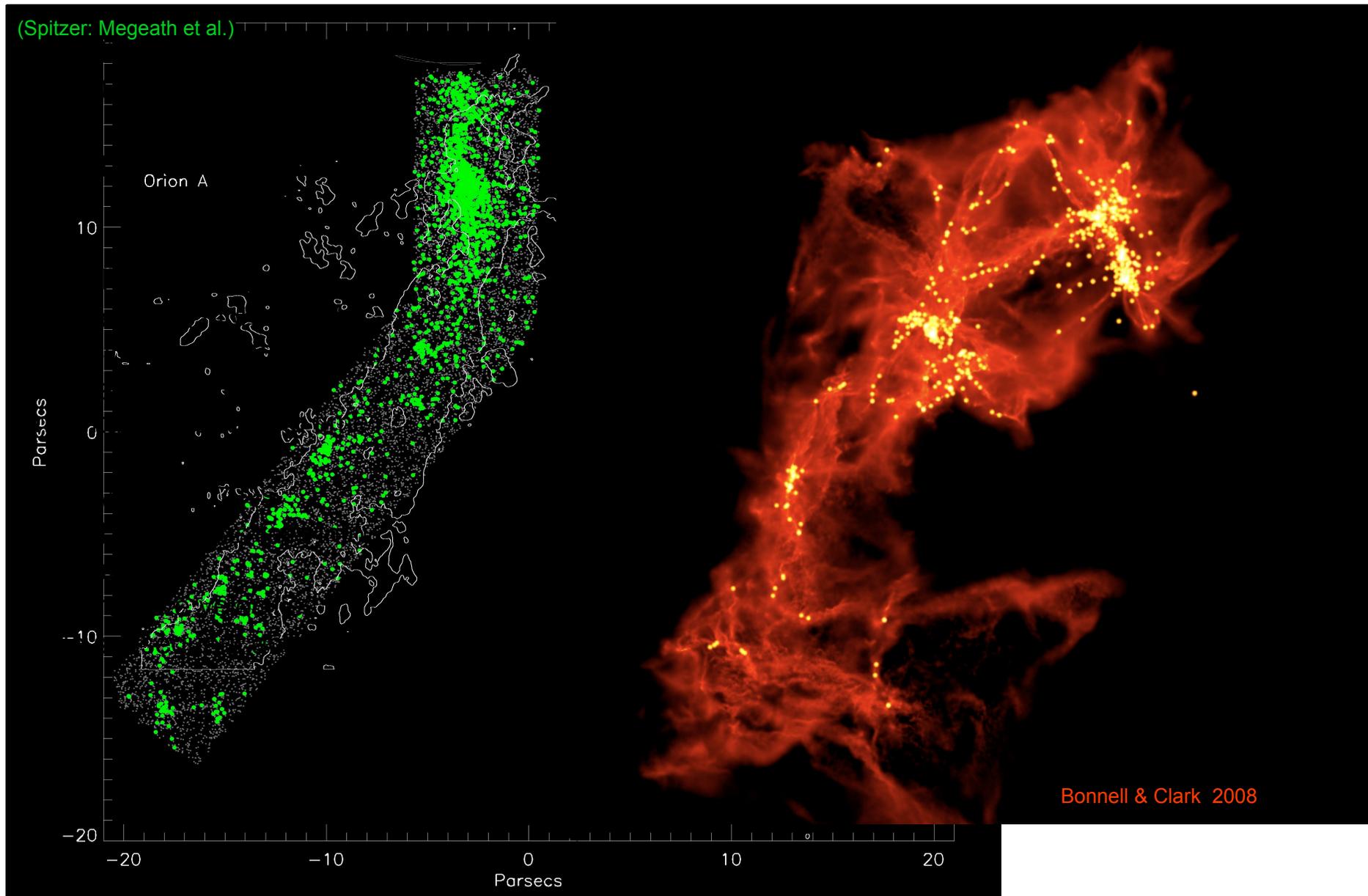
develops full IMF  
(distribution of sink particle masses)



(Bonnell & Clark 2008)



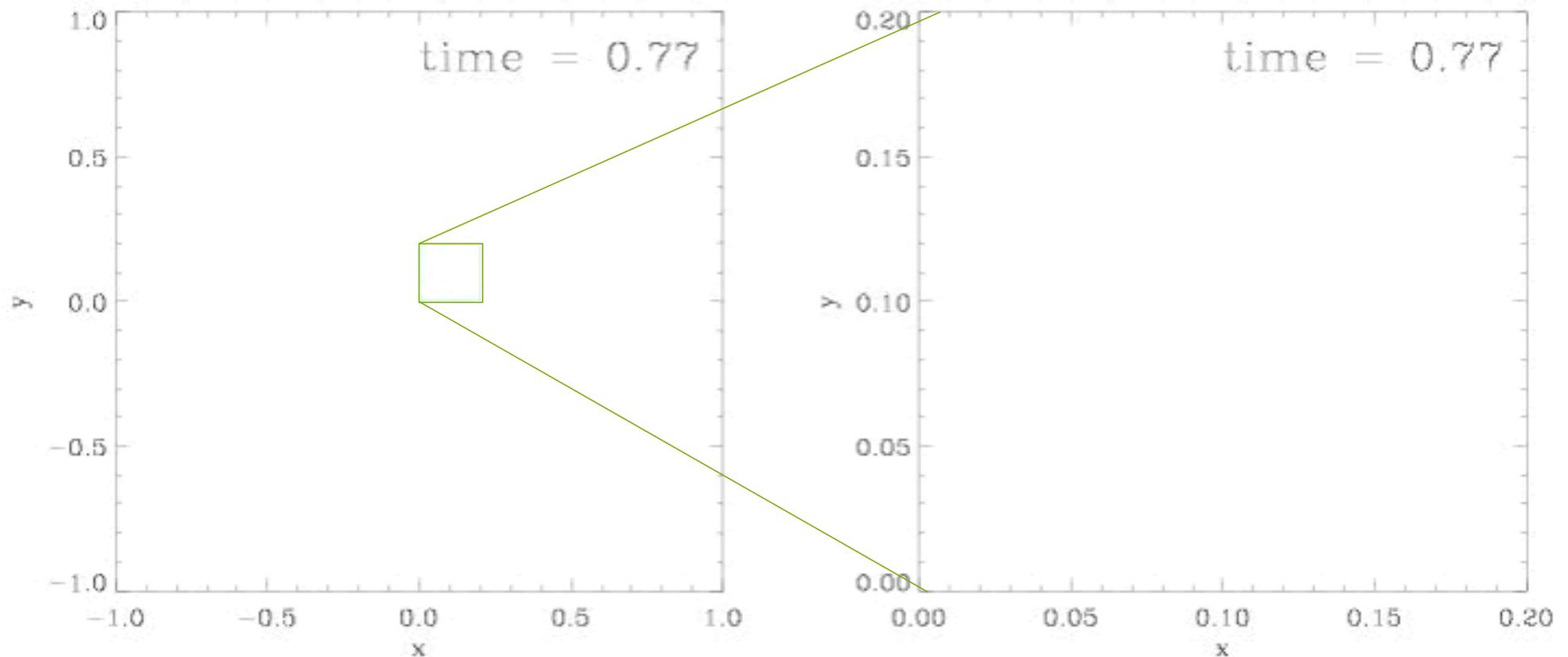
# example: model of Orion cloud





# dynamics of nascent star cluster

in dense clusters protostellar interaction may become important!

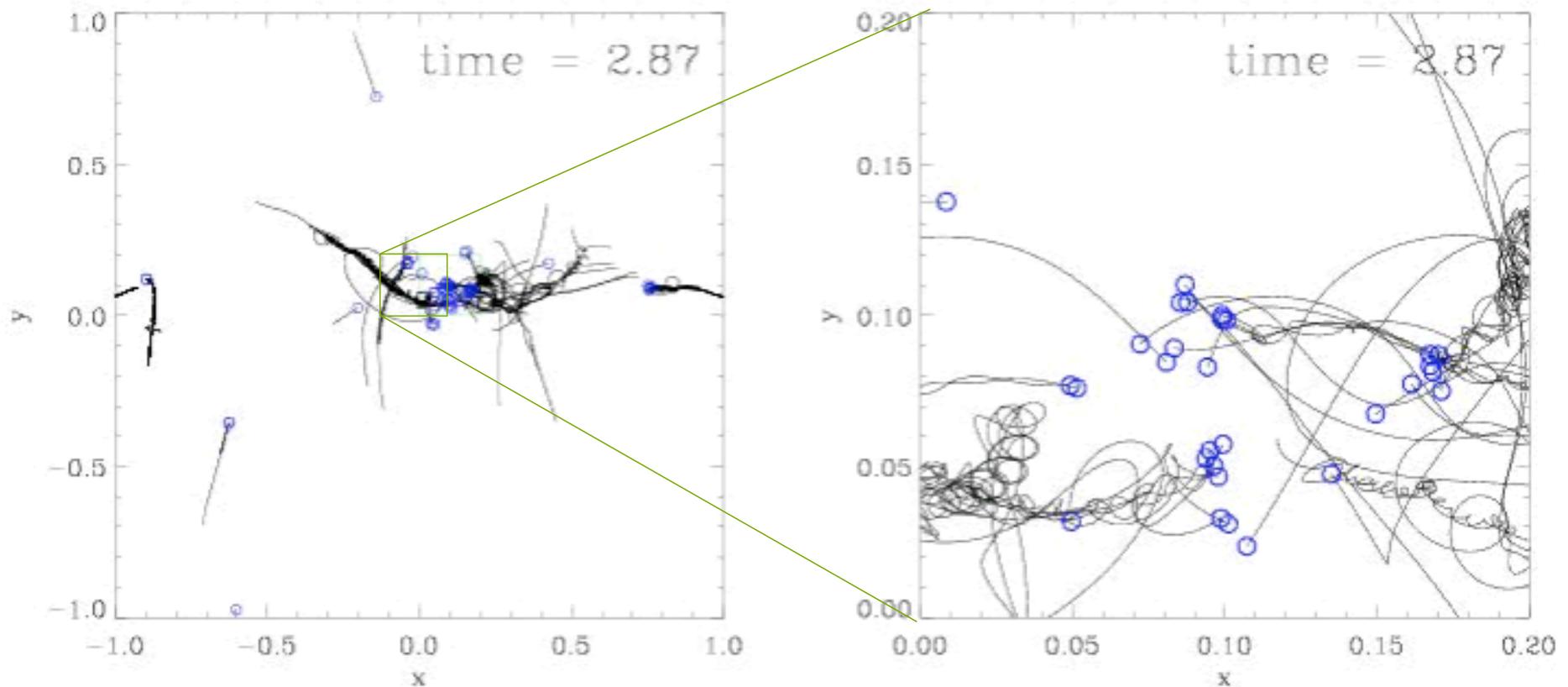


Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation  
(from Klessen & Burkert 2000, ApJS, 128, 287)



# Dynamics of nascent star cluster

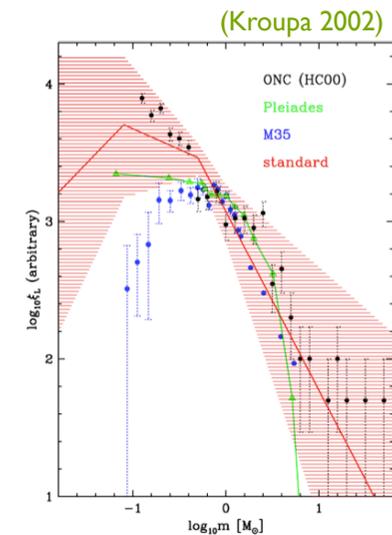
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Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation  
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# stellar masses

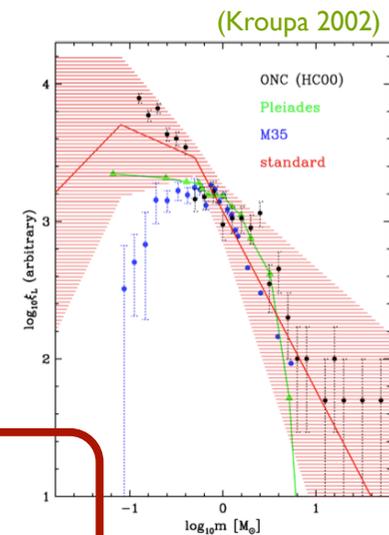
- distribution of stellar masses depends on
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# stellar masses

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application to first star formation



# thermodynamics & fragmentation

degree of fragmentation depends on *EOS!*

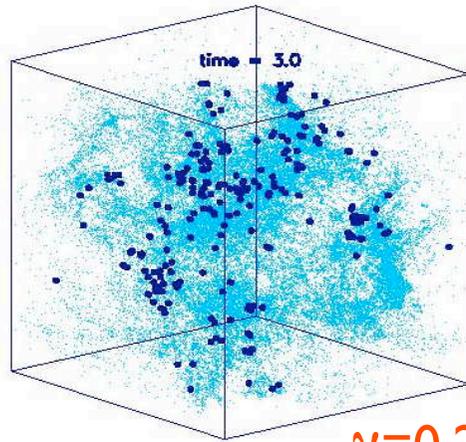
polytropic EOS:  $p \propto \rho^\gamma$

