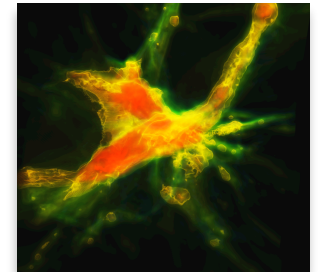
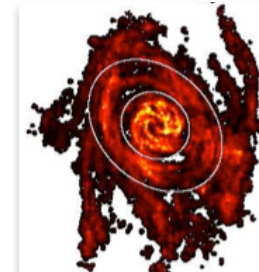
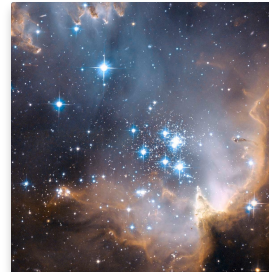
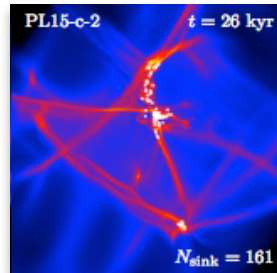
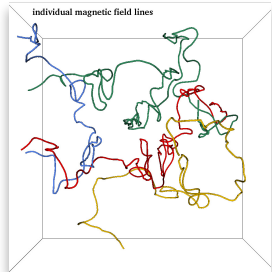


Star Formation (part 2)



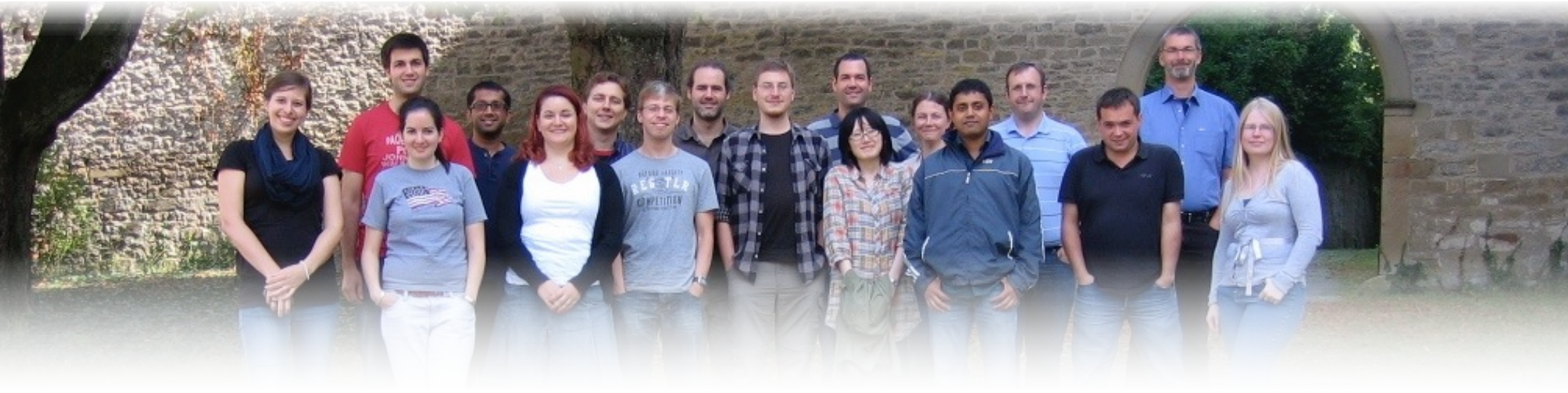
Ralf Klessen



Zentrum für Astronomie der Universität Heidelberg
Institut für Theoretische Astrophysik



thanks to ...



... people in the star formation group at Heidelberg University:

Christian Baczynski, Erik Bertram, Frank Bigiel, Andre Bubel, Diane Cormier, Volker Gaibler, Simon Glover, Dimitrious Gouliermis, Tilman Hartwig, Juan Ibanez, Christoph Klein, Lukas Konstandin, Mei Sasaki, Jennifer Schober, Rahul Shetty, Rowan Smith, László Szűcs

... former group members:

Robi Banerjee, Ingo Berentzen, Paul Clark, Christoph Federrath, Philipp Girichidis, Thomas Greif, Milica Micic, Thomas Peters, Dominik Schleicher, Stefan Schmeja, Sharanya Sur, ...

... many collaborators abroad!



Deutsche
Forschungsgemeinschaft
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**BADEN-
WÜRTTEMBERG**
STIFTUNG
Wir stiften Zukunft



European
Research
Council

agenda

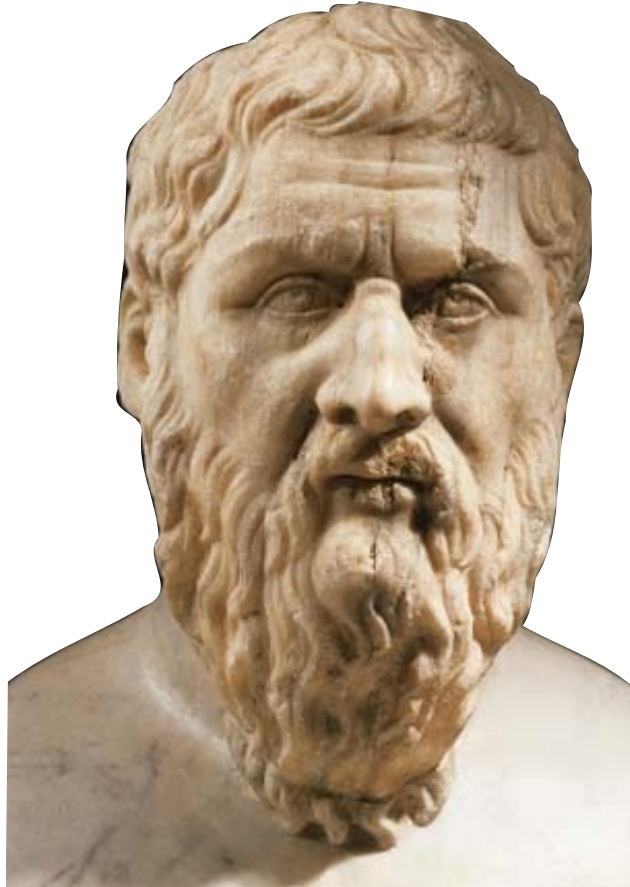
- star formation theory
 - phenomenology
 - historic remarks
 - our current understanding and its limitations
- applications
 - formation of molecular clouds
 - the stellar mass function at birth (IMF)



agenda

- star formation theory
 - phenomenology
 - historic remarks on observations
 - our current understanding and its limitations
- applications & *controversies*
 - *global star formation relations*
 - *ICs for cluster formation*
 - the stellar mass function at birth (IMF)

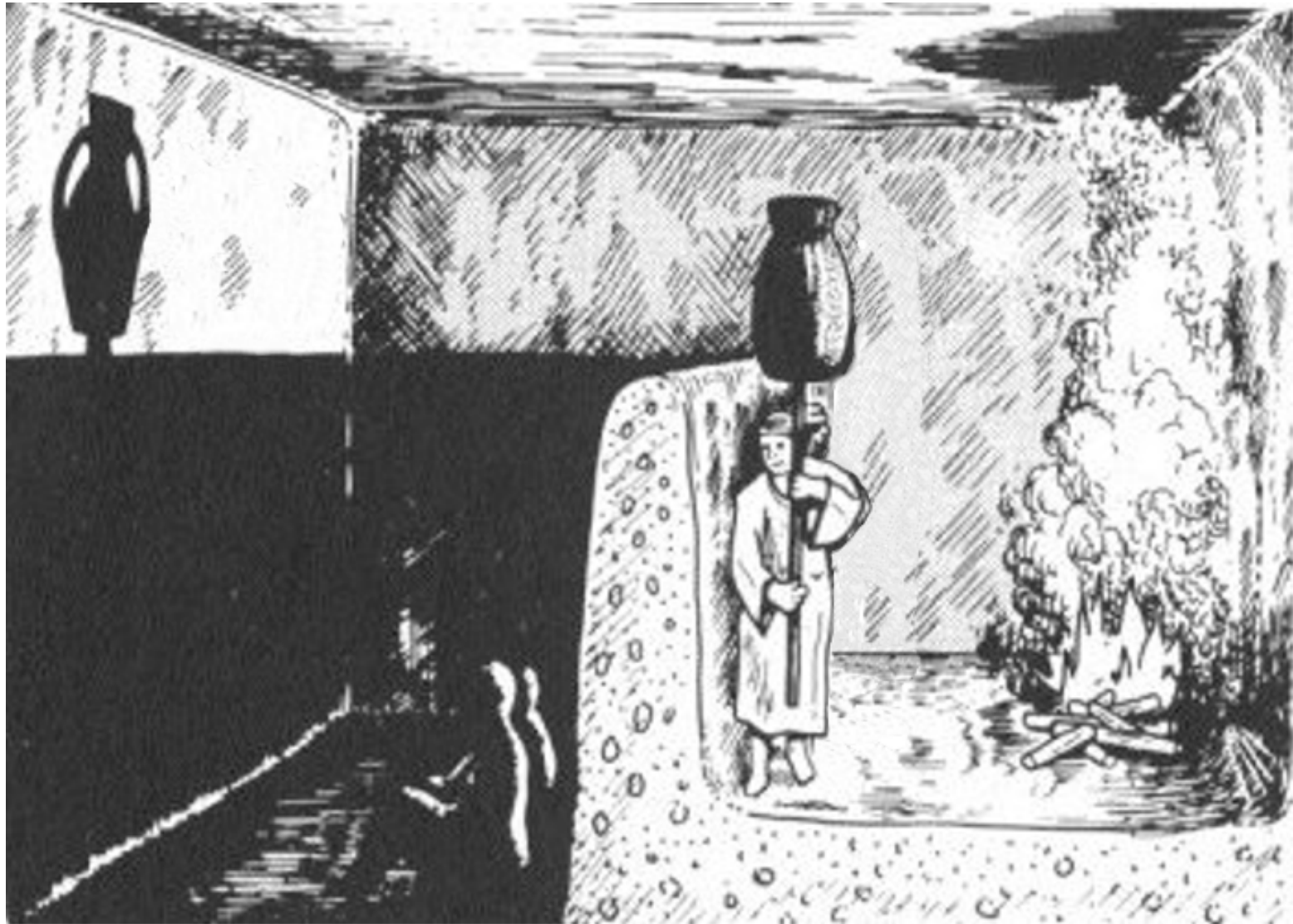




Platon

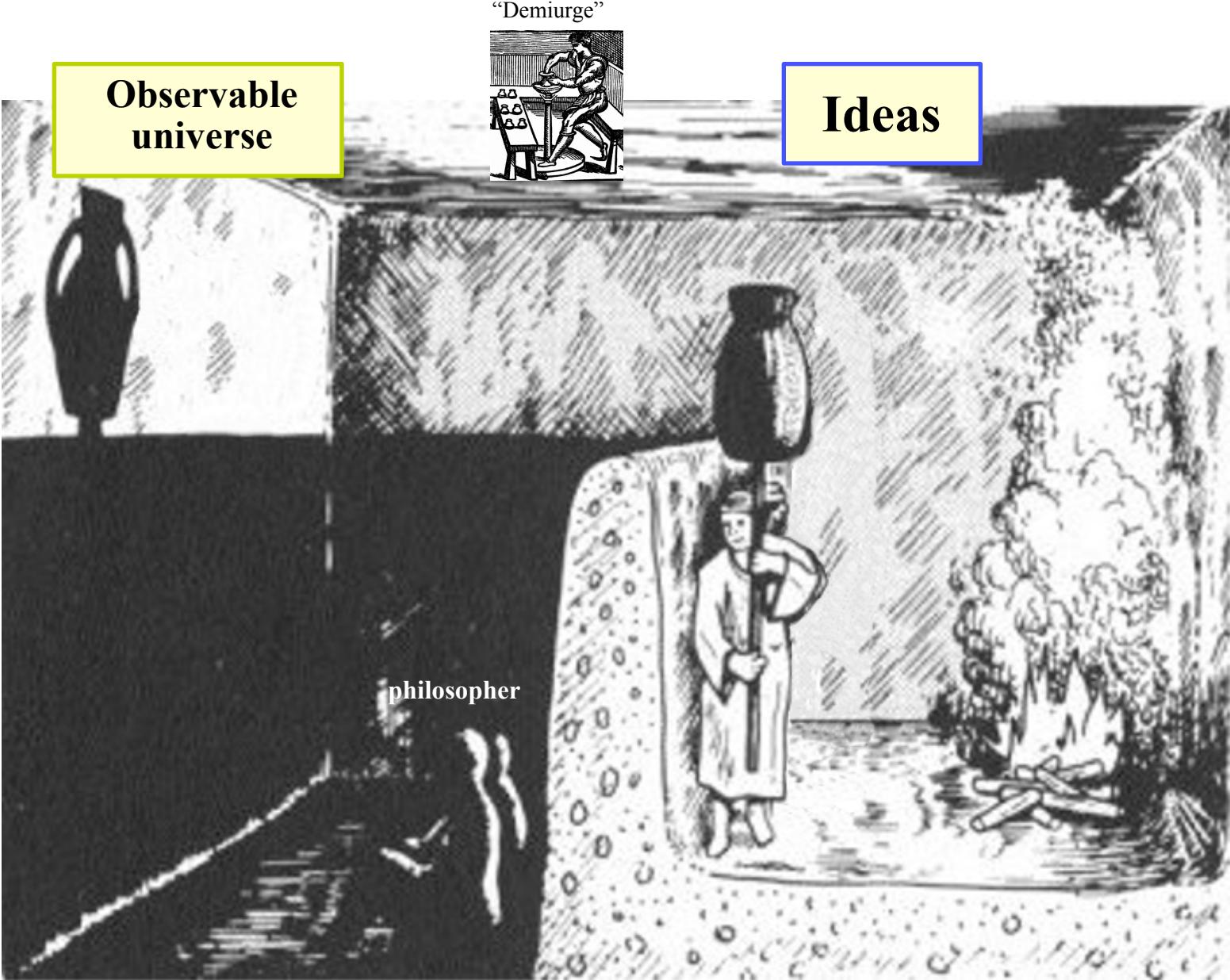
428/427–348/347 BC

Plato's allegory of the cave*



* The Republic
(514a-520a)

Plato's allegory of the cave*



Observable universe

"Demiurge"