$\gamma < 1$ : dense cluster of low-mass stars

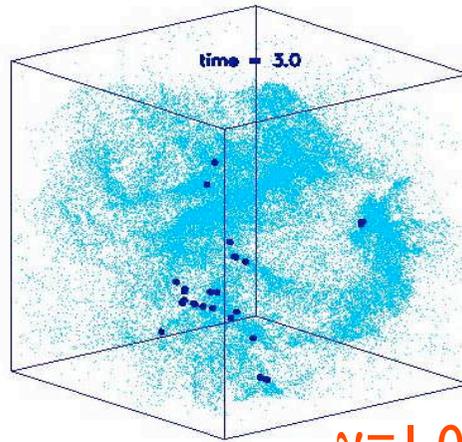
$\gamma > 1$ : isolated high-mass stars

(see Li et al. 2003; also Kawachi & Hanawa 1998, Larson 2003)

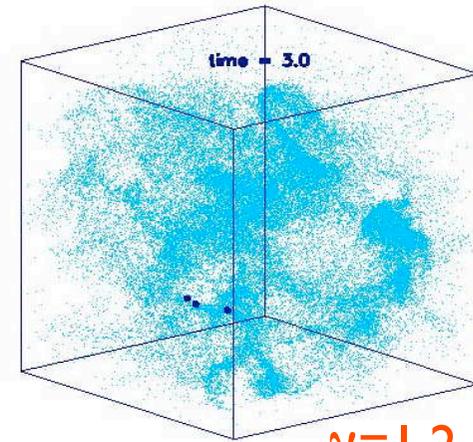
# dependency on EOS



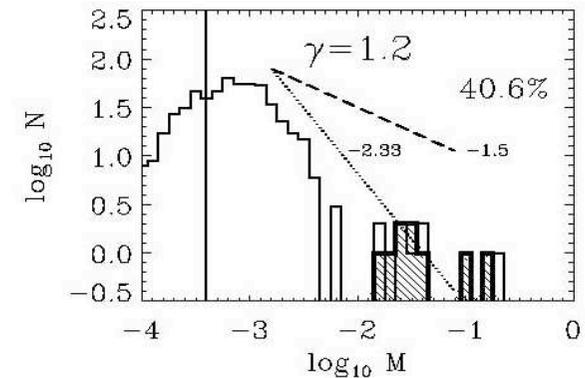
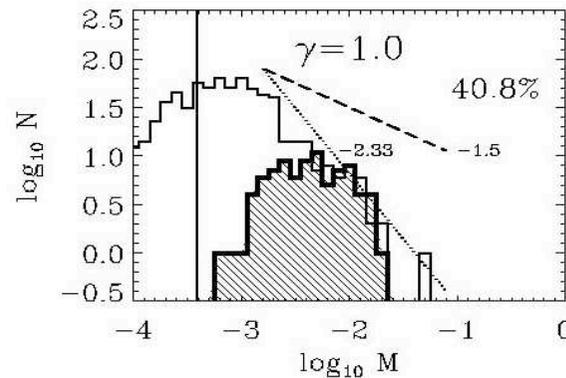
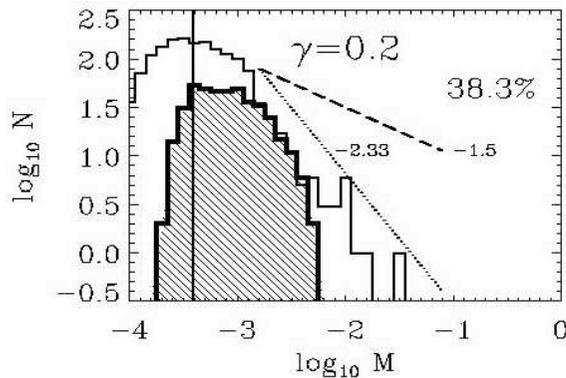
$\gamma = 0.2$



$\gamma = 1.0$



$\gamma = 1.2$

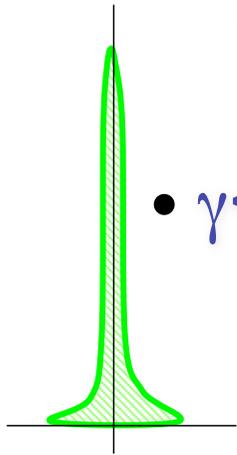


for  $\gamma < 1$  fragmentation is enhanced  $\rightarrow$  *cluster of low-mass stars*  
for  $\gamma > 1$  it is suppressed  $\rightarrow$  formation of *isolated massive stars*

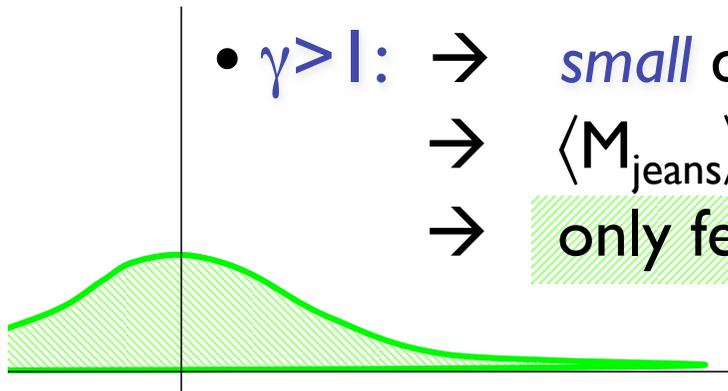
# how does that work?

$$(1) \mathbf{p} \propto \rho^\gamma \rightarrow \rho \propto \mathbf{p}^{1/\gamma}$$

$$(2) \mathbf{M}_{\text{jeans}} \propto \gamma^{3/2} \rho^{(3\gamma-4)/2}$$

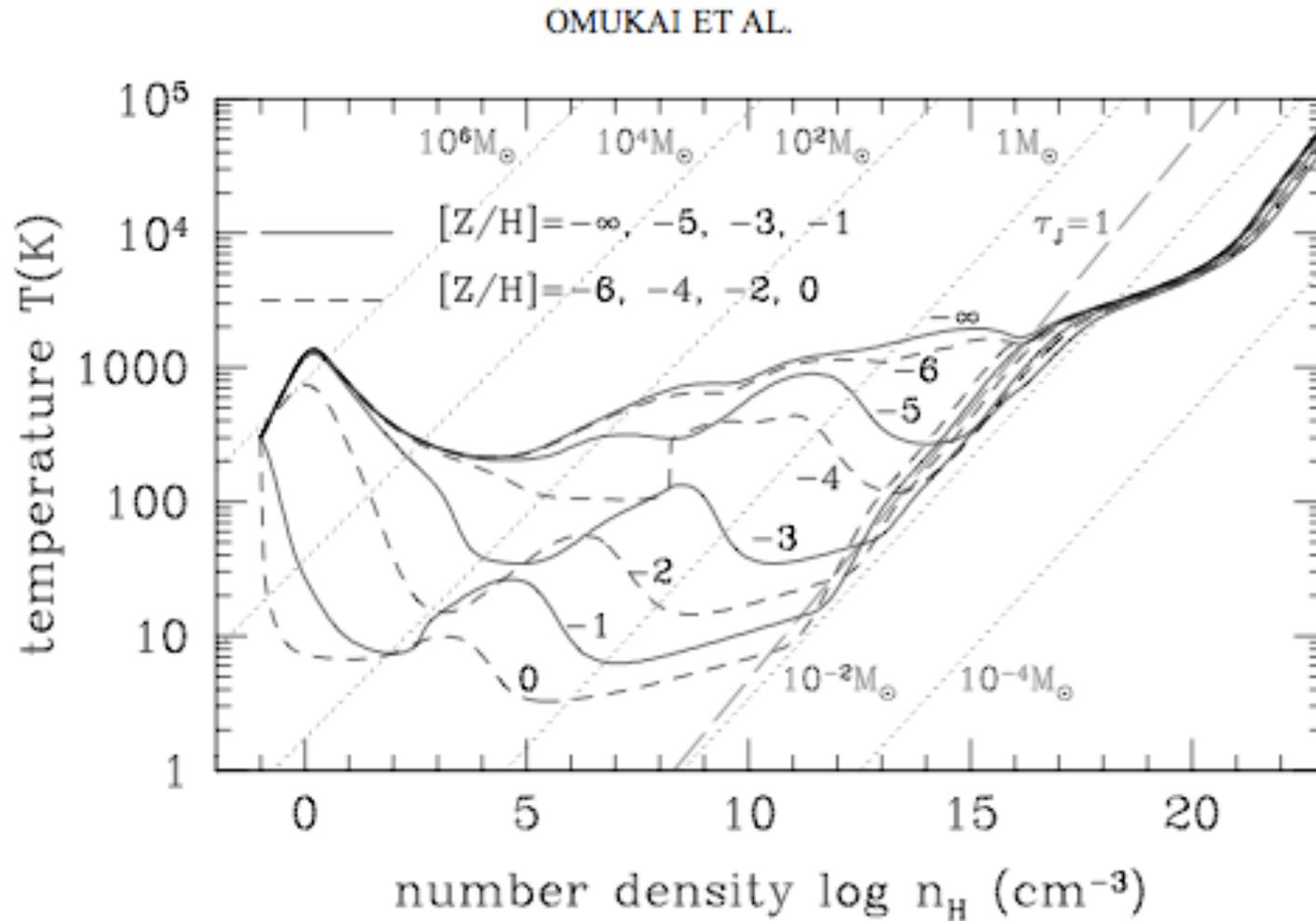


- $\gamma < 1$ :  $\rightarrow$  *large* density excursion for given pressure
  - $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  becomes small
  - $\rightarrow$  number of fluctuations with  $M > M_{\text{jeans}}$  is large



- $\gamma > 1$ :  $\rightarrow$  *small* density excursion for given pressure
  - $\rightarrow$   $\langle M_{\text{jeans}} \rangle$  is large
  - $\rightarrow$  only few and massive clumps exceed  $M_{\text{jeans}}$

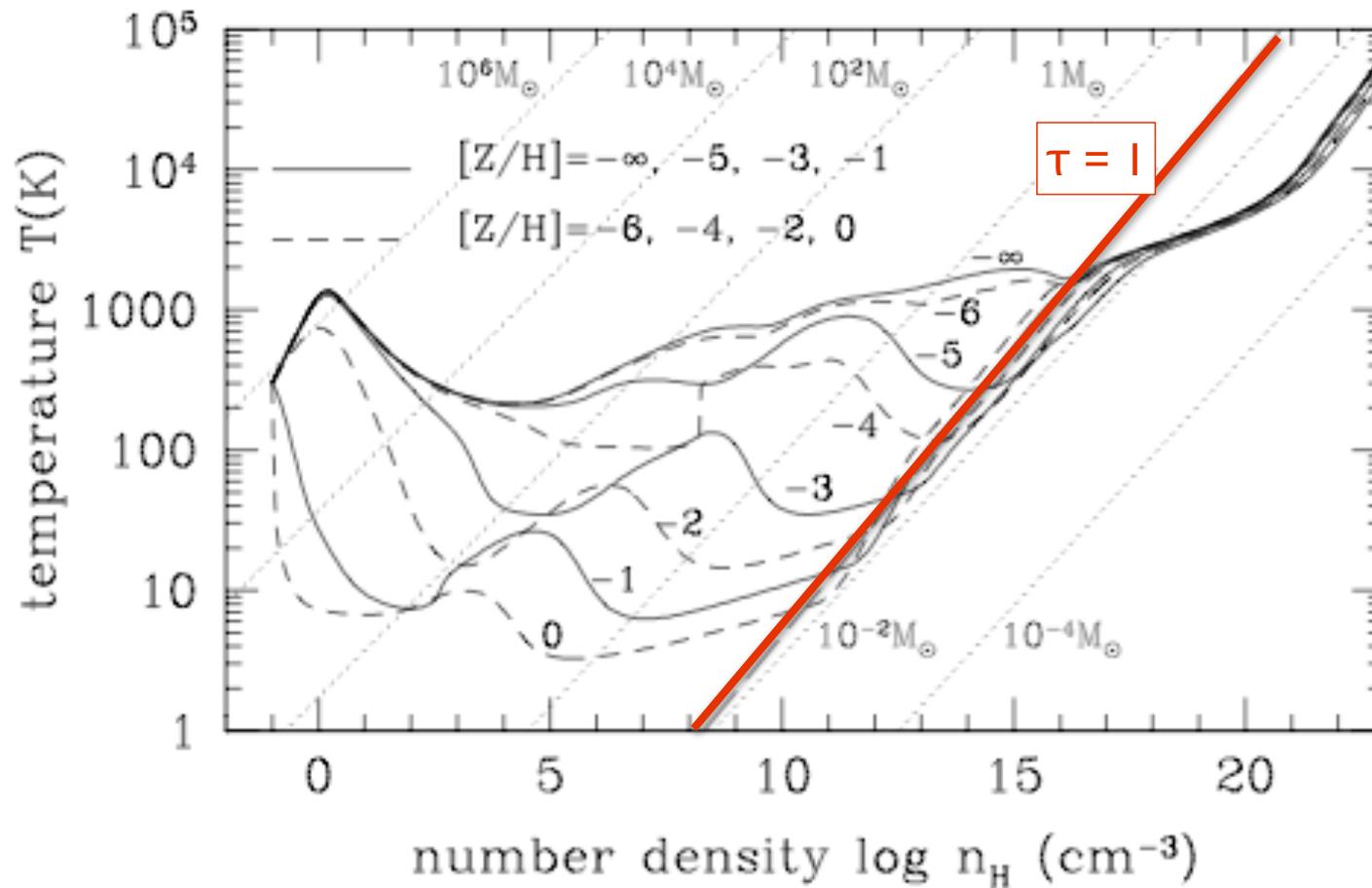
# EOS as function of metallicity



(Omukai et al. 2005)

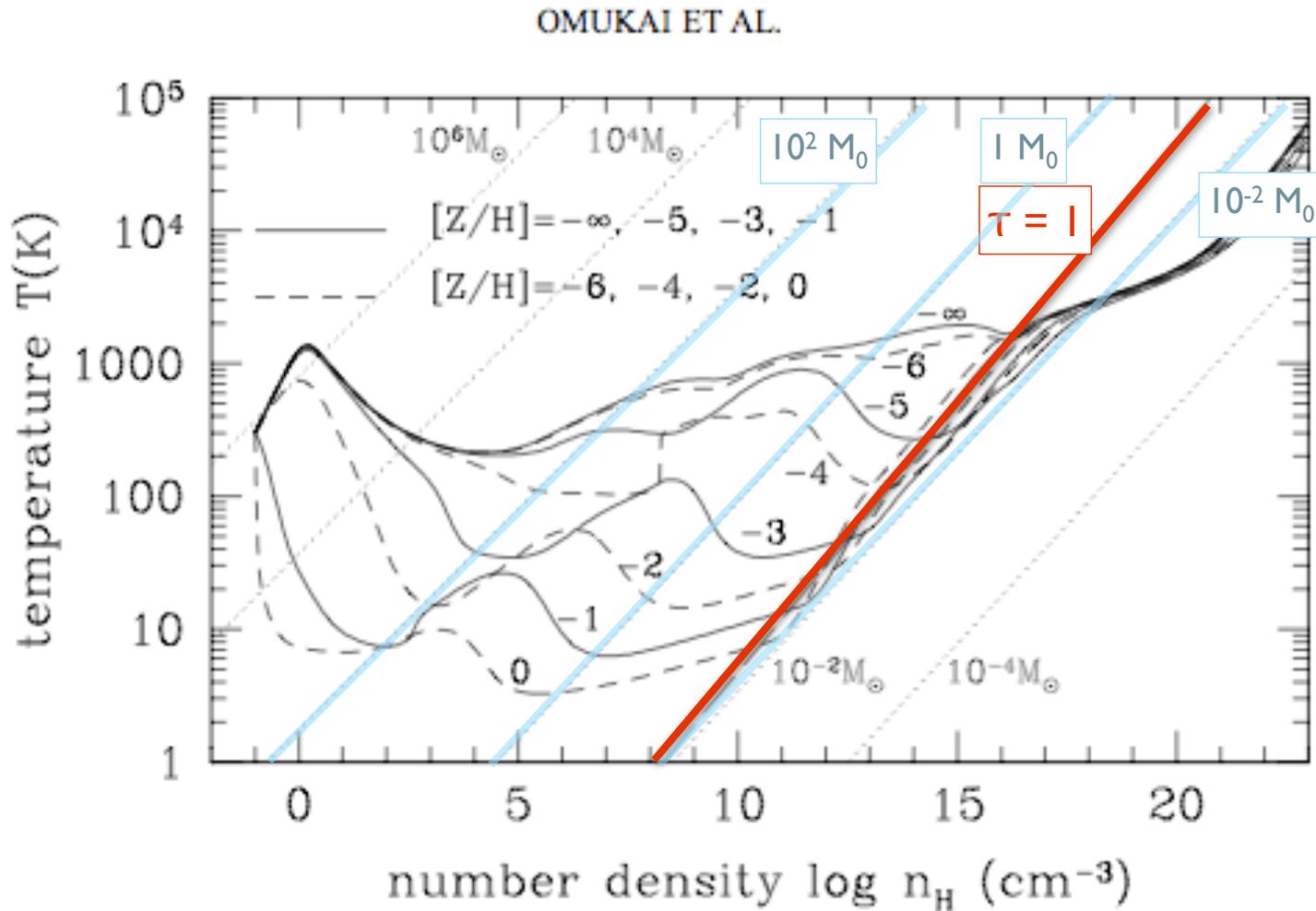
# EOS as function of metallicity

OMUKAI ET AL.



(Omukai et al. 2005)

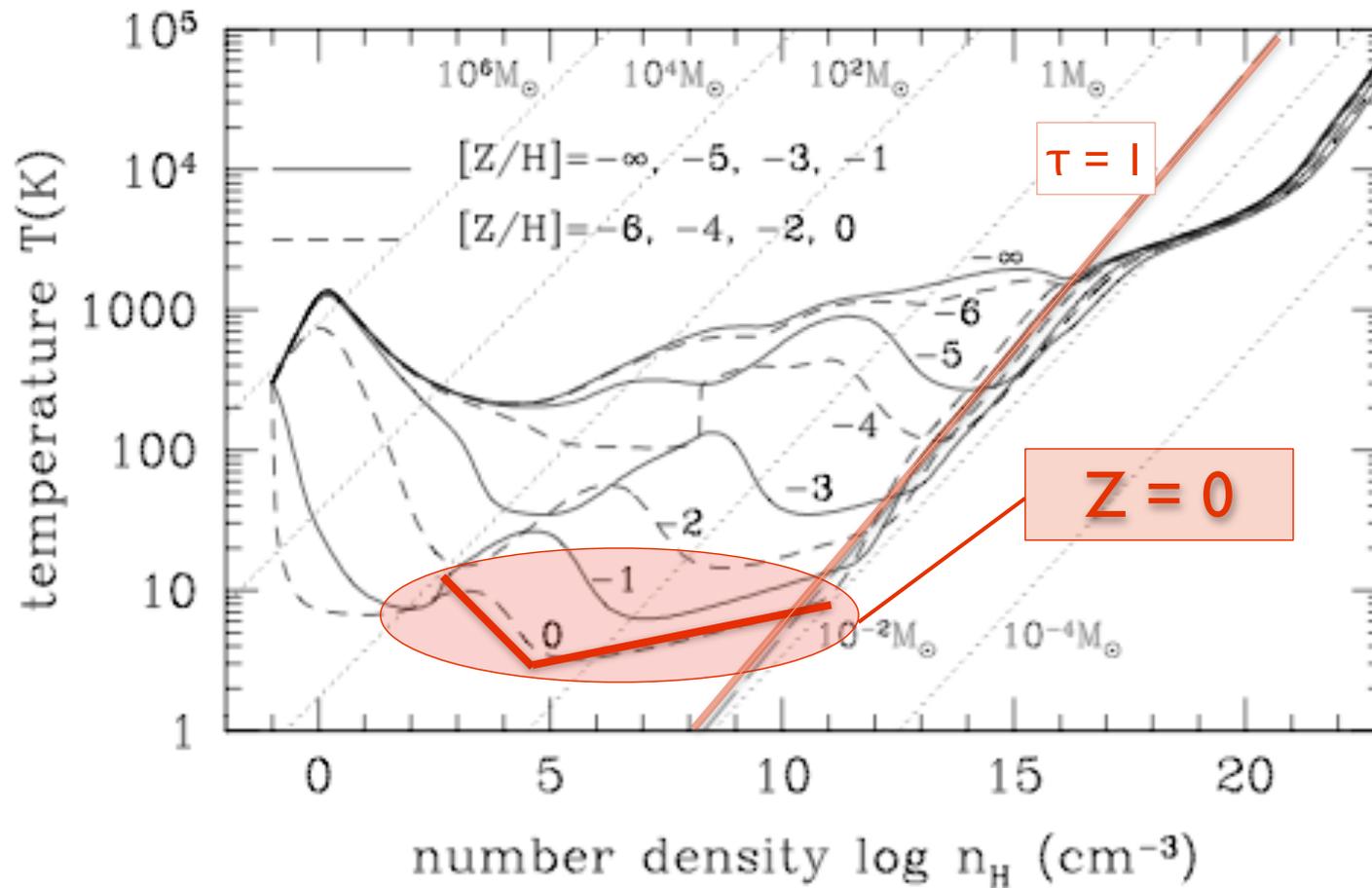
# EOS as function of metallicity



(Omukai et al. 2005)

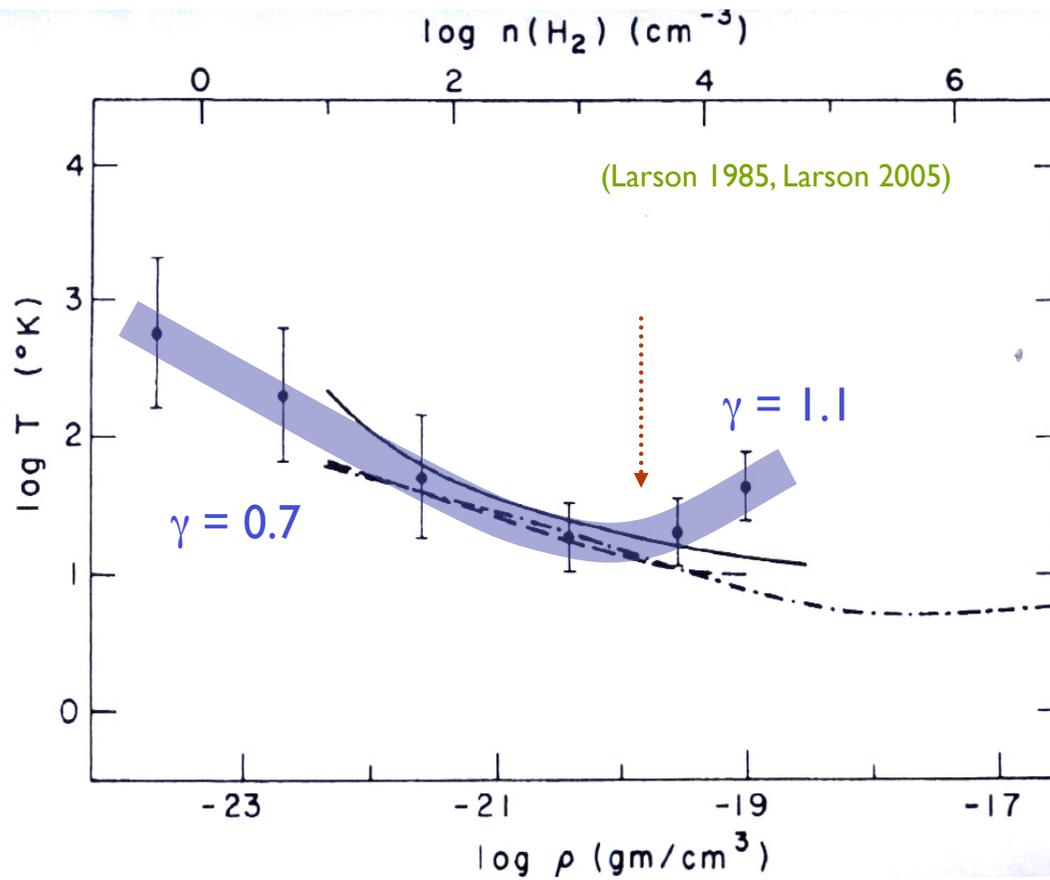
# present-day star formation

OMUKAI ET AL.



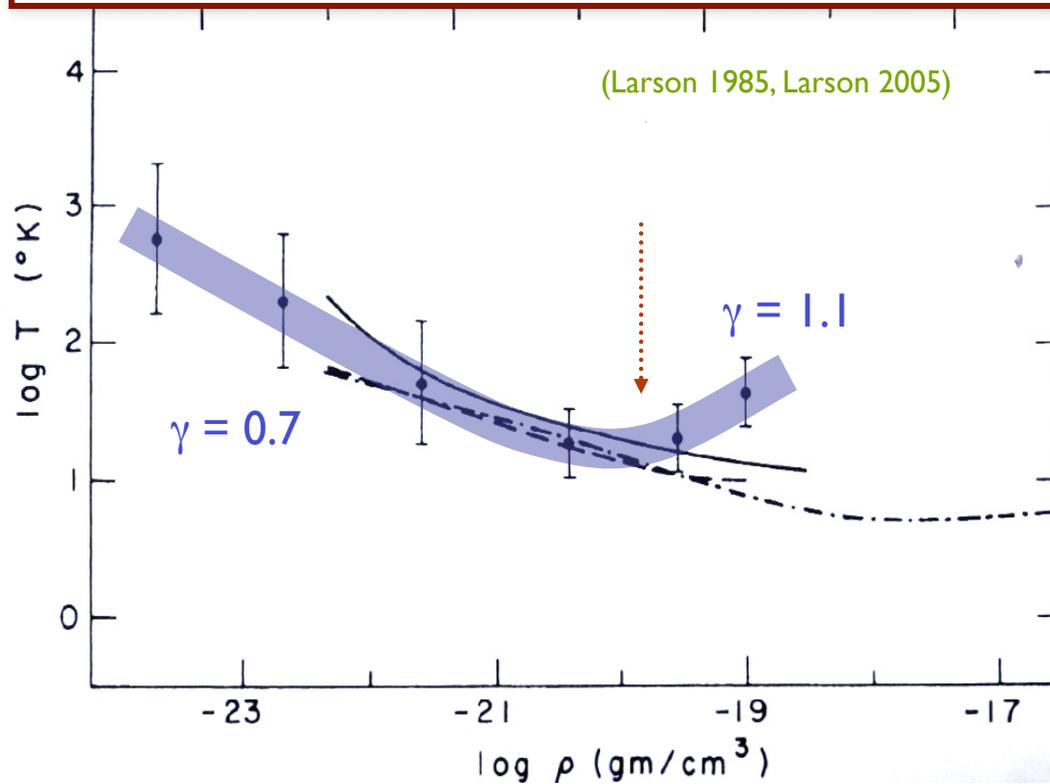
(Omukai et al. 2005, Jappsen et al. 2005, Larson 2005)