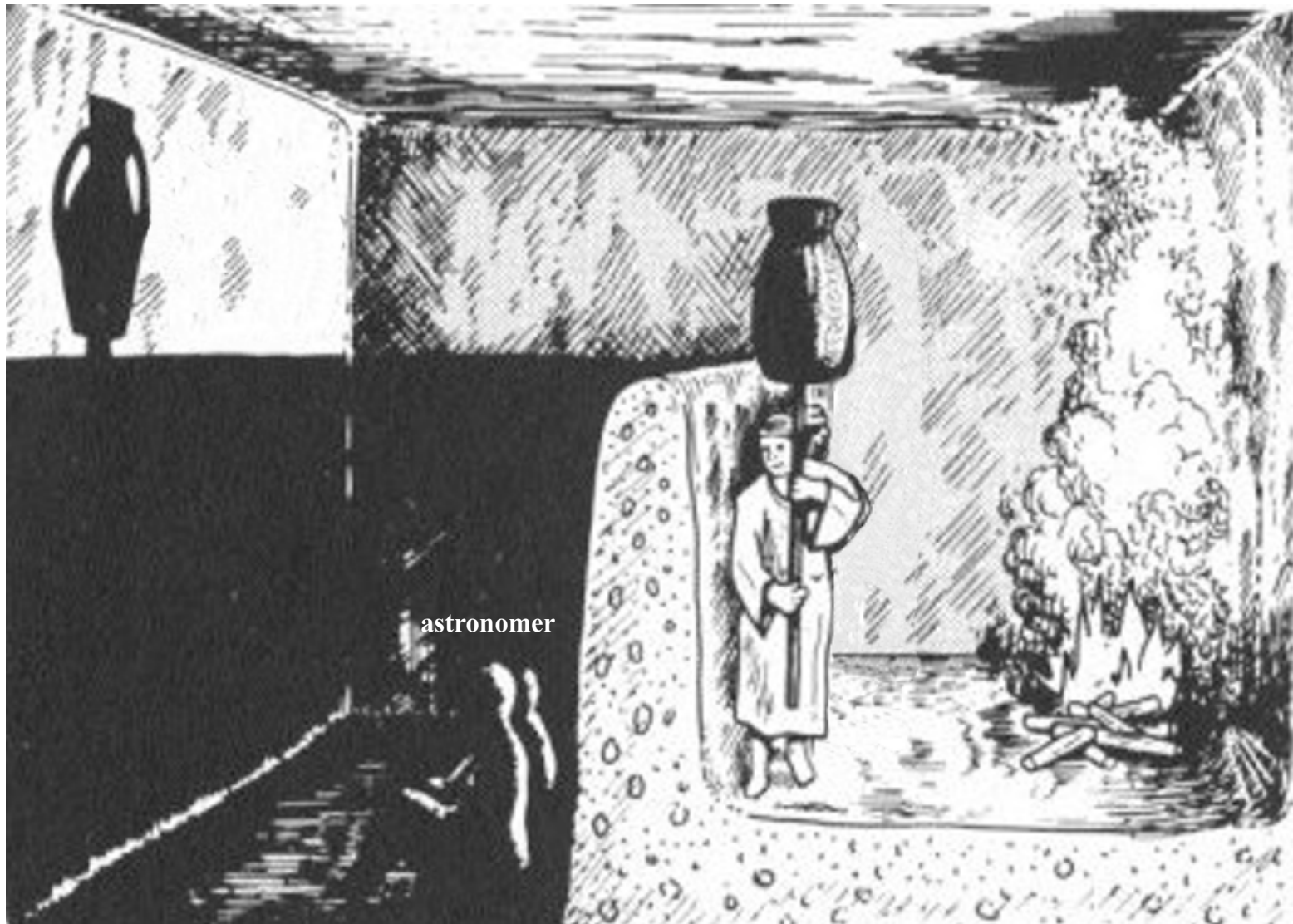


Ideas

philosopher

* The Republic (514a-520a)

Plato's allegory of the cave* ↔ Astronomical observations




* The Republic
(514a-520a)

Plato's allegory of the cave* ↔ Astronomical observations


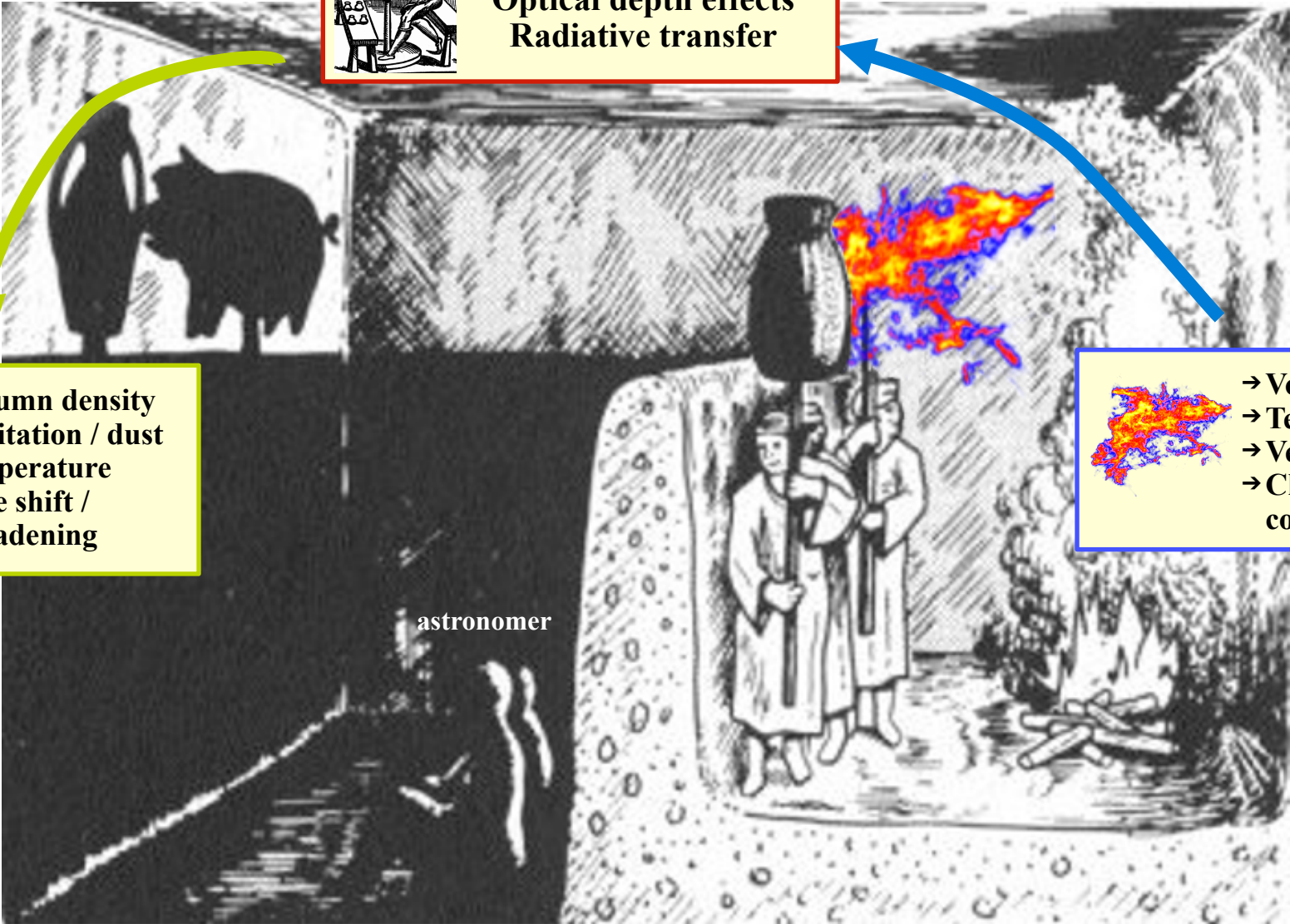


* The Republic
(514a-520a)

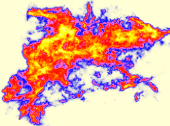
Plato's allegory of the cave* ↔ Astronomical observations



Projection effects
Optical depth effects
Radiative transfer



→ Column density
→ Excitation / dust temperature
→ Line shift / broadening

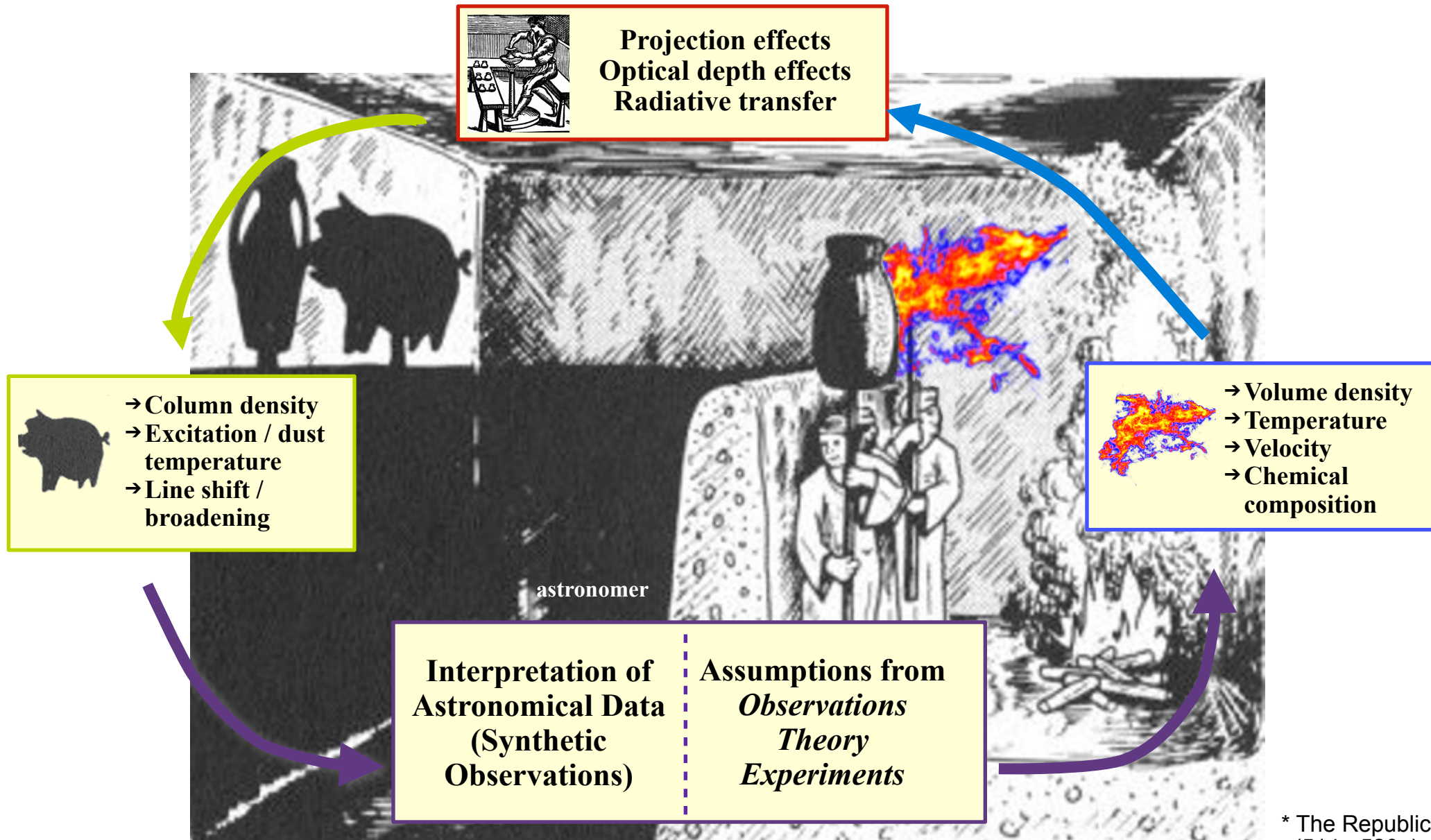


→ Volume density
→ Temperature
→ Velocity
→ Chemical composition

astronomer

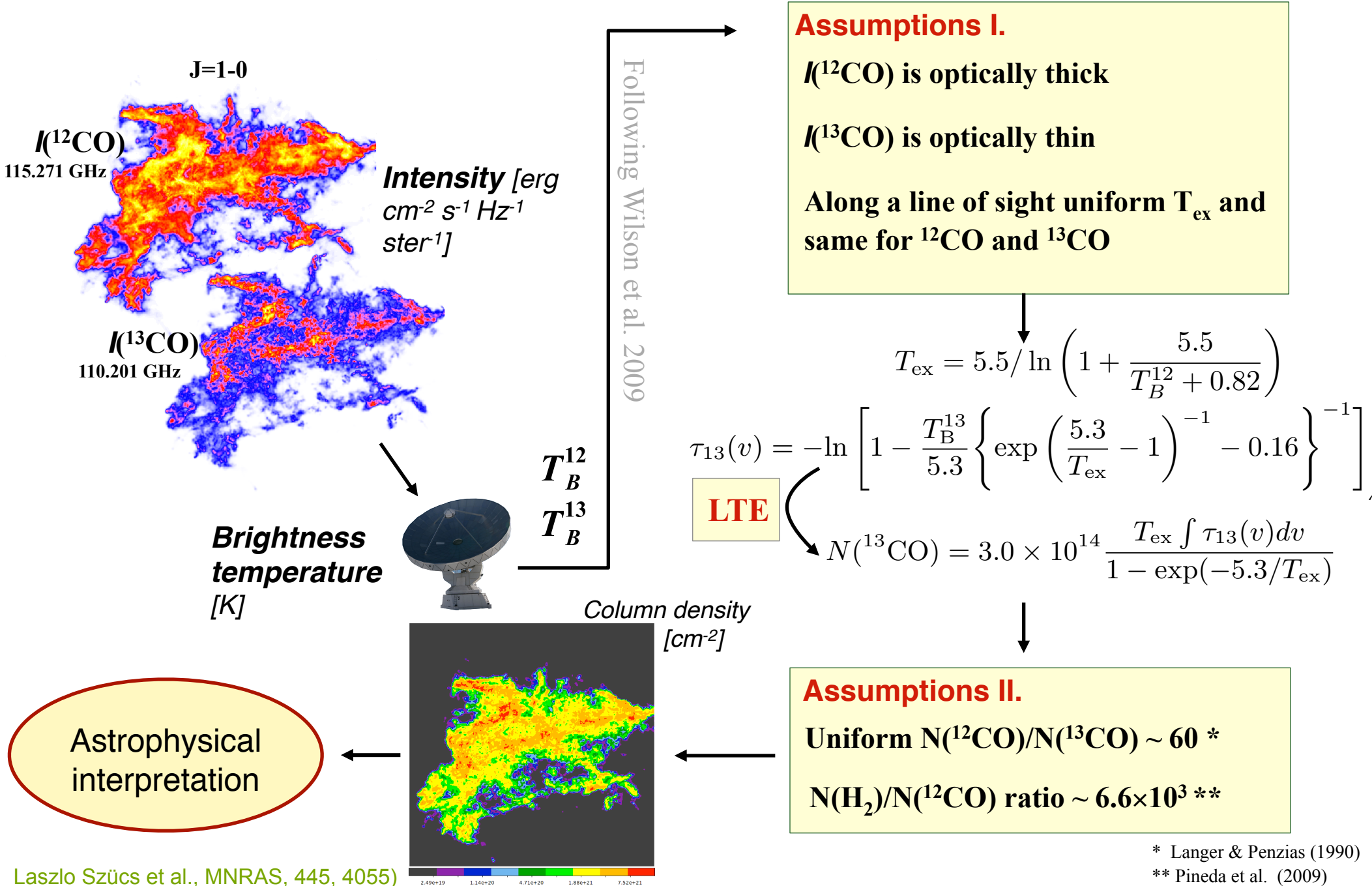
* The Republic (514a-520a)