# present-day star formation

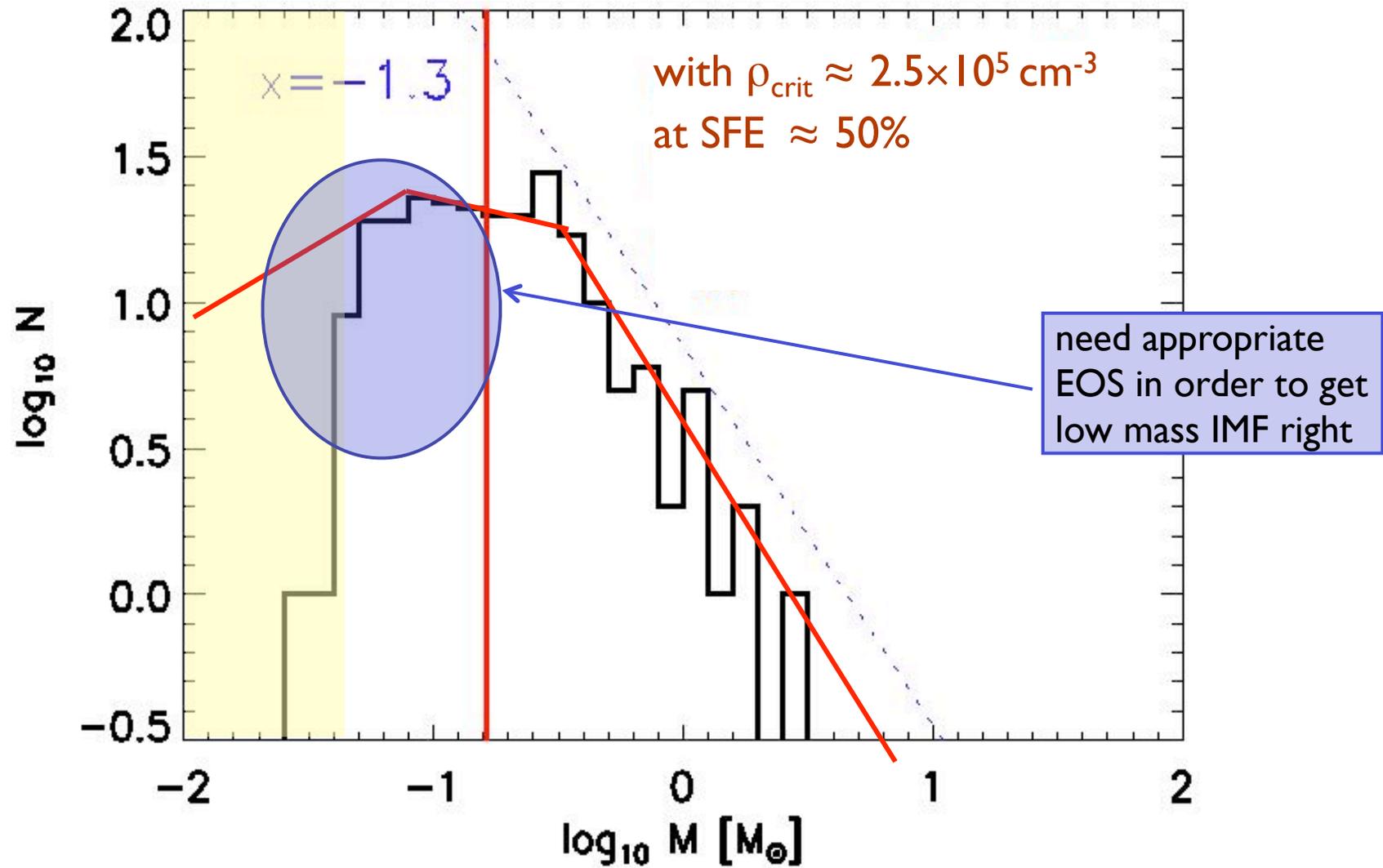


# present-day star formation

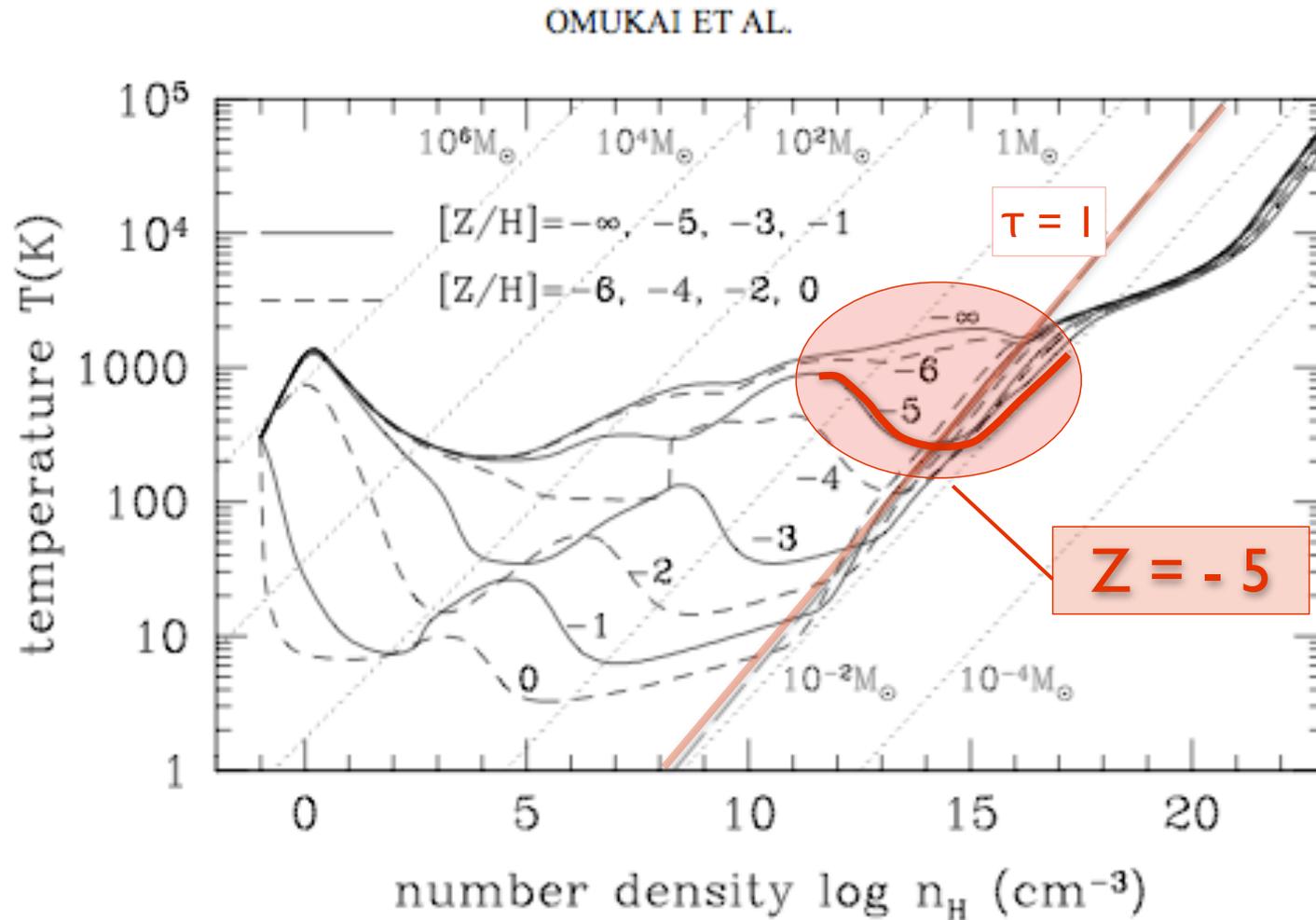
This kink in EOS is very insensitive to environmental conditions such as ambient radiation field  
--> reason for universal form of the IMF? (Elmegreen et al. 2008)



# IMF in nearby molecular clouds



# transition: Pop III to Pop II.5



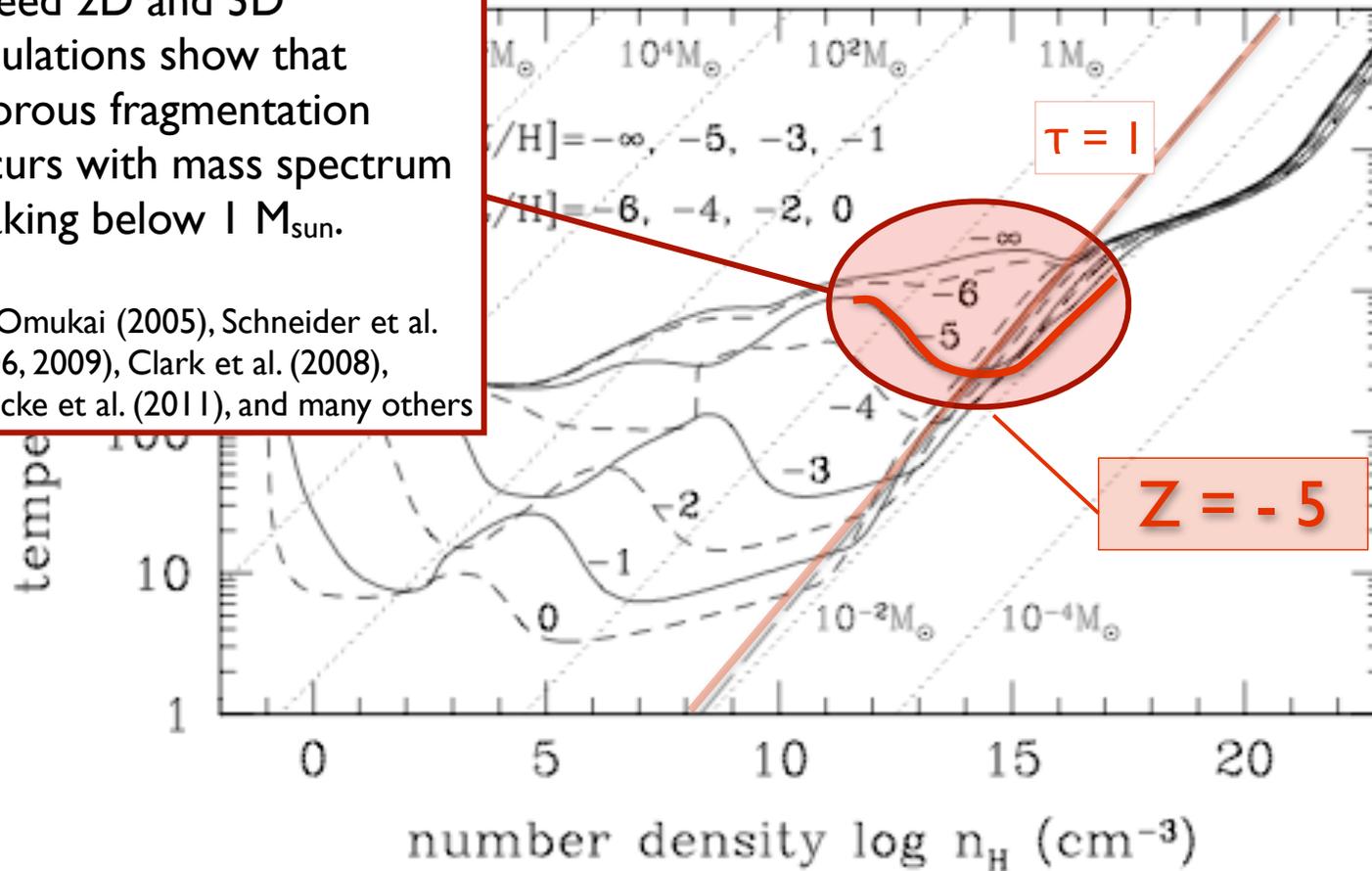
(Omukai et al. 2005)

# transition: Pop III to Pop II.5

OMUKAI ET AL.

indeed 2D and 3D simulations show that vigorous fragmentation occurs with mass spectrum peaking below  $1 M_{\text{sun}}$ .

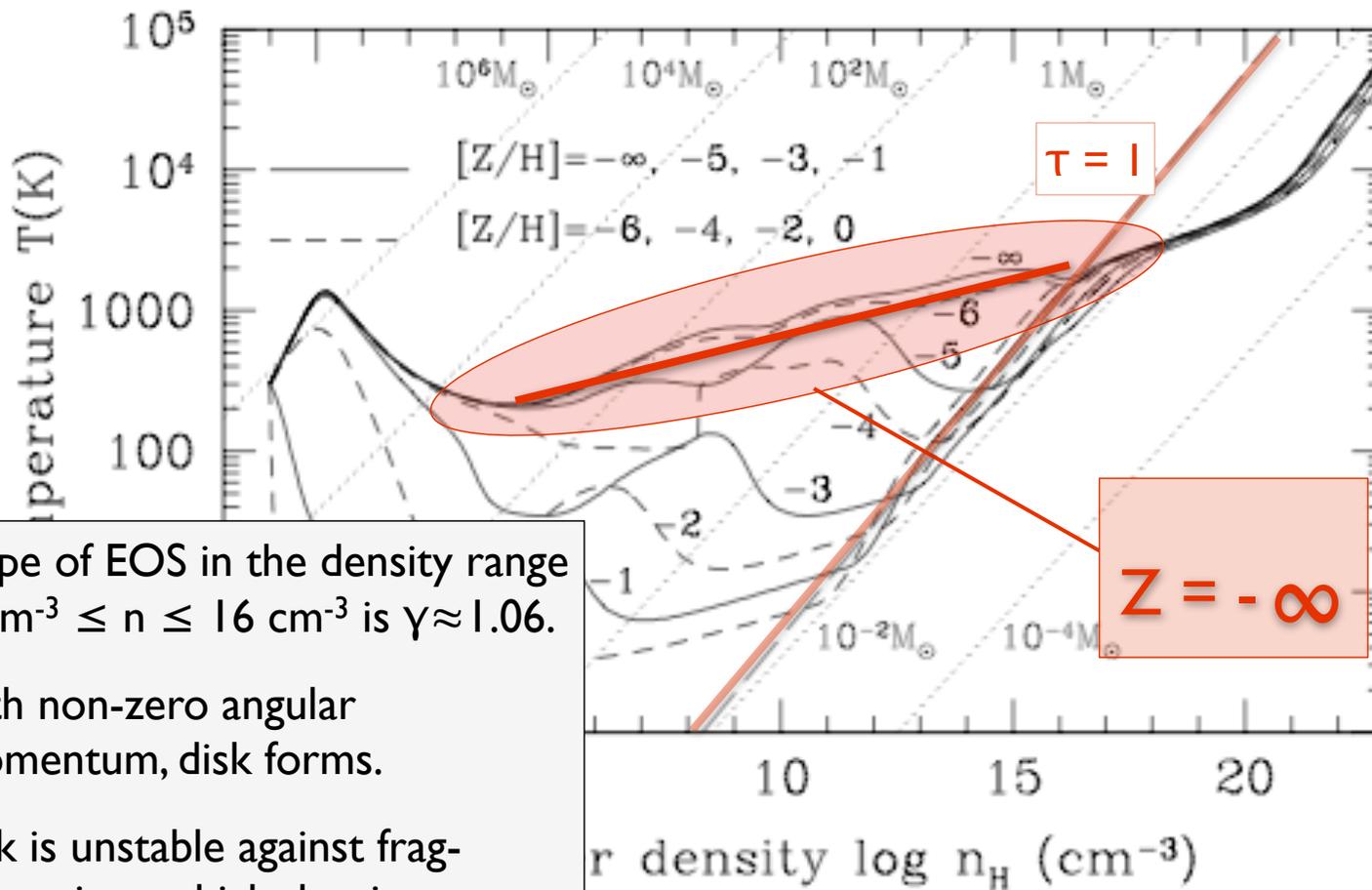
see Omukai (2005), Schneider et al. (2006, 2009), Clark et al. (2008), Dopcke et al. (2011), and many others



(Omukai et al. 2005)

# metal-free star formation

OMUKAI ET AL.

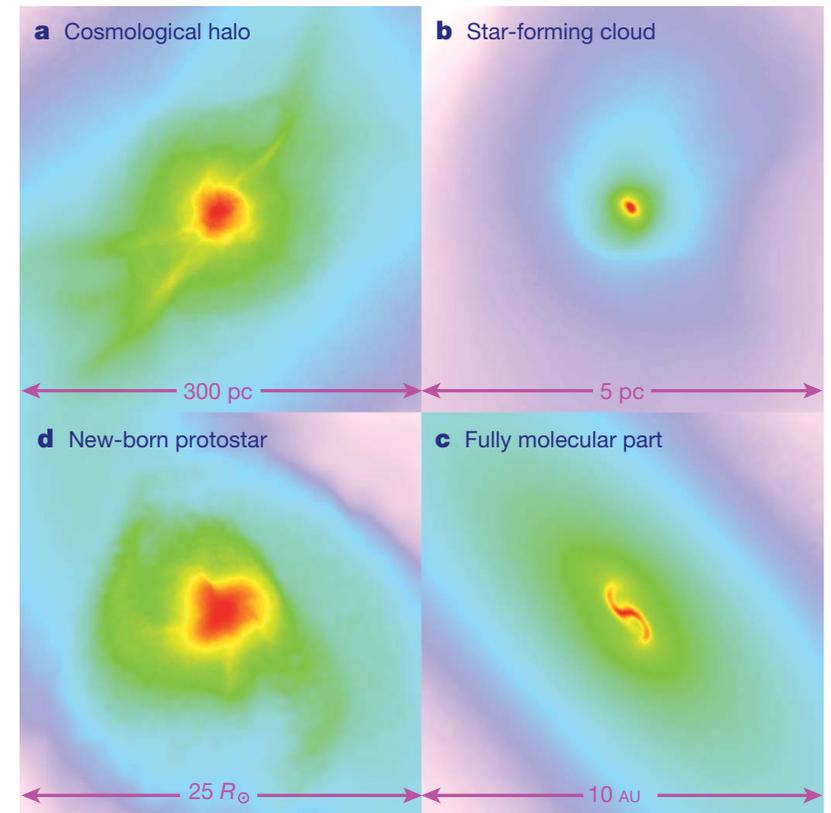


- slope of EOS in the density range  $5 \text{ cm}^{-3} \leq n \leq 16 \text{ cm}^{-3}$  is  $\gamma \approx 1.06$ .
- with non-zero angular momentum, disk forms.
- disk is unstable against fragmentation at high density

(Omukai et al. 2005)

# metal-free star formation

- most current numerical simulations of Pop III star formation predict very massive objects (e.g. Abel et al. 2002, Yoshida et al. 2008, Bromm et al. 2009)
- similar for theoretical models (e.g. Tan & McKee 2004)
- there are some first hints of fragmentation, however (Turk et al. 2009, Stacy et al. 2010)

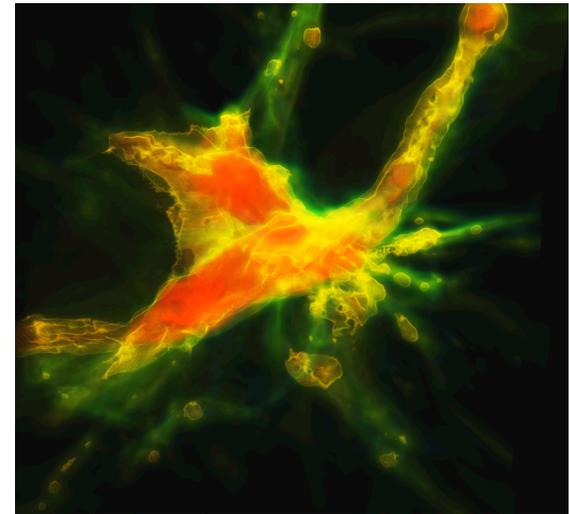


**Figure 1 | Projected gas distribution around a primordial protostar.** Shown is the gas density (colour-coded so that red denotes highest density) of a single object on different spatial scales. **a**, The large-scale gas distribution around the cosmological minihalo; **b**, a self-gravitating, star-forming cloud; **c**, the central part of the fully molecular core; and **d**, the final protostar. Reproduced by permission of the AAAS (from ref. 20).

(Yoshida et al. 2008, *Science*, 321, 669)

# turbulence in Pop III halos

- star formation will depend on *degree of turbulence* in protogalactic halo
- speculation: *differences in stellar mass function*, just like in present-day star formation

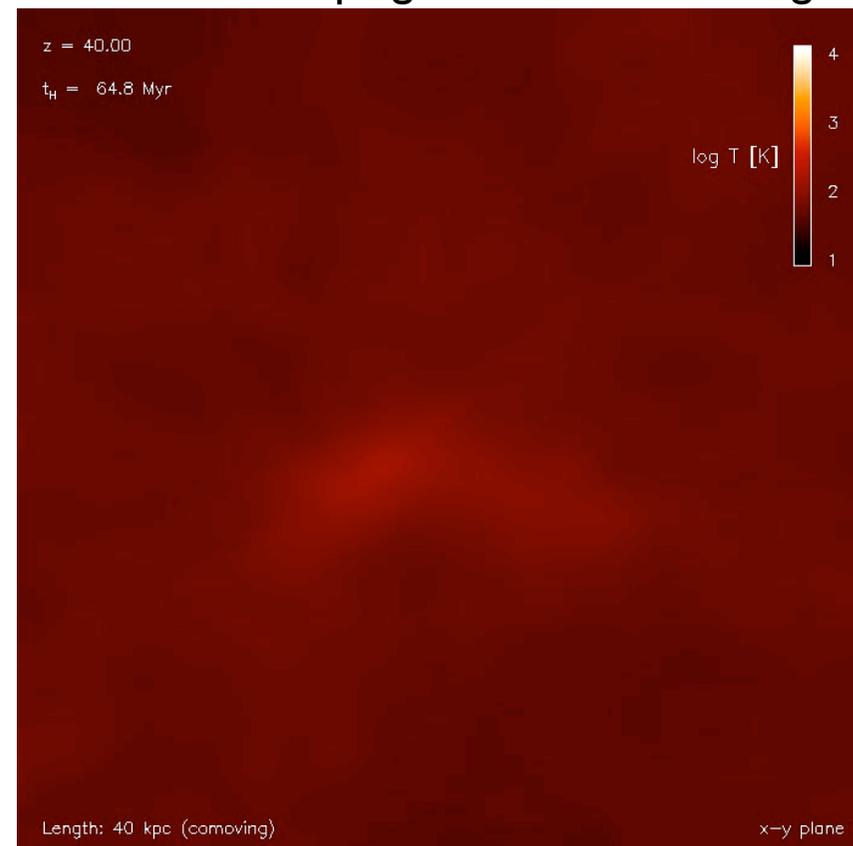


(Greif et al. 2008)

# turbulence in Pop III halos

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- speculation: *differences in stellar mass function*, just like in present-day star formation

turbulence developing in an atomic cooling halo



(Greif et al. 2008)

# multiple Pop III stars in halo

- parameter study with different strength of turbulence using SPH: study Pop III.1 and Pop III.2 case (Clark et al., 2011a, ApJ, 727, 110)
- 2 very high resolution studies of Pop III star formation in cosmological context
  - SPH: Clark et al. 2011b, Science, 311, 1040
  - Arepo: Greif et al. 2011a, ApJ, in press (arXiv:1101.5491)
  - complementary approaches with interesting similarities and differences....

# SPH study: face on look at accretion disk

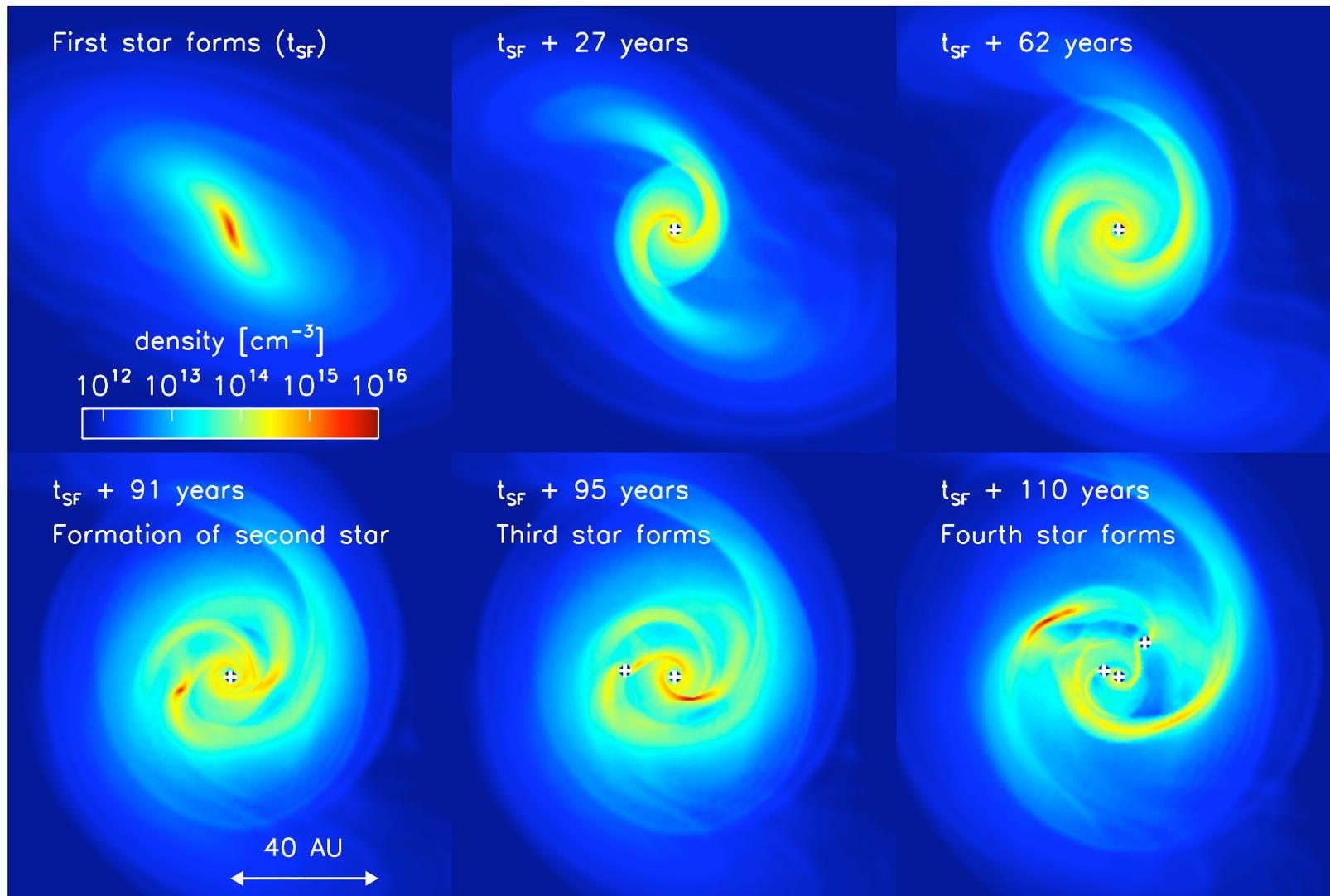


Figure 1: Density evolution in a 120 AU region around the first protostar, showing the build-up of the protostellar disk and its eventual fragmentation. We also see ‘wakes’ in the low-density regions, produced by the previous passage of the spiral arms.