Plato's allegory of the cave* ↔ Astronomical observations



* The Republic (514a-520a)

Example: from CO emission to total column density





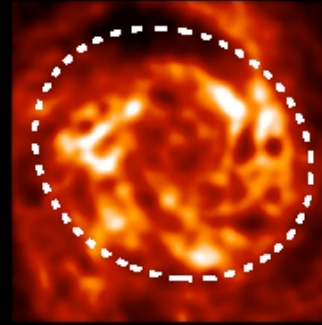
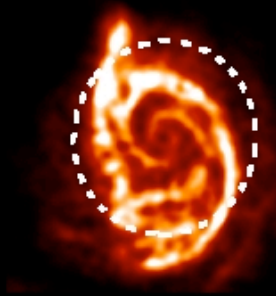
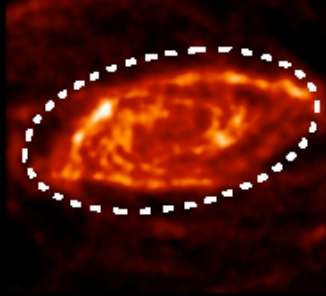
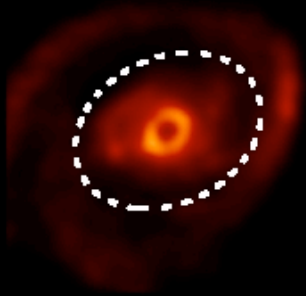
global SF relations

NGC 4736

NGC 5055

NGC 5194

NGC 6946



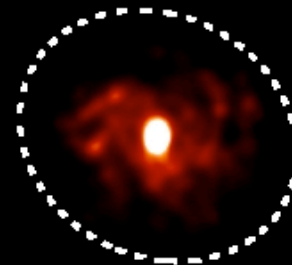
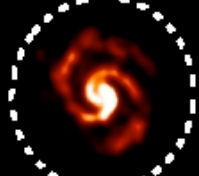
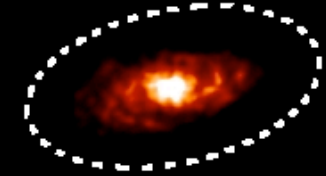
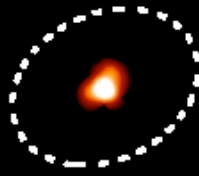
atomic
hydrogen

NGC 4736

NGC 5055

NGC 5194

NGC 6946



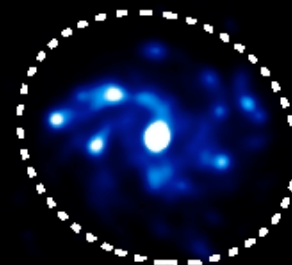
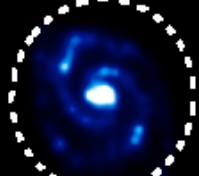
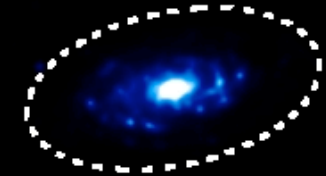
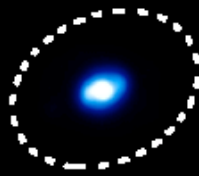
molecular
hydrogen

NGC 4736

NGC 5055

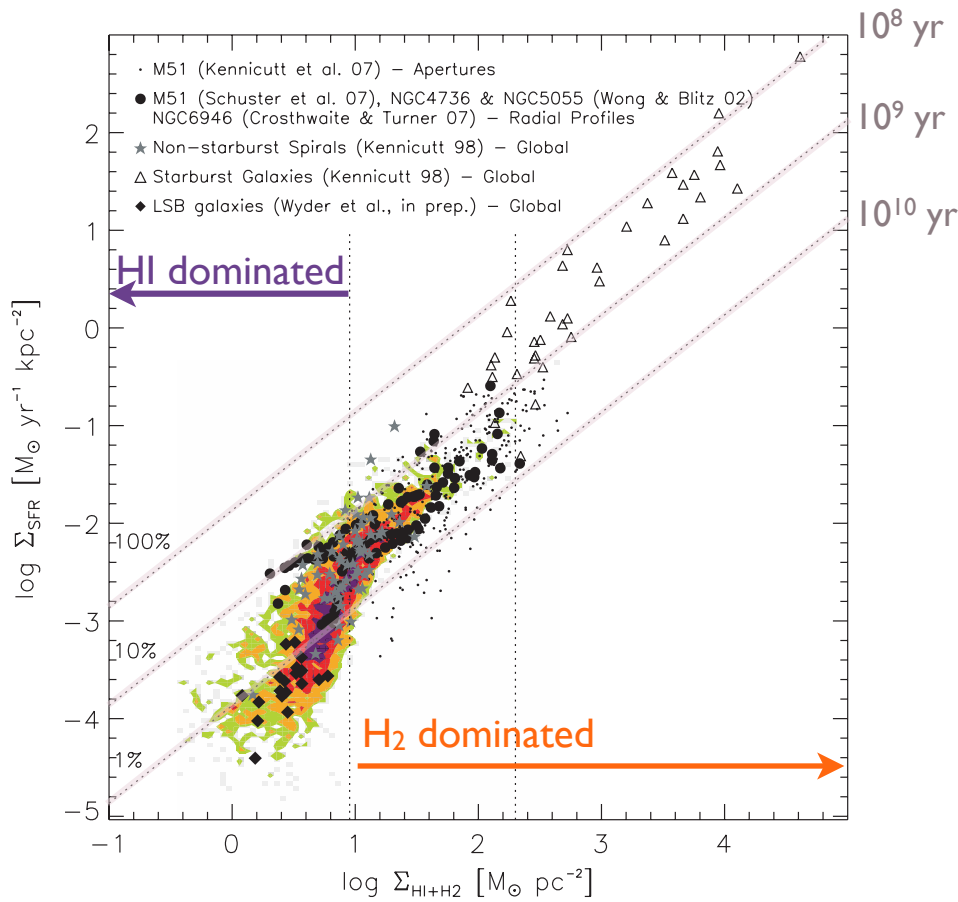
NGC 5194

NGC 6946

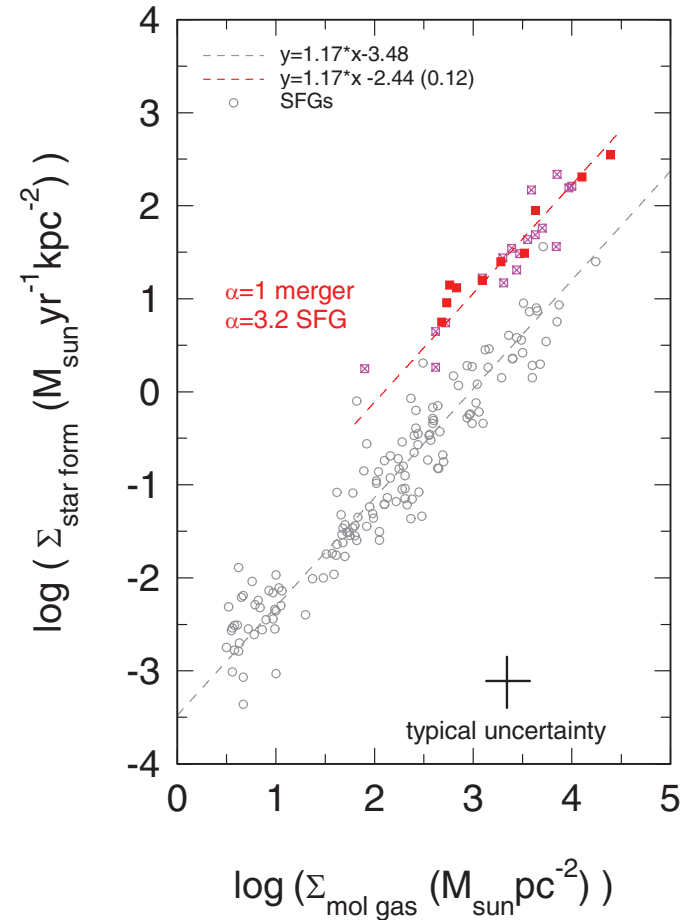


star
formation

- HI gas more extended
- H2 and SF well correlated



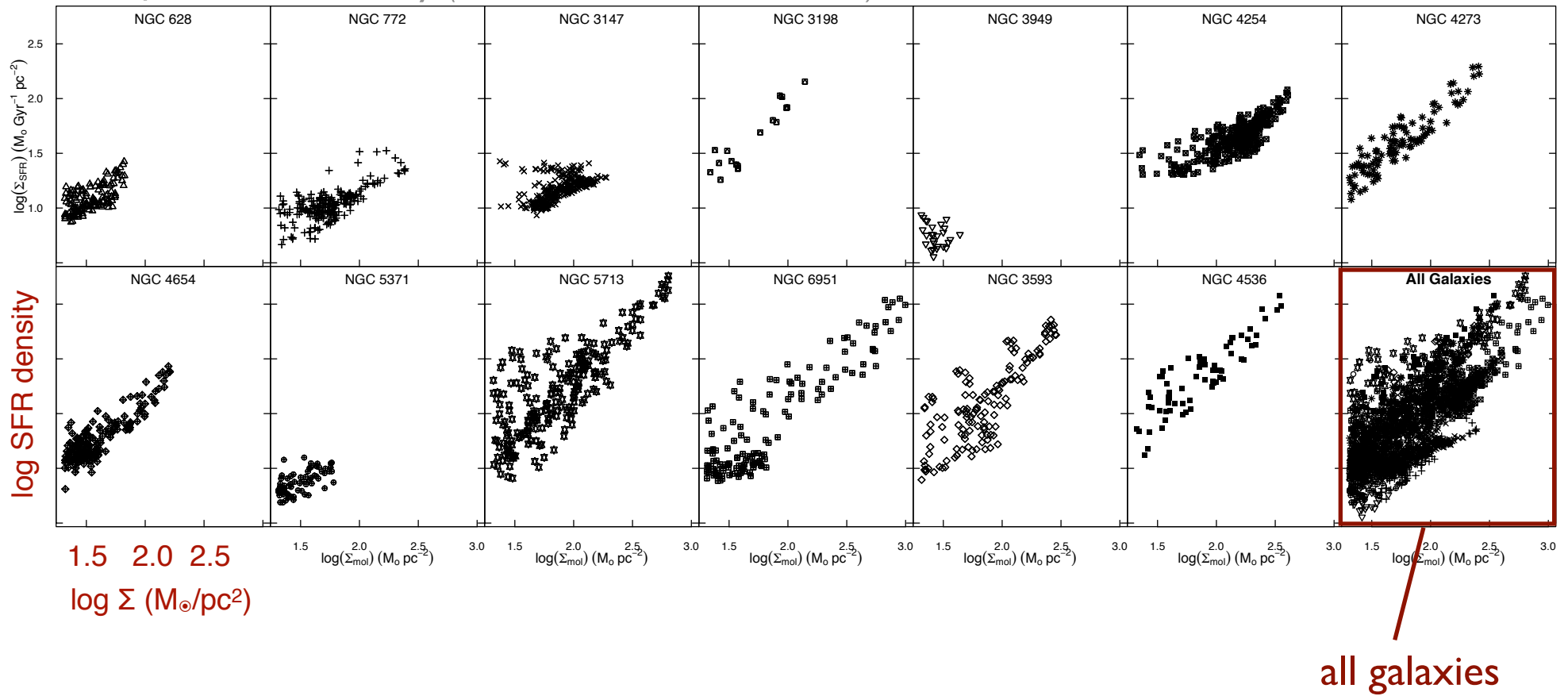
Bigiel et al. (2008, AJ, 136, 2846)



Genzel et al. (2010, MNRAS, AJ, 407, 2091)

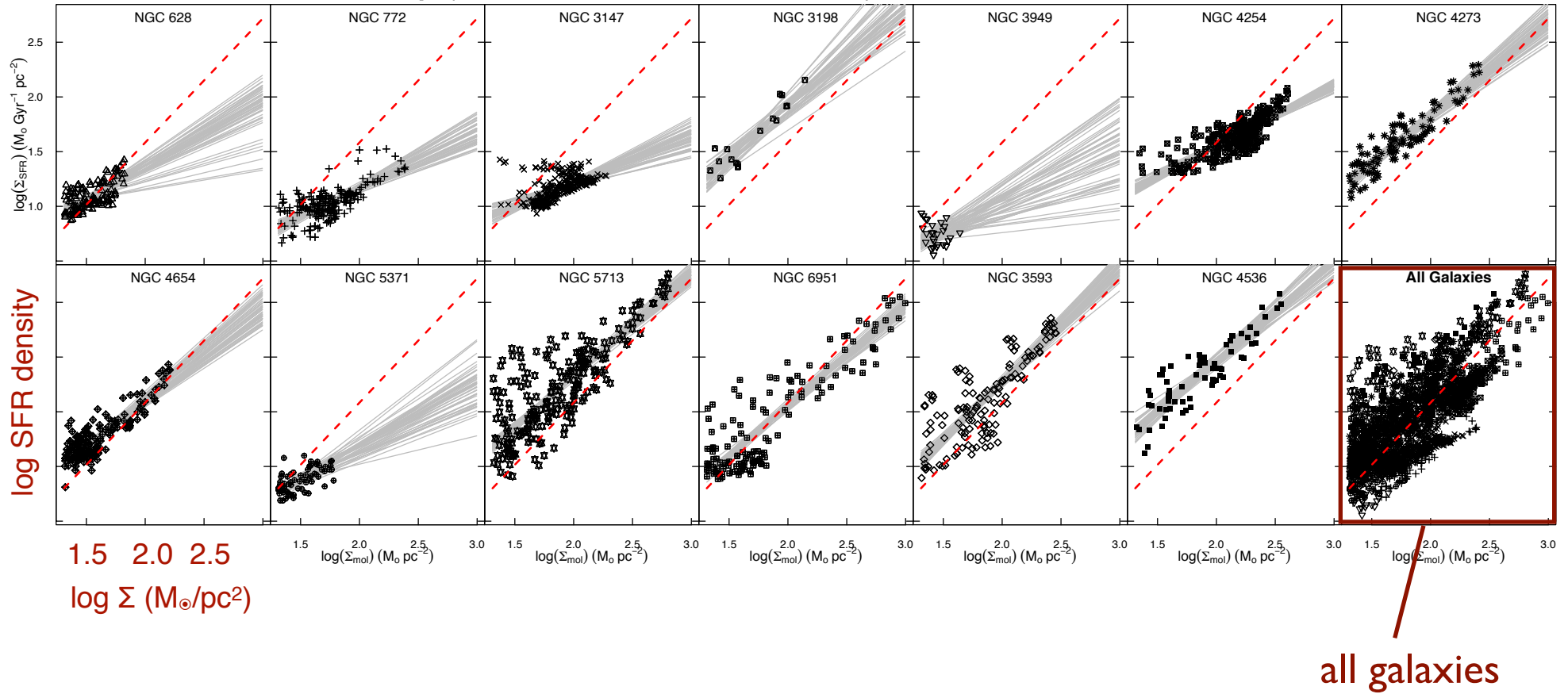
- standard model: roughly linear relation between H
- standard model: roughly constant depletion time: few x 10
- super linear relation between total gas and SFR

data from STING survey (Rahman et al. 2011, 2012)



- QUIZ: do you see a universal

data from STING survey (Rahman et al. 2011, 2012)



- QUIZ: do you see a universal
- ANSWER: - probably not
- in addition, the relation often is sublinear

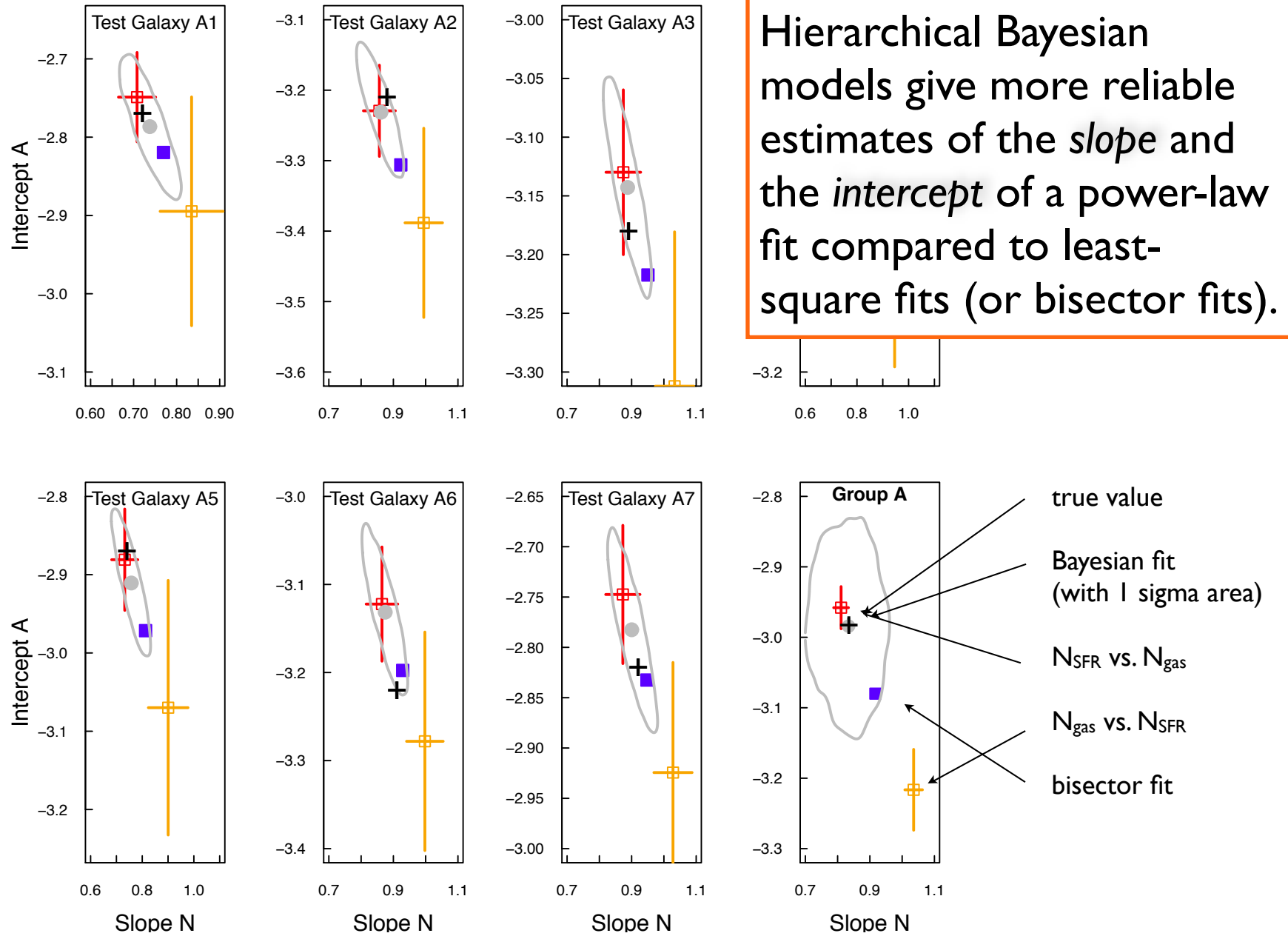
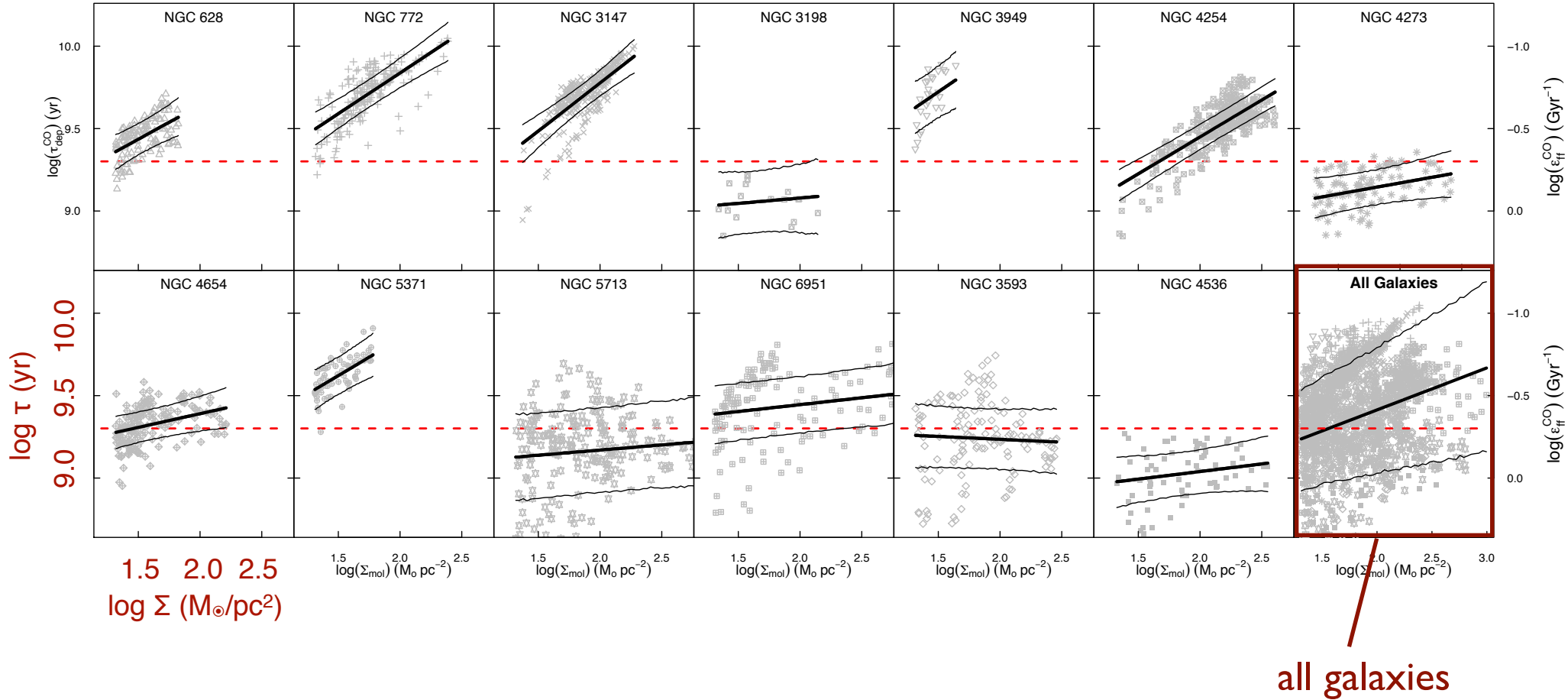


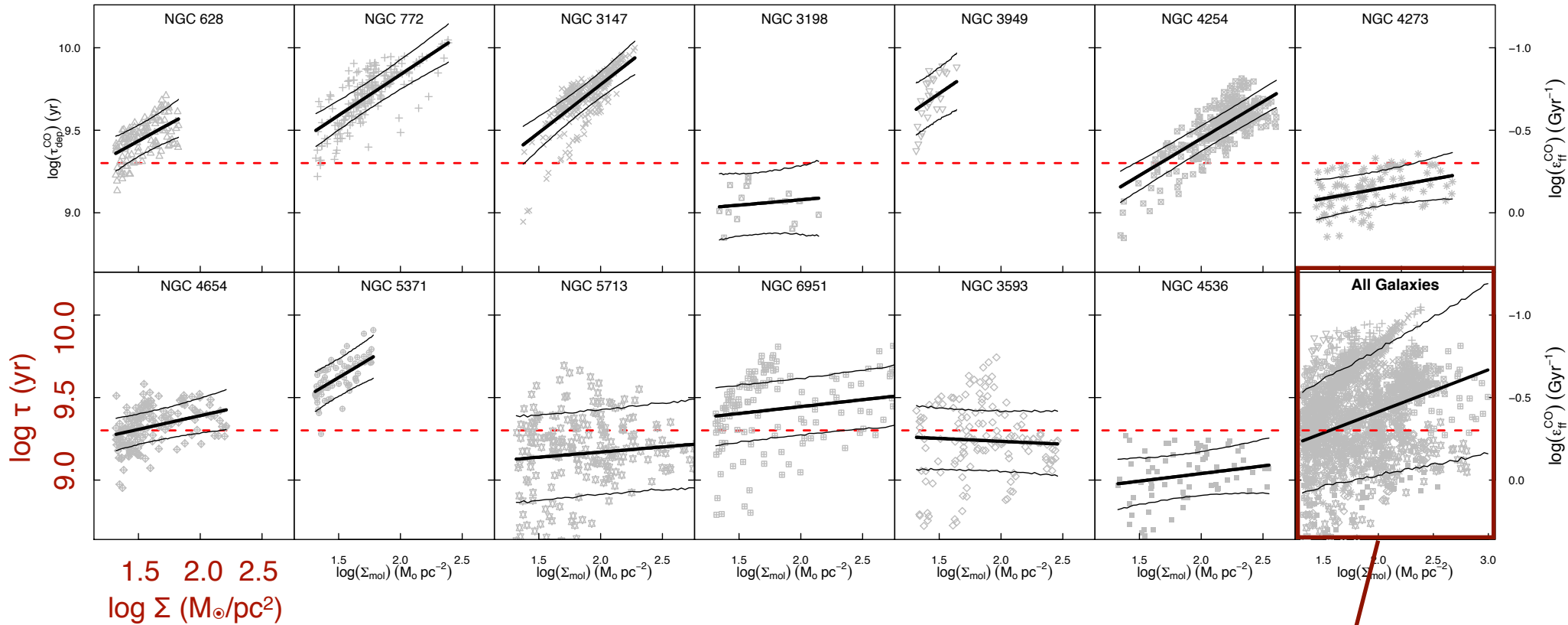
Figure 1. Slope and intercept of test galaxies in Group A. Black cross shows the true values. Red and orange squares show the $OLS(\Sigma_{SFR}|\Sigma_{mol})$ and $OLS(\Sigma_{mol}|\Sigma_{SFR})$ results, with their 1σ uncertainties, respectively. The gray circles indicate the estimate provided by the median of hierarchical Bayesian posterior result, and the contours mark the 1σ deviation. The filled blue squares mark the bisector estimates. The last panel on the bottom row shows the group parameters and fit estimates.

data from STING survey (Rahman et al. 2011, 2012)



Hierarchical Bayesian model for STING galaxies indicate *varying depleting times*.