# SPH study: some disk parameters

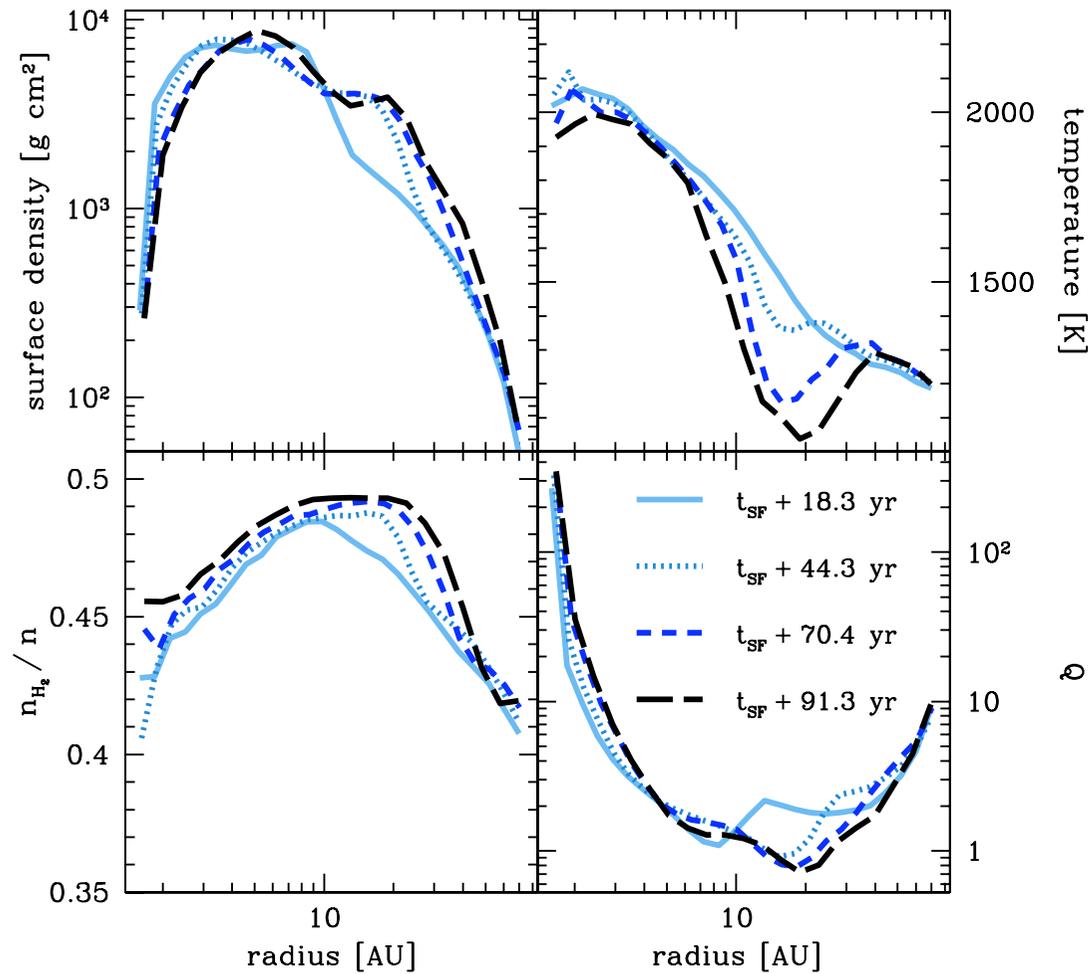


Figure 2: Radial profiles of the disk's physical properties, centered on the first protostellar core to form. The quantities are mass-weighted and taken from a slice through the midplane of the disk. In the lower right-hand plot we show the radial distribution of the disk's Toomre parameter,  $Q = c_s \kappa / \pi G \Sigma$ , where  $c_s$  is the sound speed and  $\kappa$  is the epicyclic frequency. Because our disk is Keplerian, we adopted the standard simplification, and replaced  $\kappa$  with the orbital frequency. The molecular fraction is defined as the number density of hydrogen molecules ( $n_{\text{H}_2}$ ), divided by the number density of hydrogen nuclei ( $n$ ), such that fully molecular gas has a value of 0.5

# SPH study: mass accretion onto disk and onto protostars

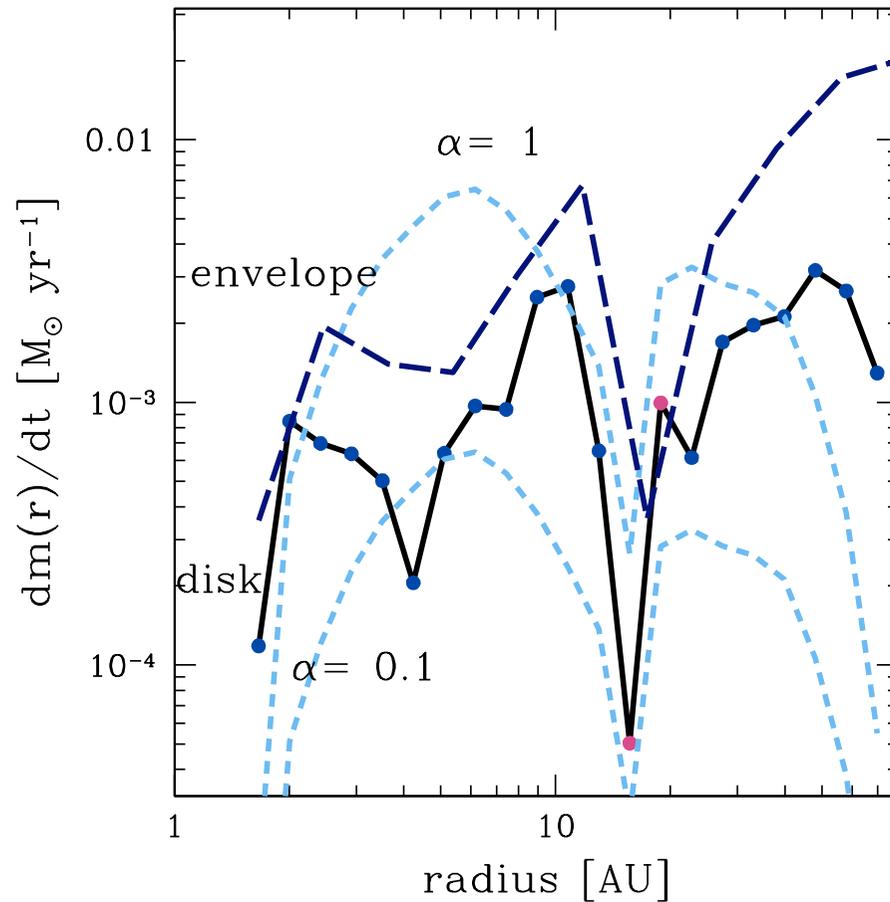


Figure 3: The mass transfer rate through the disk is denoted by the solid black line, while the mass infall rate through spherical shells with the specified radius is shown by the dark blue dashed line. The latter represents the total amount of material flowing through a given radius, and is thus a measure of the material flowing through *and onto* the disk at each radius. Both are shown at the onset of disk fragmentation. In the case of the disk accretion we have denoted annuli that are moving towards the protostar with blue dots, and those moving away in pink (further details can be found in Section 6 of the online material). The light blue dashed lines show the accretion rates expected from an ‘alpha’ (thin) disk model, where  $\dot{M}(r) = 3\pi\alpha c_s(r)\Sigma(r)H(r)$ , with two global values of alpha and where  $c_s(r)$ ,  $\Sigma(r)$ , and  $H(r)$  are (respectively) the sound speed, surface density and disk thickness at radius  $r$ .

# SPH study: comparison of all relevant heating and cooling processes

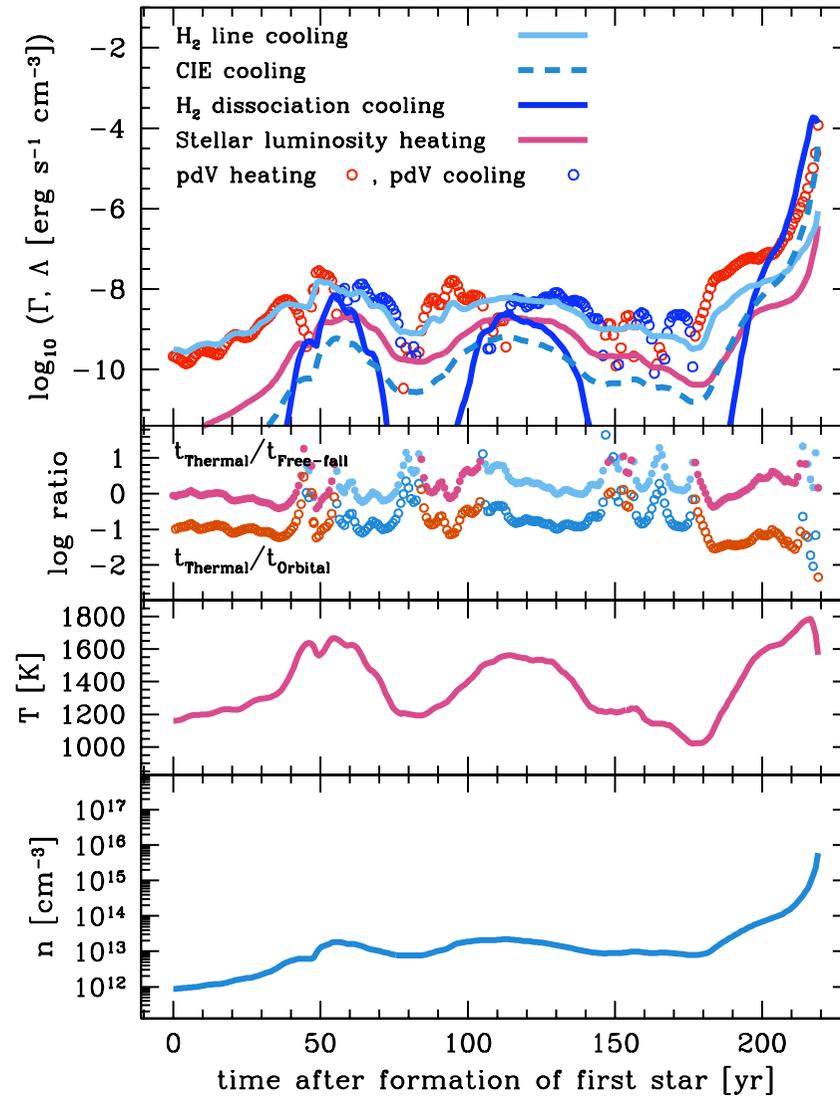
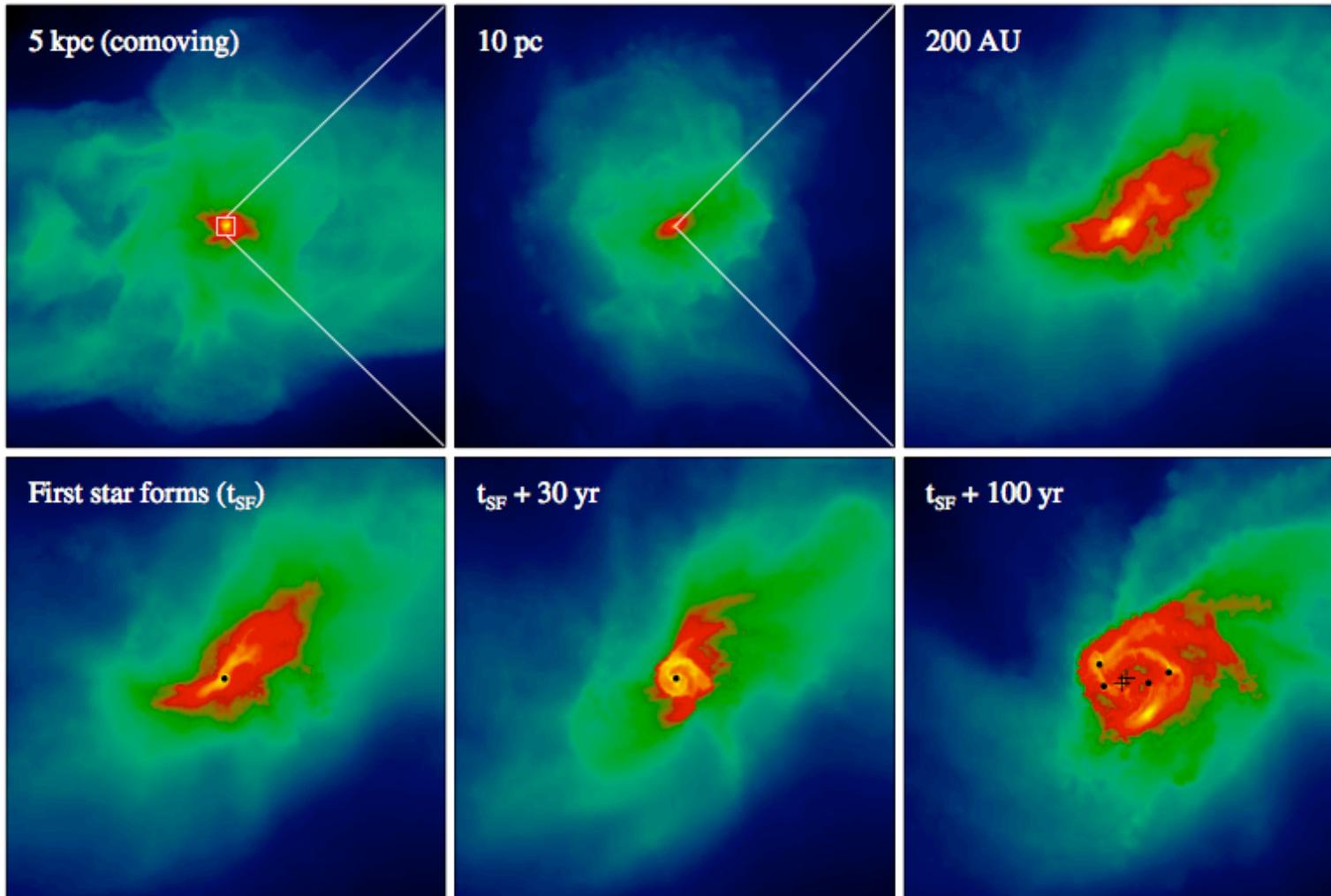
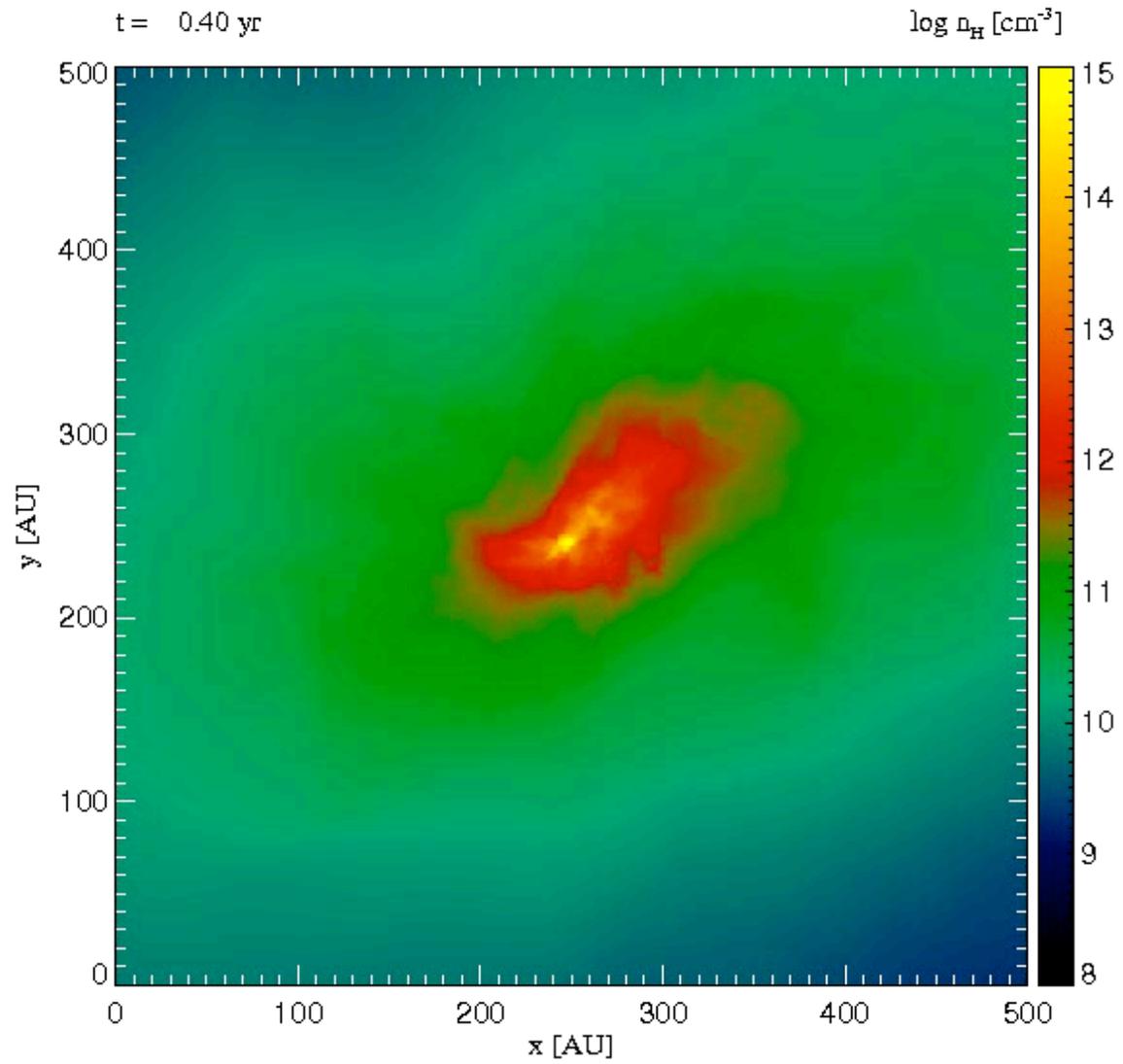


Figure 7: (a) Dominant heating and cooling processes in the gas that forms the second sink particle. (b) Upper line: ratio of the thermal timescale,  $t_{\text{thermal}}$ , to the free-fall timescale,  $t_{\text{ff}}$ , for the gas that forms the second sink particle. Periods when the gas is cooling are indicated in blue, while periods when the gas is heating are indicated in red. Lower line: ratio of  $t_{\text{thermal}}$  to the orbital timescale,  $t_{\text{orbital}}$ , for the same set of SPH particles (c) Temperature evolution of the gas that forms the second sink (d) Density evolution of the gas that forms the second sink

## Arepo study: surface density at different times

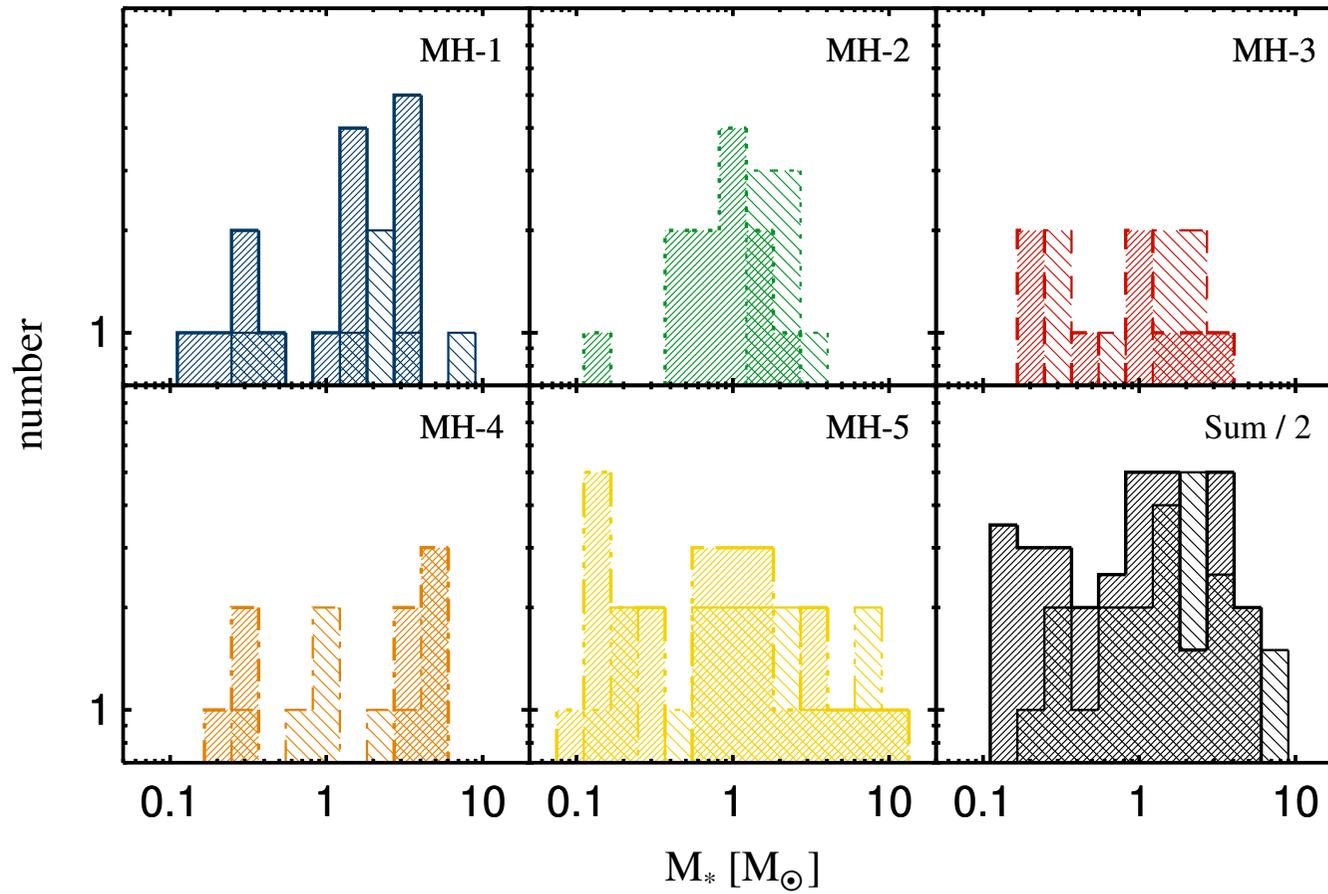


one out of five halos



(Greif et al. 2011a, ApJ, in press, arXiv:1101.5491)

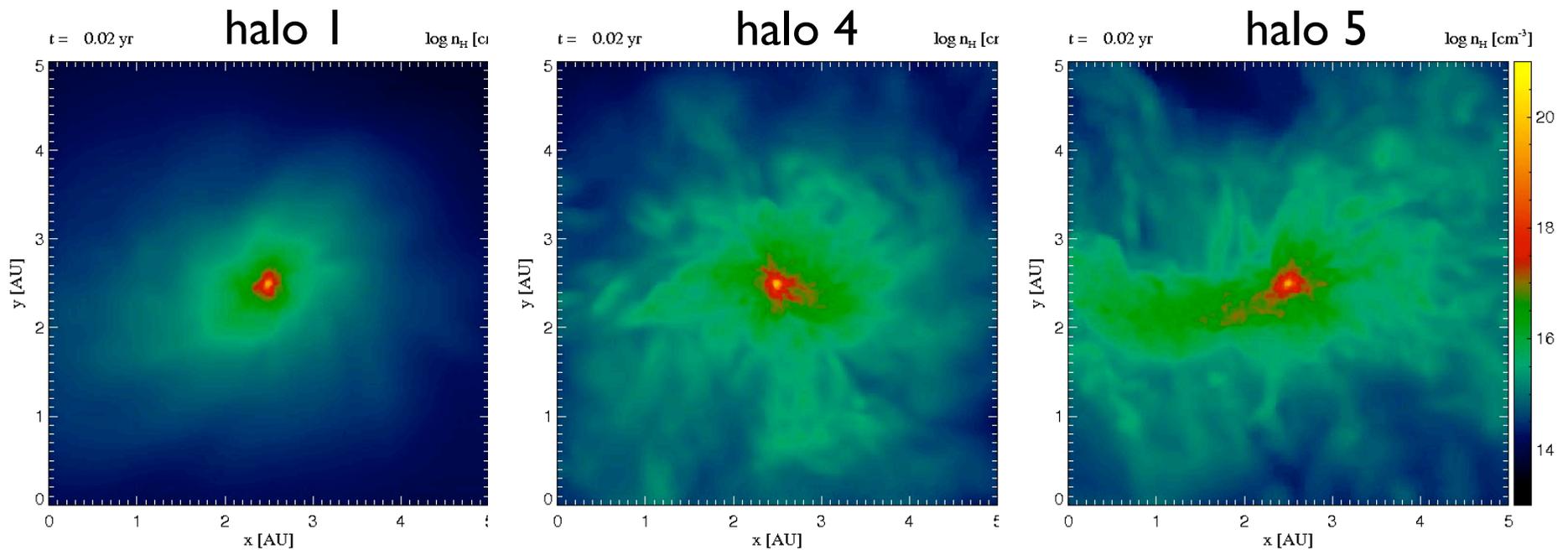
# Arepo study: mass spectrum of fragments