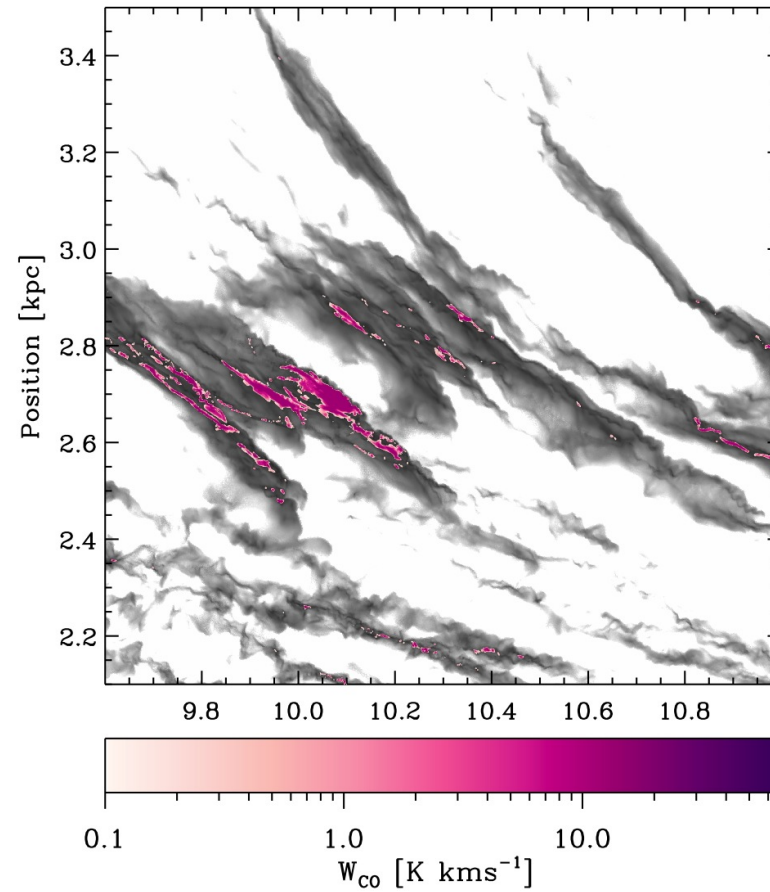
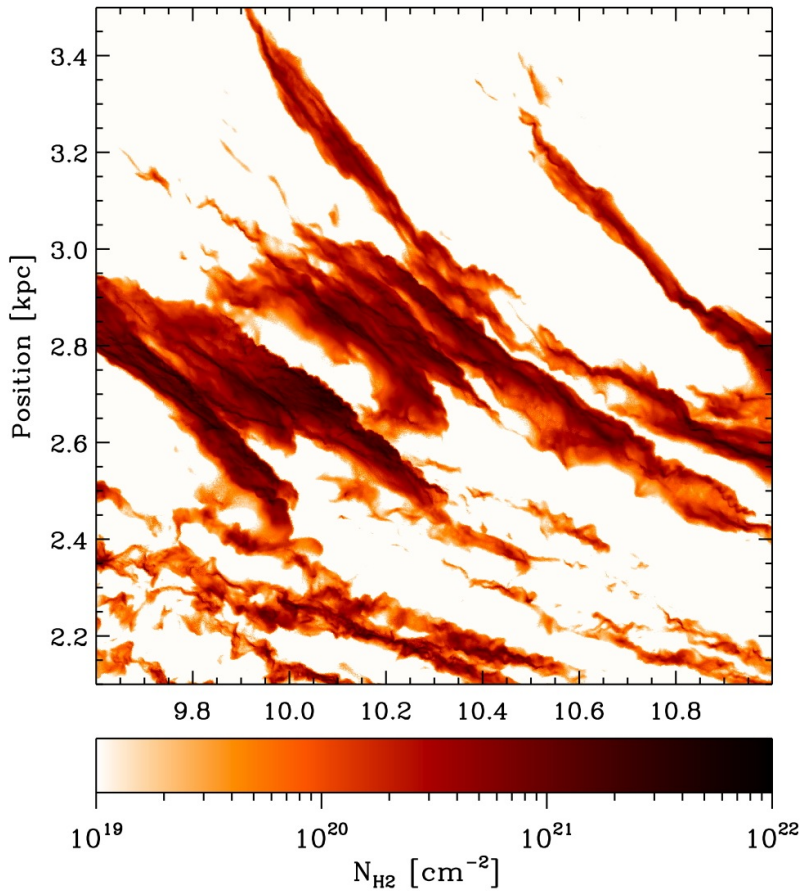
data from STING survey (Rahman et al. 2011, 2012)



all galaxies

physical origin of this behavior?

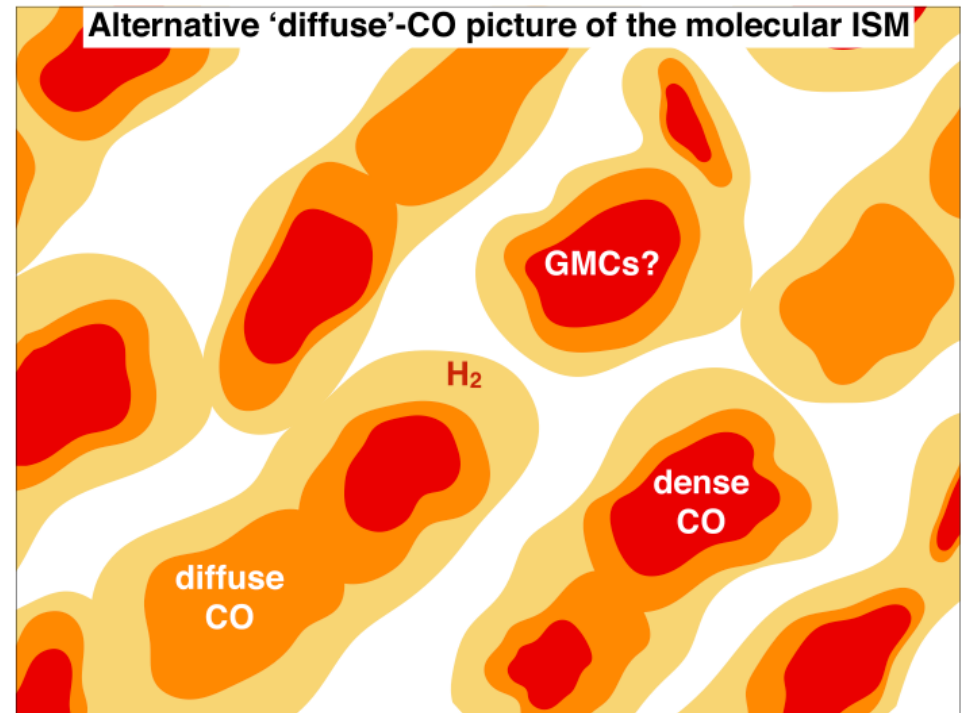
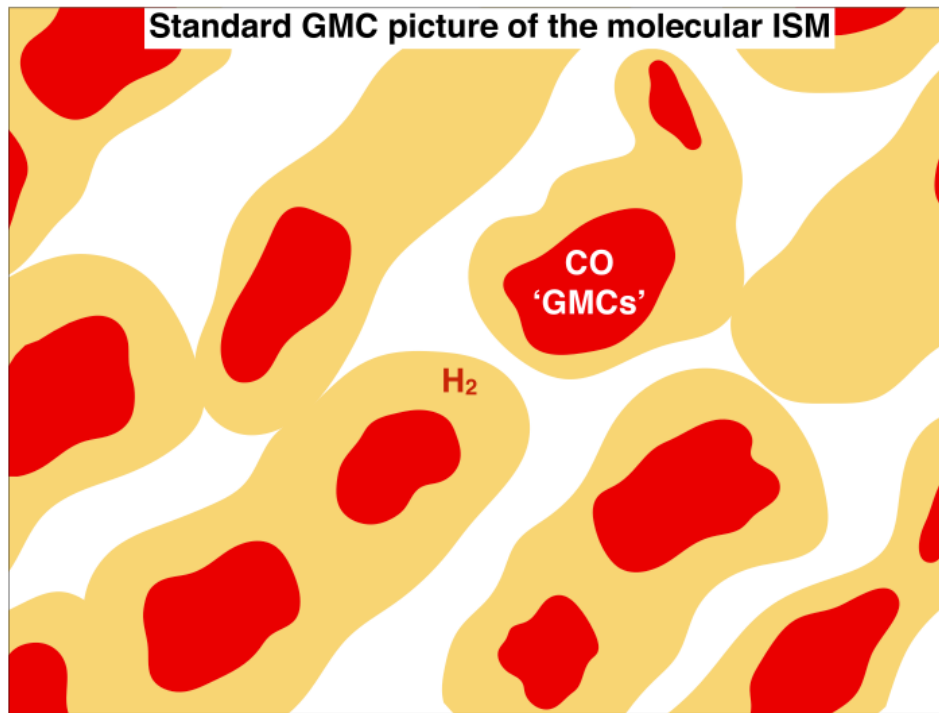
- maybe strong shear in dense arms (example M51, Meidt et al. 2013)...
- maybe non-star forming H densities (recall H [see part I on December 19, 2014])



Smith et al. (2014, MNRAS, 441, 1628)

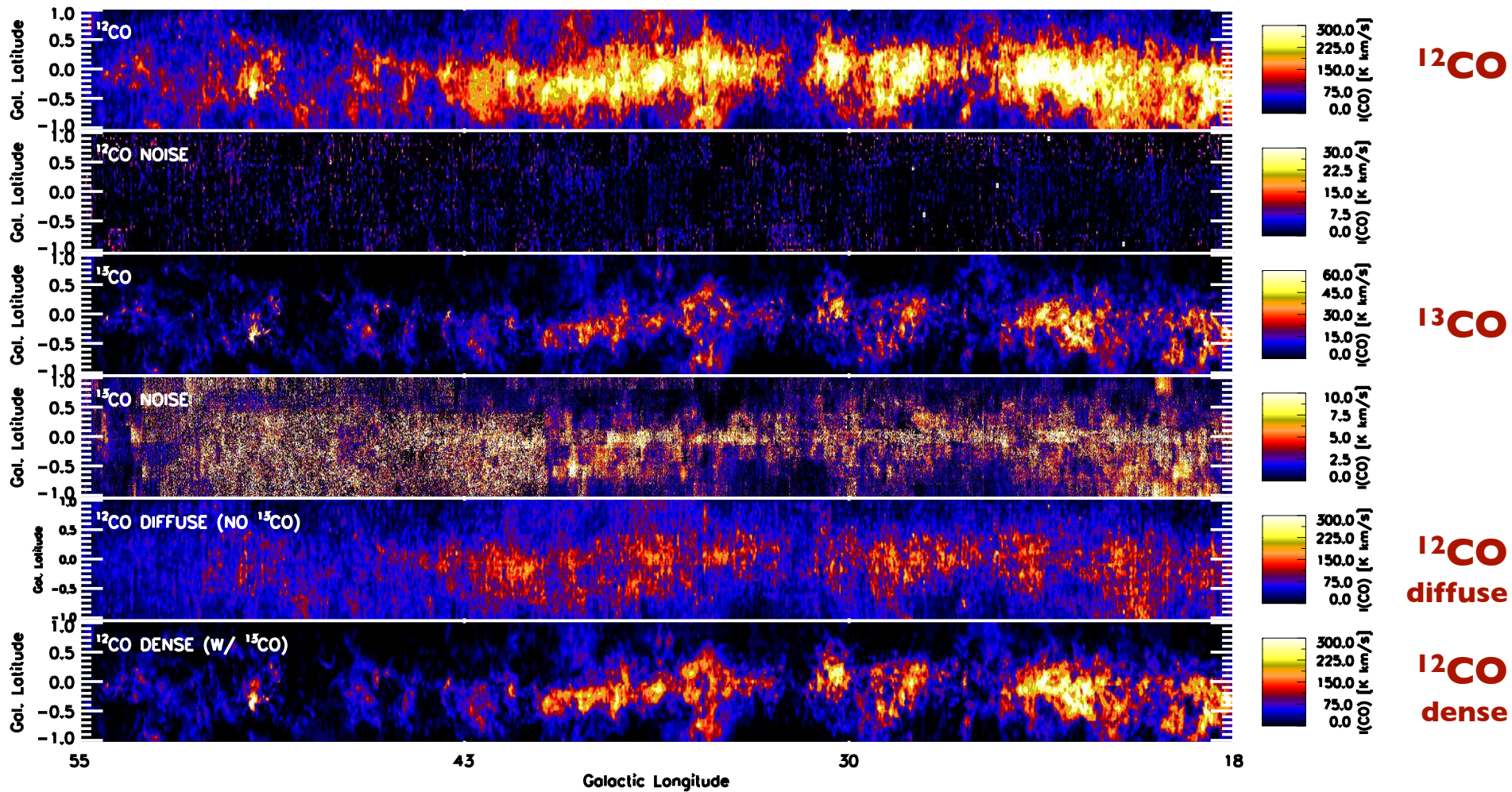
physical origin of this behavior?

- maybe strong shear in dense arms (example M5 I, Meidt et al. 2013)...
- maybe non-star forming H densities (recall H [see part I on December 19, 2014])



in addition:

- maybe a large fraction of H₂ dense clouds, but in a diffuse state!



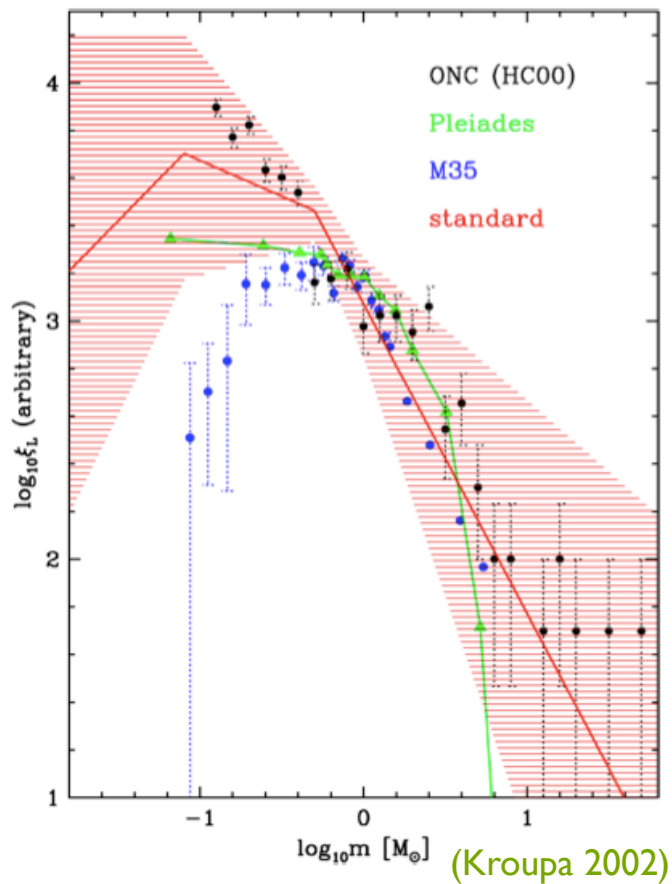
Galactic Ring Survey (GRS)