# brand-new “sinkless” calculations

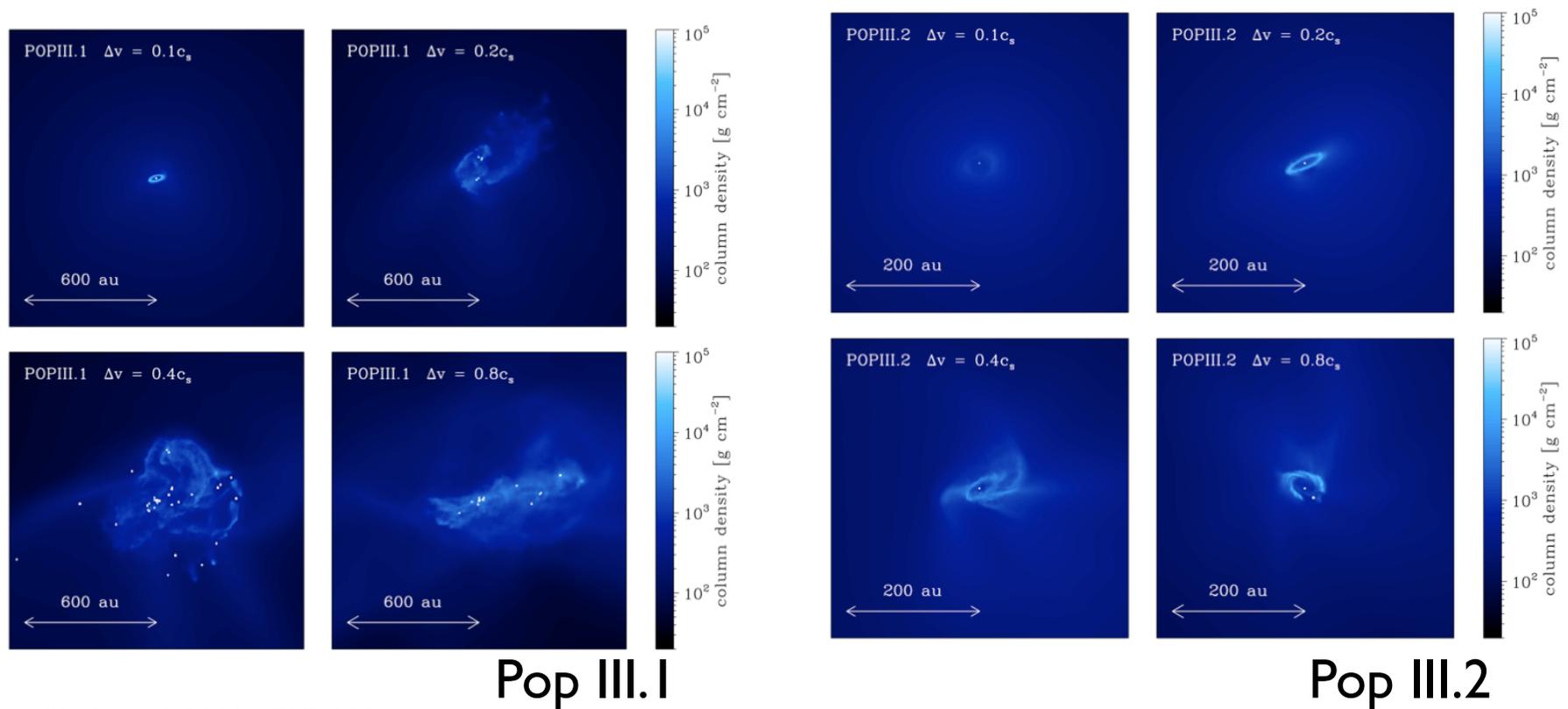
10 years need 1 month on the computer

--> we will never be able to follow full accretion history



(Greif et al. in preparation)

# fragmentation continues to larger scales



(Clark et al, 2011a, 727, 110)

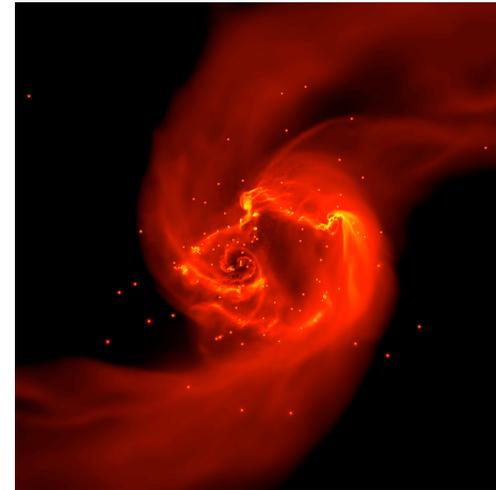
# primordial star formation

- just like in present-day SF, we expect

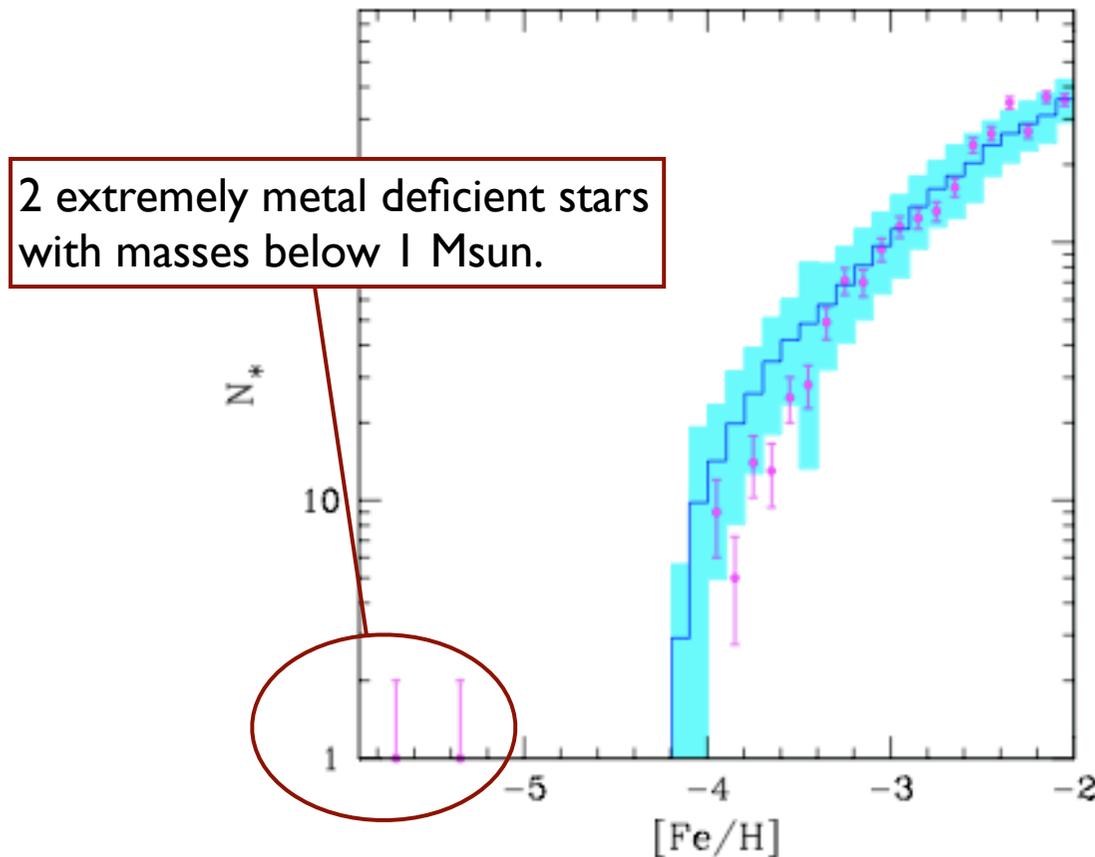
- *turbulence*
- *thermodynamics*
- *feedback*
- *magnetic fields*

to influence Pop III/II star formation.

- masses of Pop III stars still *uncertain* (surprises from new generation of high-resolution calculations that go beyond first collapse)
- disks unstable: Pop III stars should be *binaries* or *part of small clusters*
- effects of feedback less important than in present-day SF



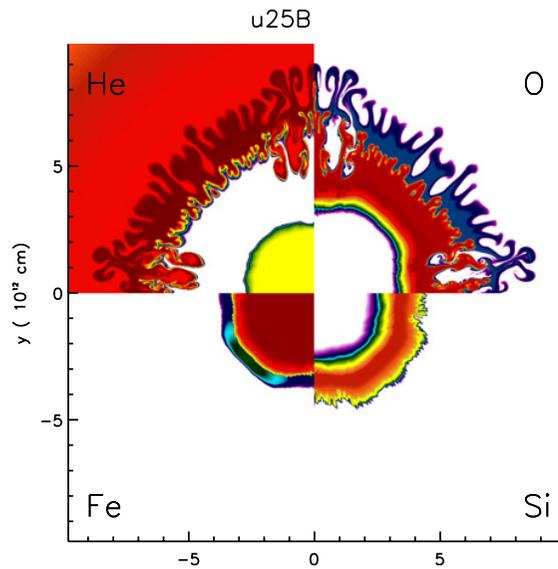
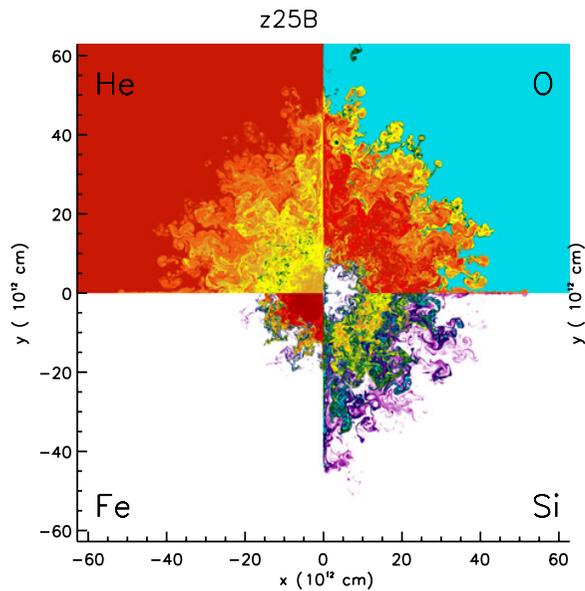
# constraints from EMP stars in halo



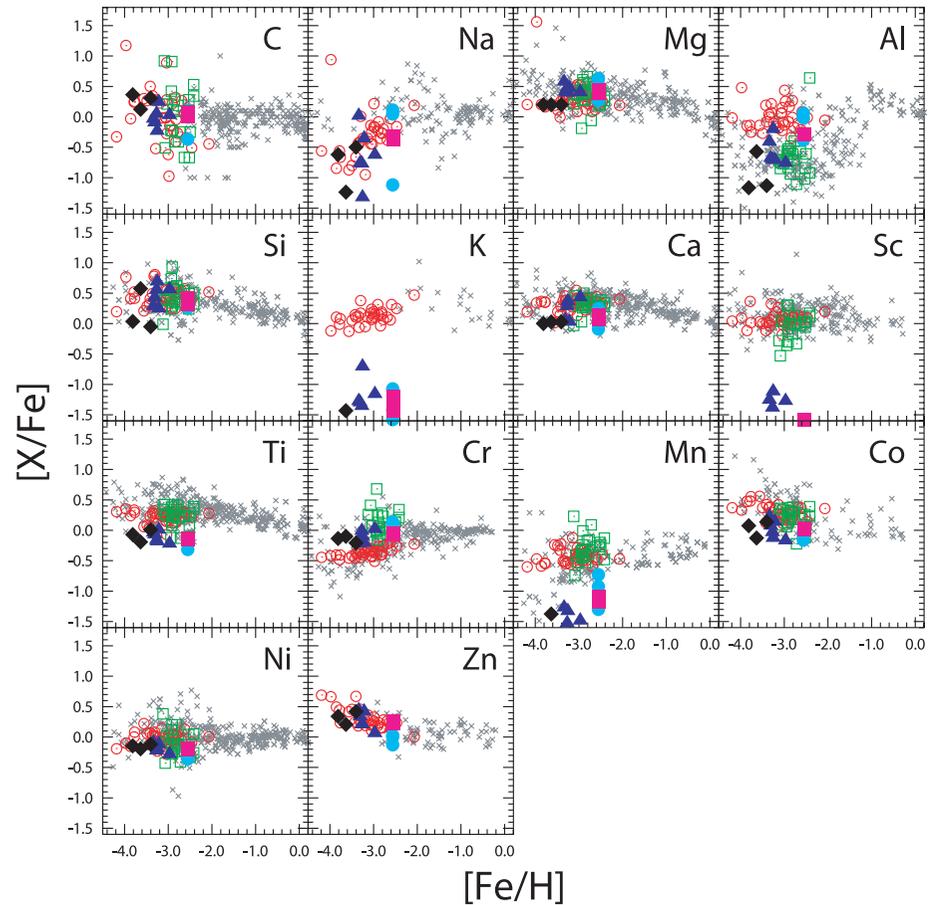
there are many extremely metal-poor stars in the halo (Beers & Christlieb 2005, ARA&A)

- mass range can be explained by dust-induced fragmentation (Clark et al. 2008)
- can use abundance pattern to learn about properties (yields) of progenitor stars

(plot from Salvadori et al. 2006, data from Frebel et al. 2005)



(Joggerst et al. 2009, 2010)



(Tominaga et al. 2007)

The metallicities of extremely metal-poor stars in the halo are consistent with the yields of core-collapse supernovae, i.e. progenitor stars with 20 - 40  $M_{\odot}$

(e.g. Tominaga et al. 2007, Izutani et al. 2009, Joggerst et al. 2009, 2010)

# reasons for debate

- is claim of Pop III stars with  $M \sim 0.5 M_{\odot}$  really justified?
  - stellar collisions
  - magnetic fields
  - radiative feedback
- how would we find them?
  - spectral features
- where should we look?
- what about magnetic fields?

# some more thoughts ...

- magnetic field amplification in primordial collapse (Sur et al. 2010, Federrath et al. 2010, 2011)
- fragmentation-induced starvation as key to understand final stellar masses (Peters et al. 2010abc, 2011)

# conclusions

- primordial star formation exhibits the *same complexity* as stellar birth at present days
    - *turbulence*
    - *thermodynamics*
    - *feedback*
    - *magnetic fields*
- } all influence Pop III and Pop II.5 star formation.



NGC 3324 (Hubble, NASA/ESA)