- comparison of tracing all the gas (including the more diffuse component)

stellar mass
function

stellar mass function

stars seem to follow a universal mass function at birth --> IMF



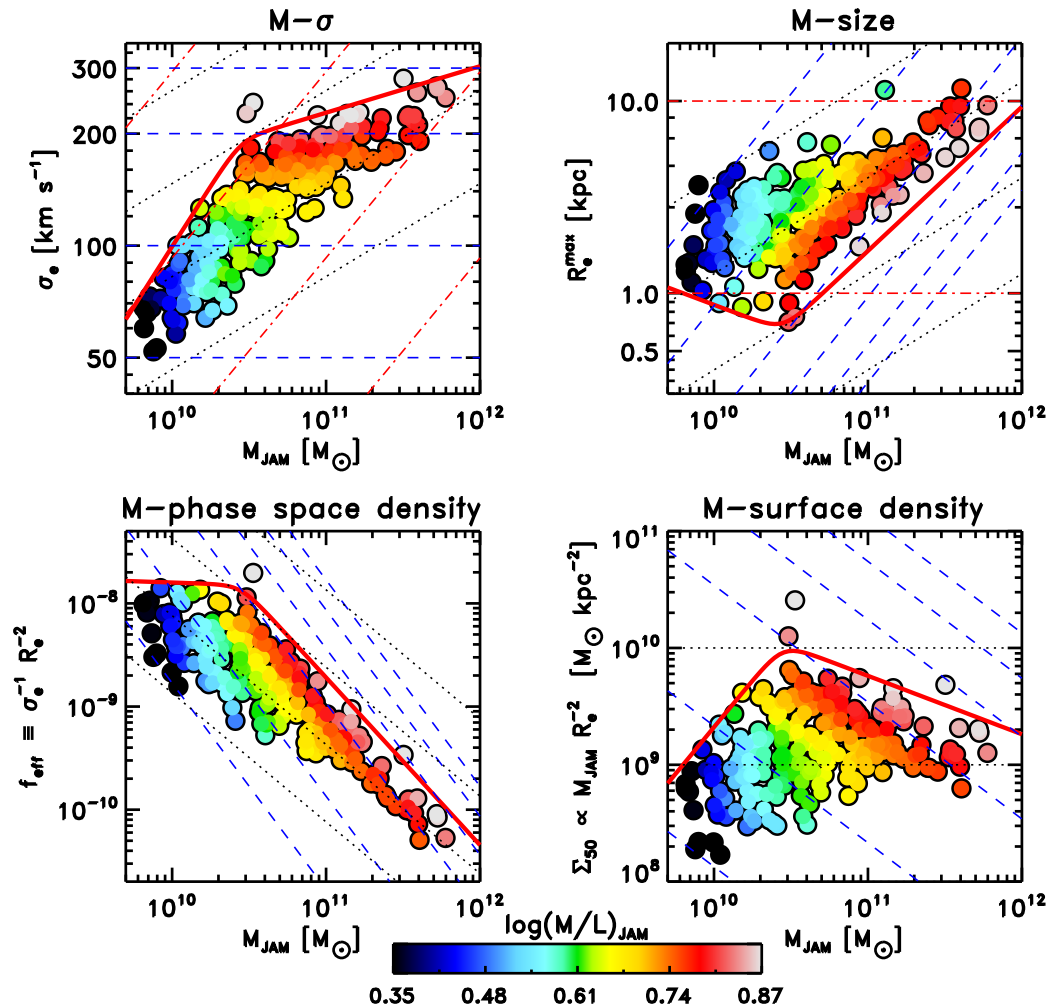
Orion, NGC 3603, 30 Doradus
(Zinnecker & Yorke 2007)

stellar mass function

BUT: maybe variations with galaxy type (bottom heavy in the centers of large ellipticals)

from JAM (Jeans anisotropic multi Gaussian expansion) modeling

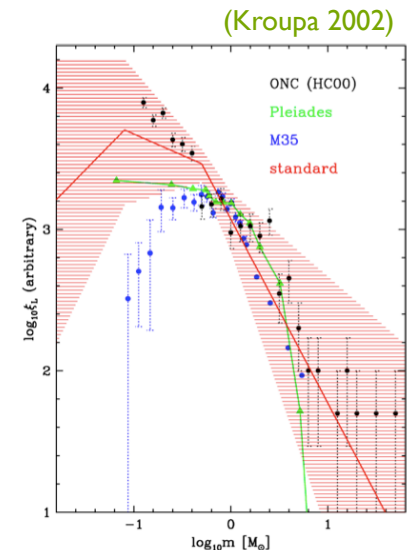
inferred excess of low-mass stars compared to Kroupa IMF



(Cappellari et al. 2012, Nature, 484, 485, Cappellari et al. 2012ab, MNRAS, submitted, also van Dokkum & Conroy 2010, Nature, 468, 940, Wegner et al. 2012, AJ, 144, 78, and others)

stellar masses

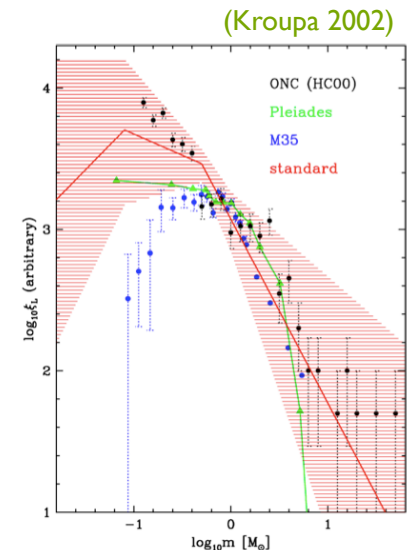
- distribution of stellar masses depends on
 - turbulent initial conditions
 - > mass spectrum of prestellar cloud cores
 - collapse and interaction of prestellar cores
 - > accretion and N -body effects
 - thermodynamic properties of gas
 - > balance between heating and cooling
 - > EOS (determines which cores go into collapse)
 - (proto) stellar feedback terminates star formation
 - ionizing radiation, bipolar outflows, winds, SN



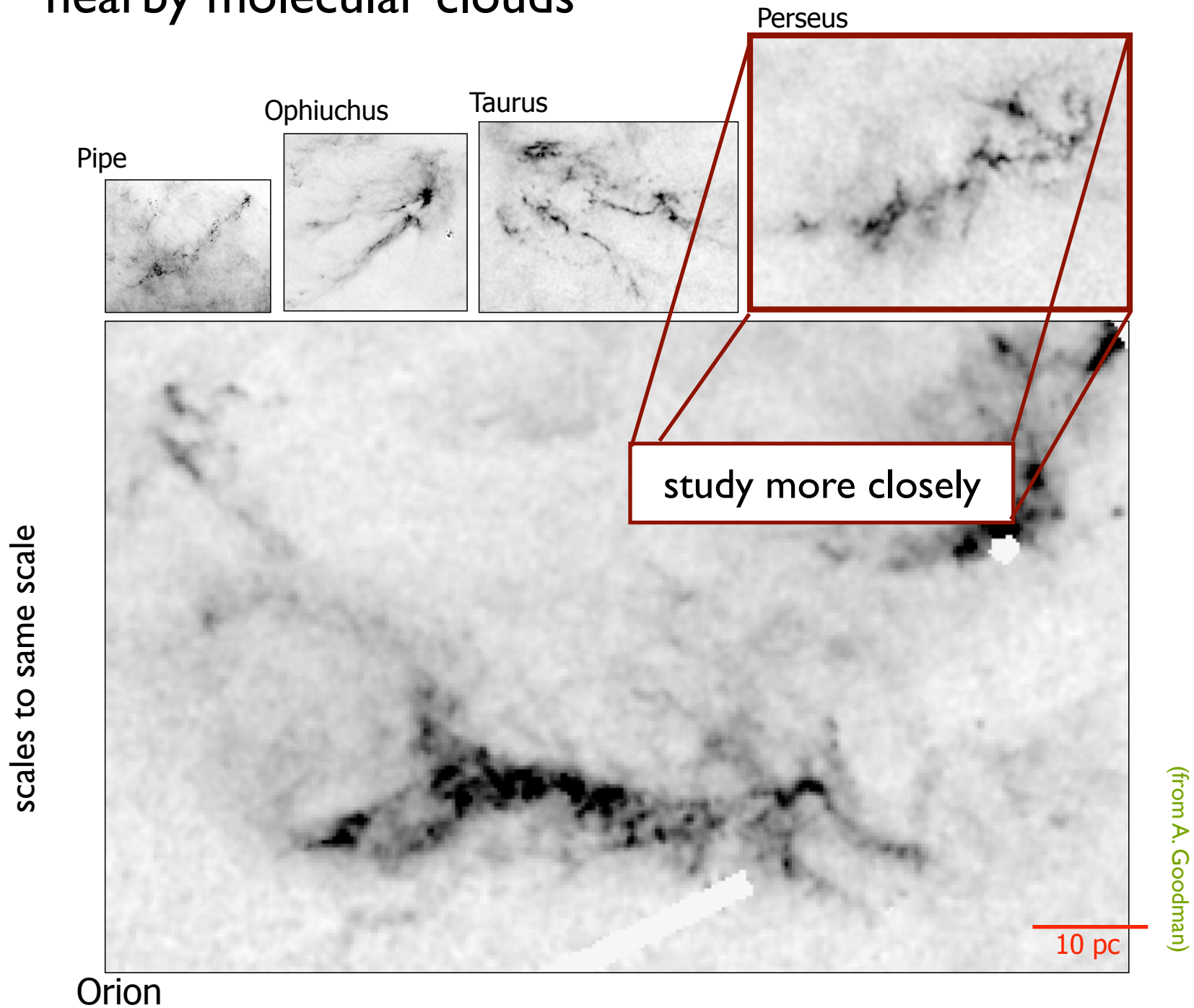
stellar masses

- distribution of stellar masses depends on

- turbulent initial conditions
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- collapse and interaction of prestellar cores
--> accretion and N -body effects
- thermodynamic properties of gas
--> balance between heating and cooling
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nearby molecular clouds



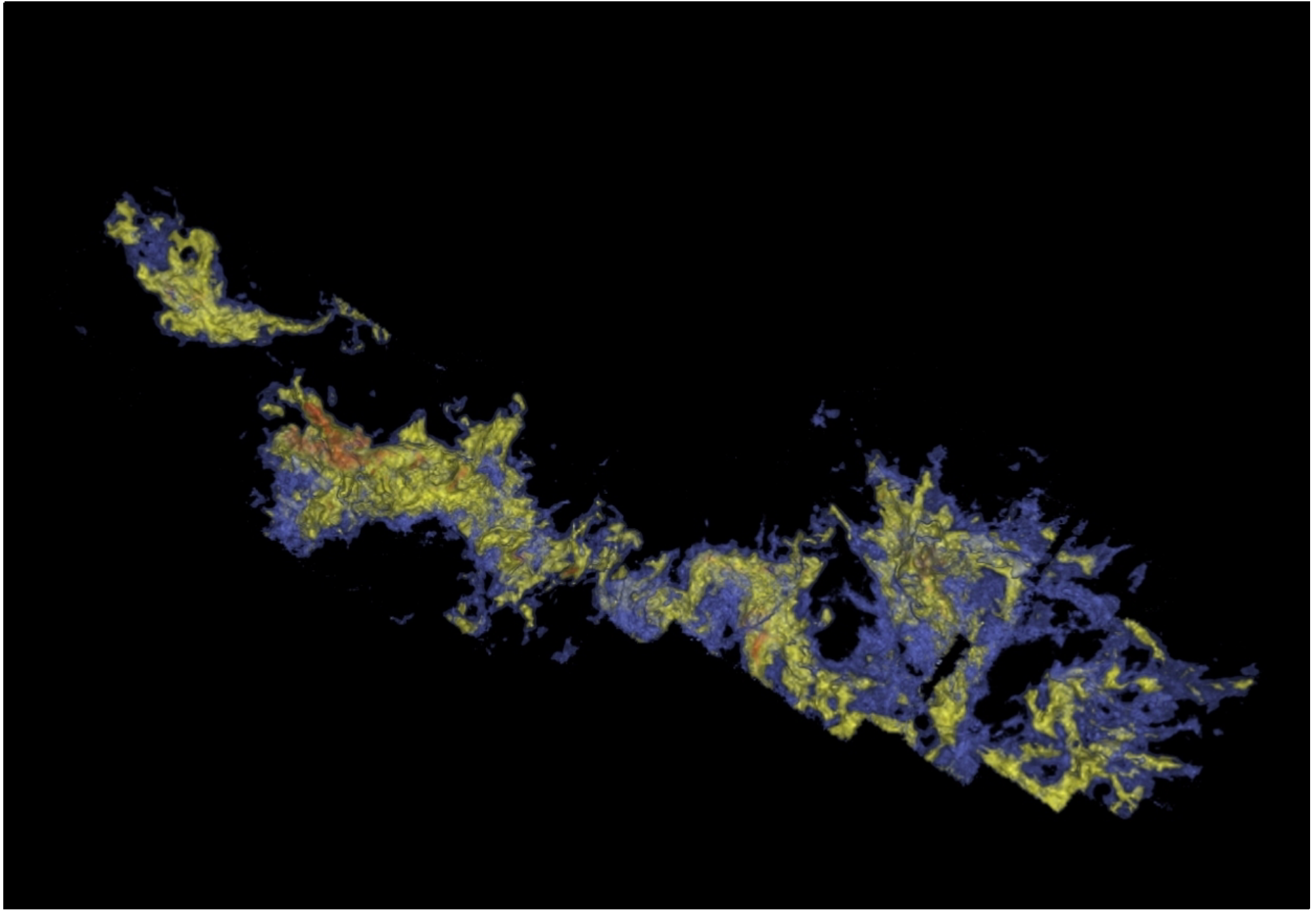
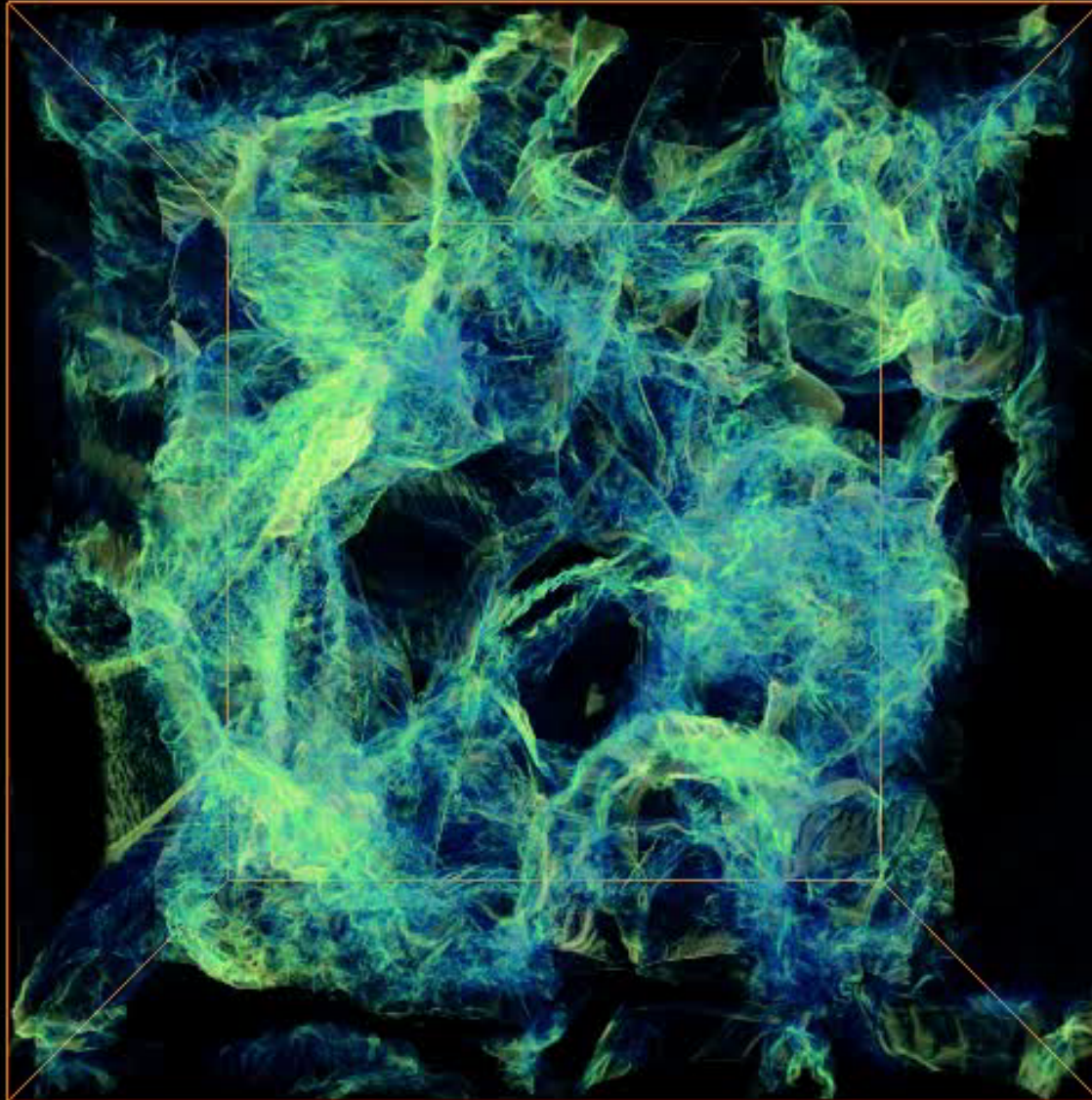


image from Alyssa Goodman: COMPLETE survey



Schmidt et al. (2009, A&A, 494, 127)



example: model of Orion cloud

„model“ of Orion cloud:

15.000.000 SPH particles,

$10^4 M_{\text{sun}}$ in 10 pc, mass resolution

$0,02 M_{\text{sun}}$, forms ~ 2.500

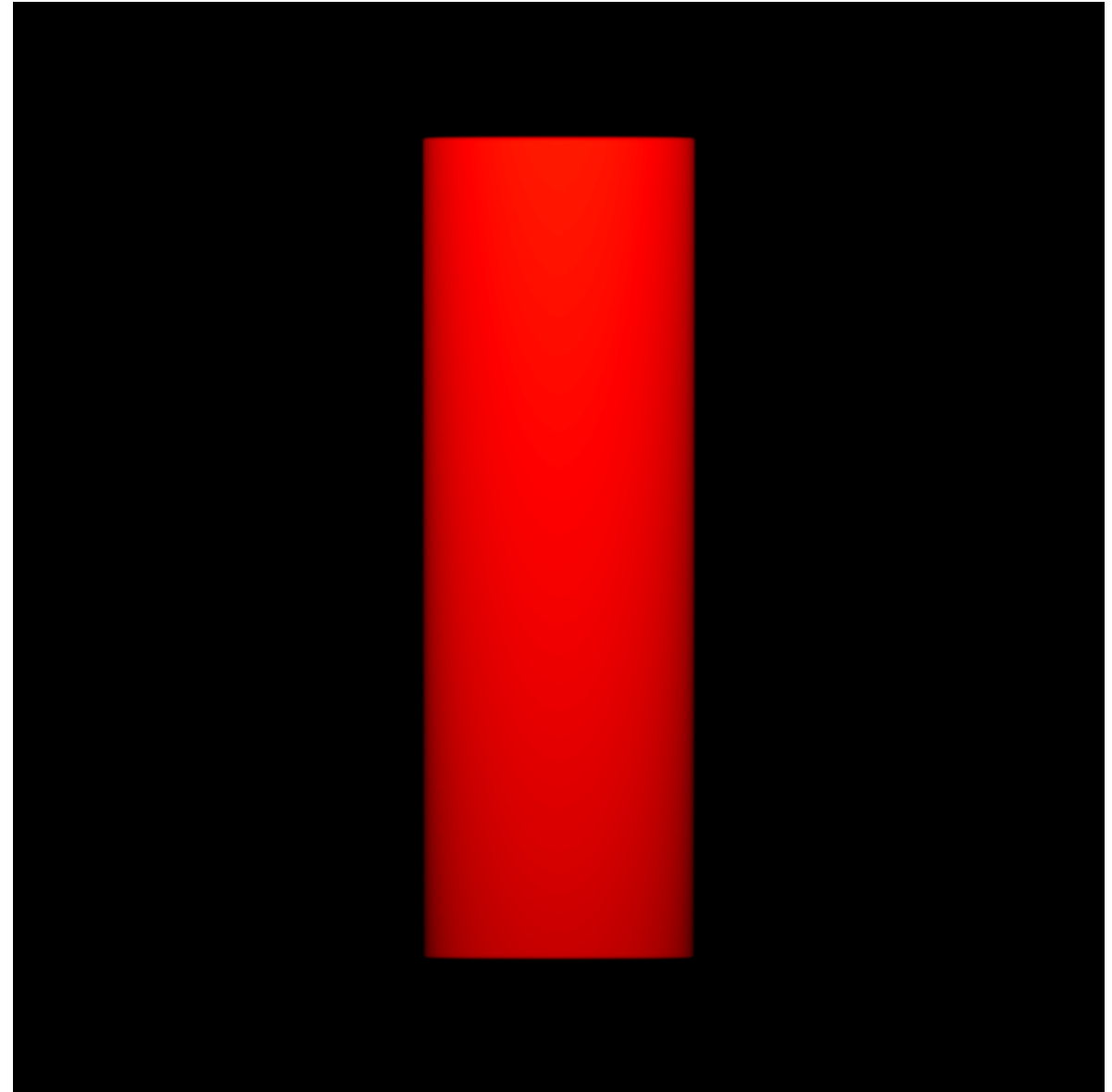
„stars“ (sink particles)

isothermal EOS, top bound, bottom unbound

has clustered as well as distributed „star“ formation

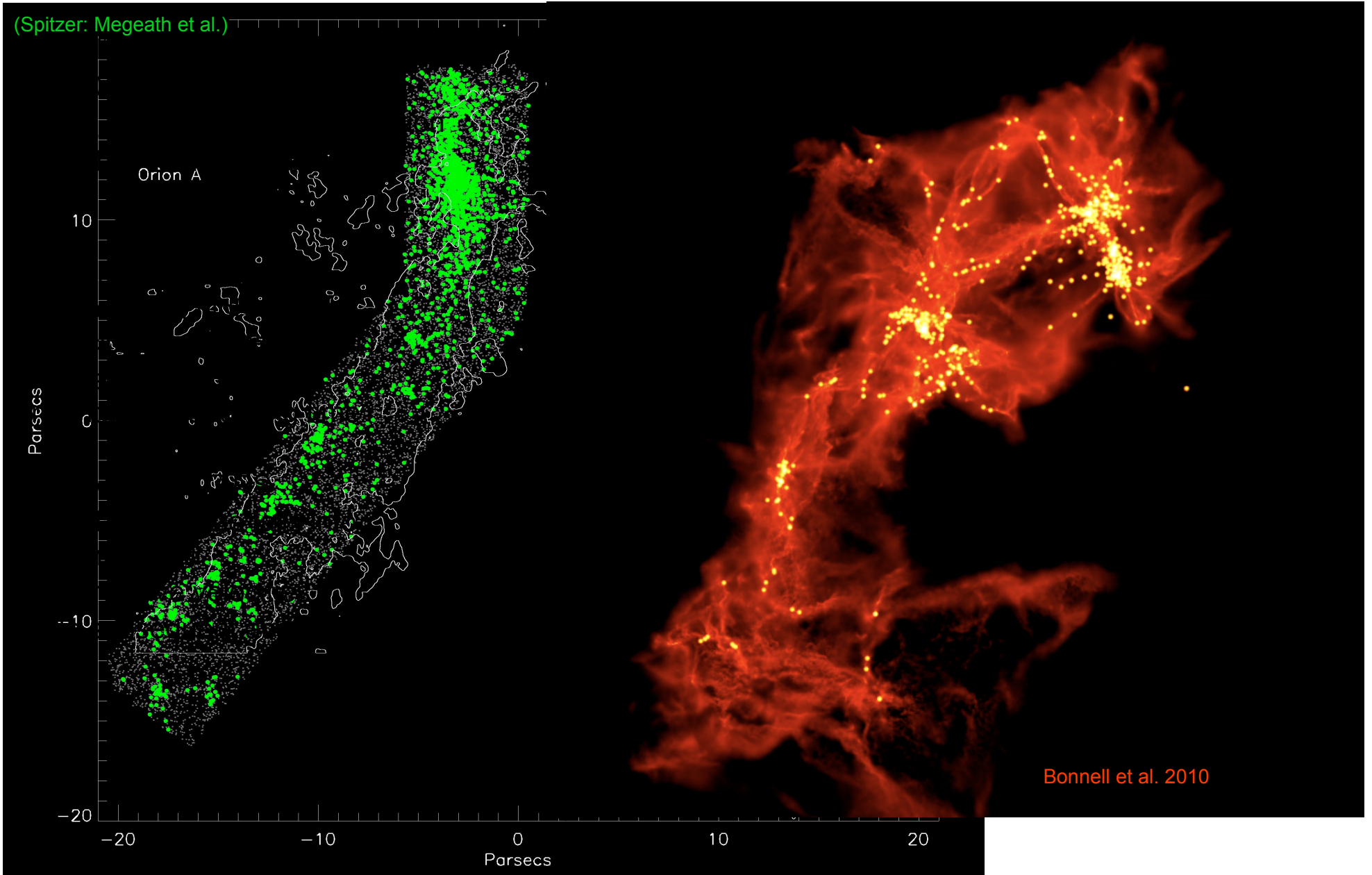
efficiency varies from 1% to 20%

develops full IMF
(distribution of sink particle masses)





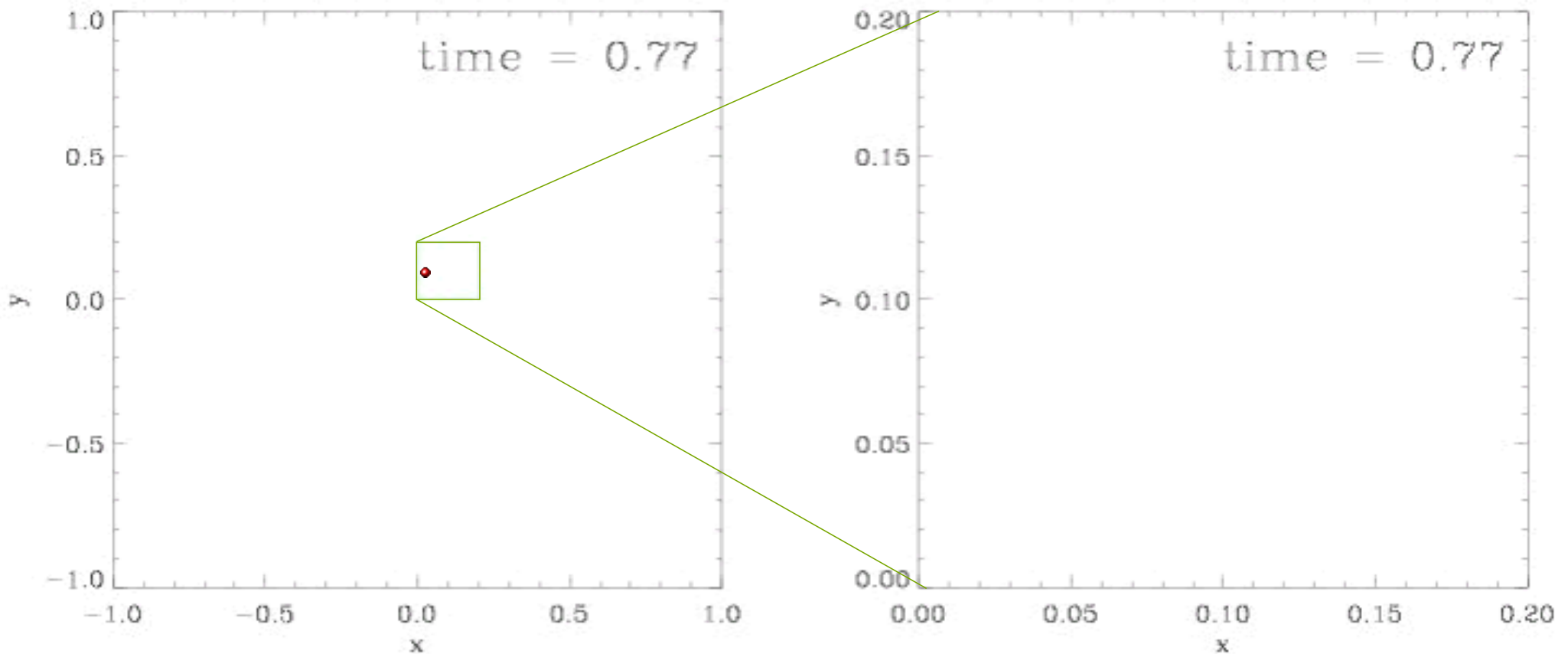
example: model of Orion cloud





dynamics of nascent star cluster

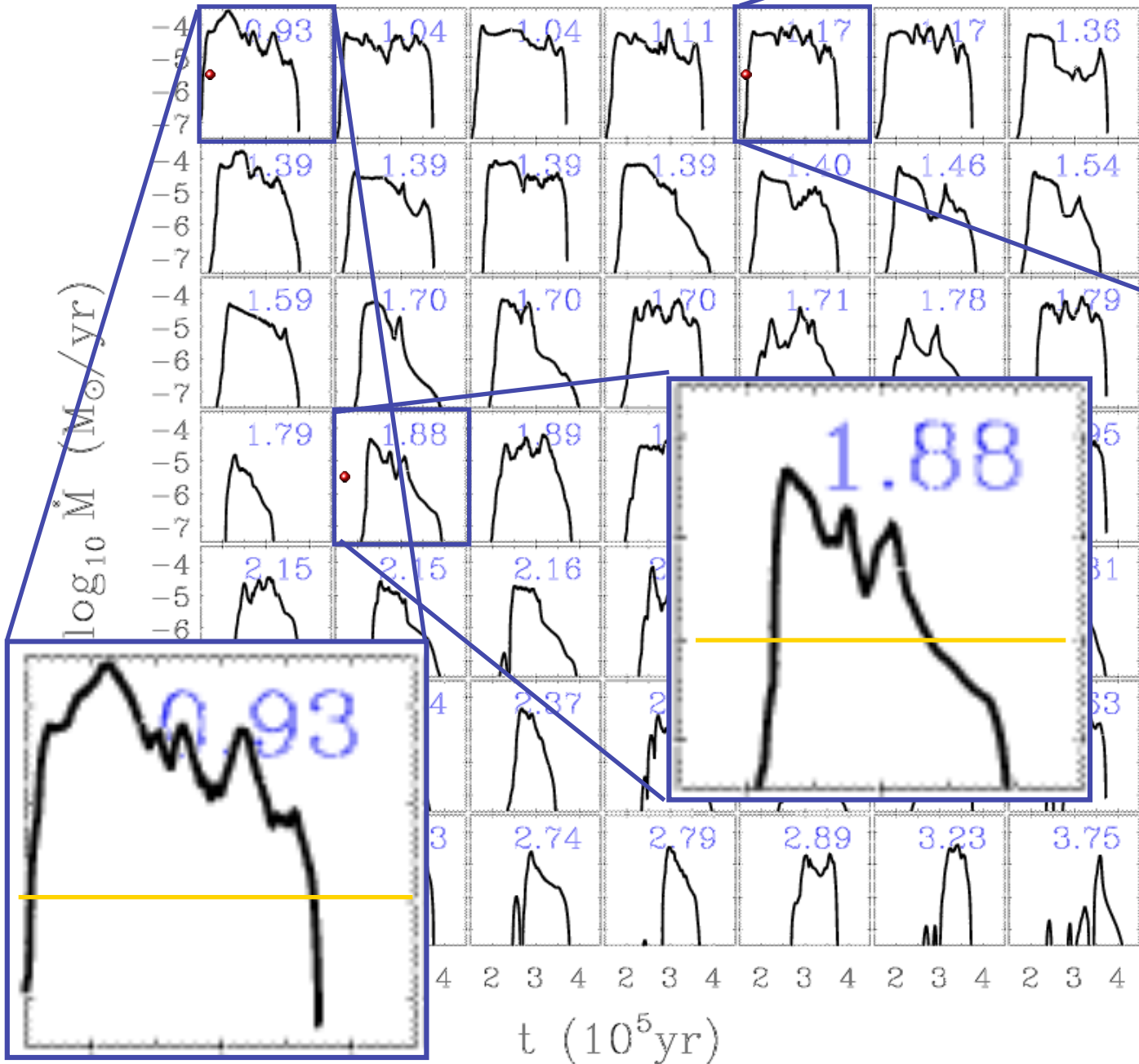
in dense clusters protostellar interaction may become important!



Trajectories of protostars in a nascent dense cluster created by gravoturbulent fragmentation
(from Klessen & Burkert 2000, *ApJS*, 128, 287)

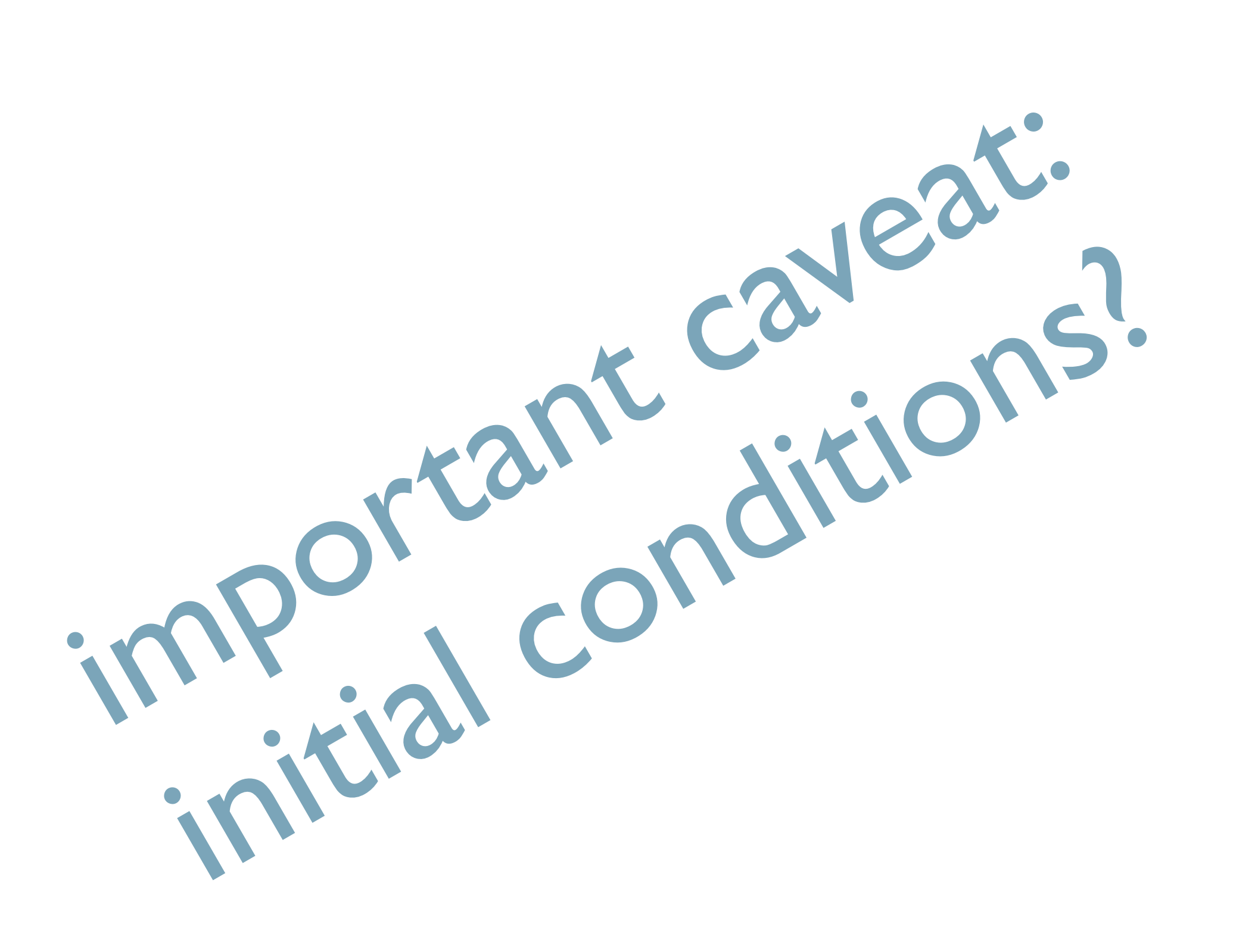


accretion rates in clust



Mass accretion rates *vary with time* and are strongly *influenced* by the *cluster environment*.

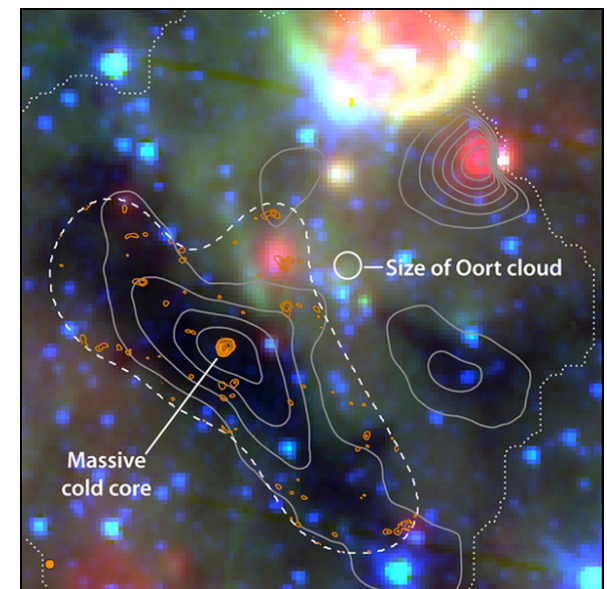
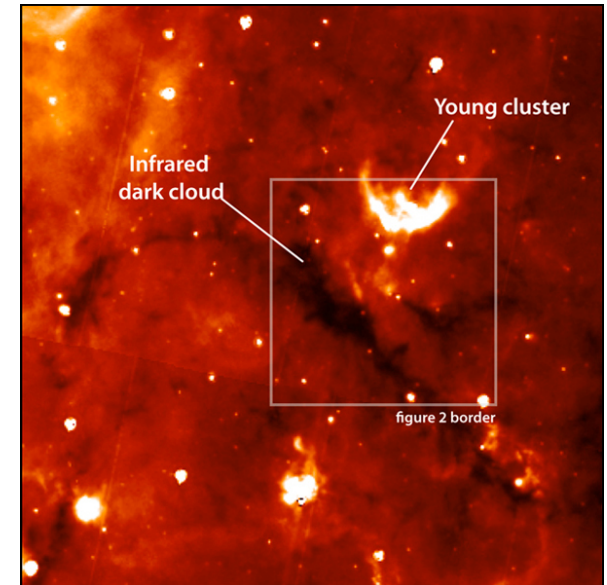
(Klessen 2001, ApJ, 550, L77; also Schmeja & Klessen, 2004, A&A, 419, 405)



important caveat:
initial conditions?

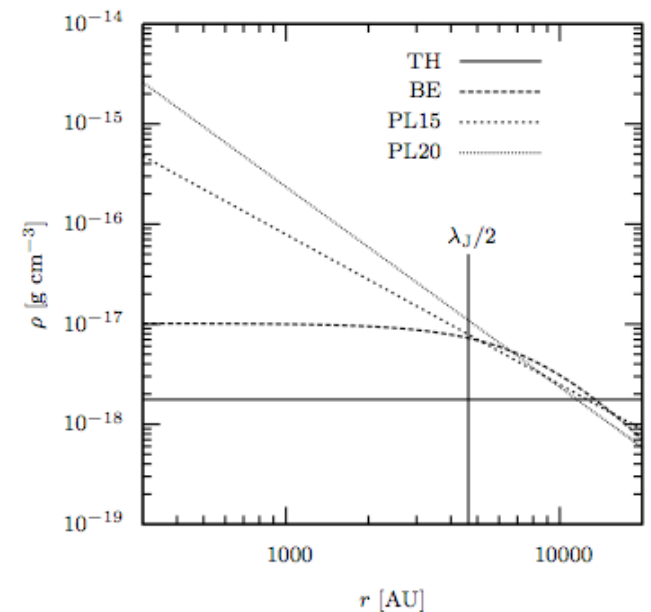
ICs of star cluster formation

- key question:
 - what is the initial density profile of cluster forming cores? how does it compare low-mass cores?
- observers answer:
 - very difficult to determine!
 - ▶ most high-mass cores have some SF inside
 - ▶ infra-red dark clouds (IRDCs) are difficult to study
 - but, new results with Herschel



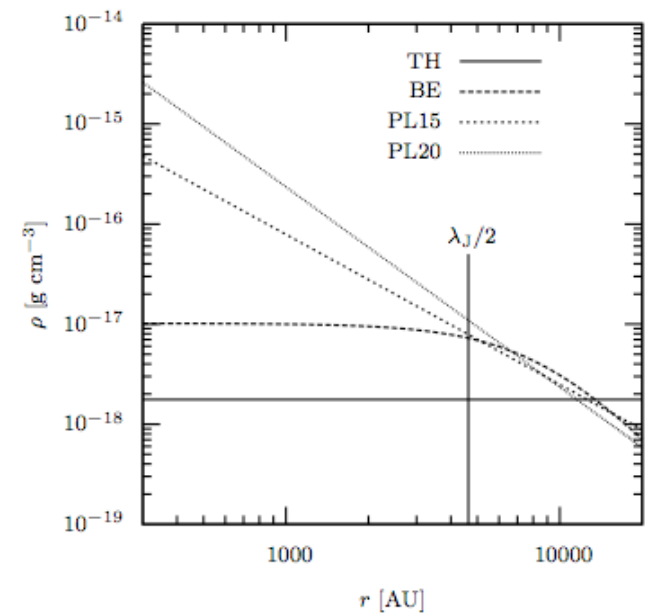
ICs of star cluster formation

- key question:
 - what is the initial density profile of cluster forming cores? how does it compare low-mass cores?
- theorists answer:
 - top hat (Larson Penston)
 - Bonnor Ebert (like low-mass cores)
 - power law $\rho \propto r^{-1}$ (logotrop)
 - power law $\rho \propto r^{-3/2}$ (Krumholz, McKee, et
 - power law $\rho \propto r^{-2}$ (Shu)
 - and many more



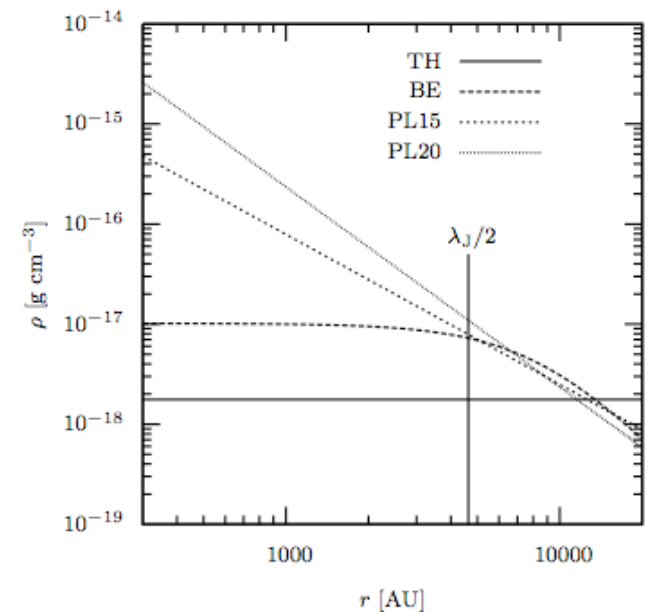
different density profiles

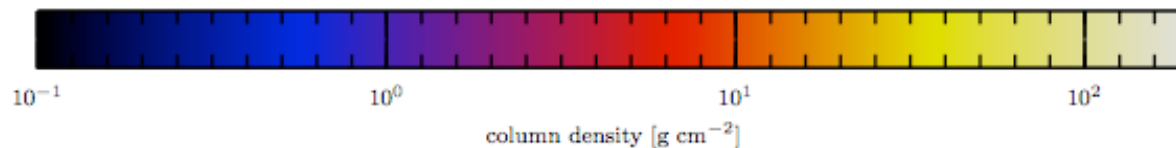
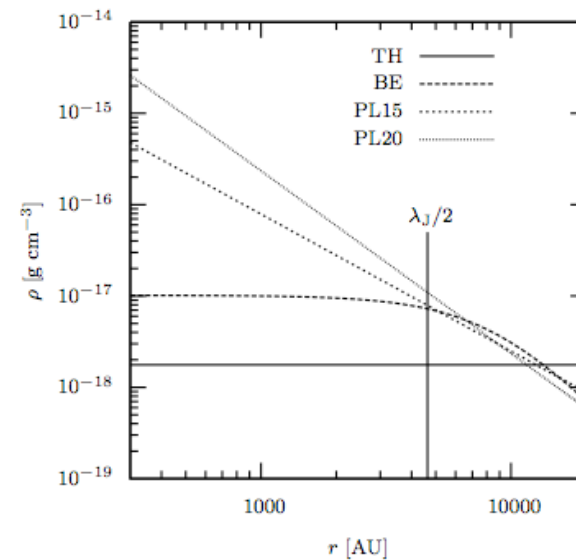
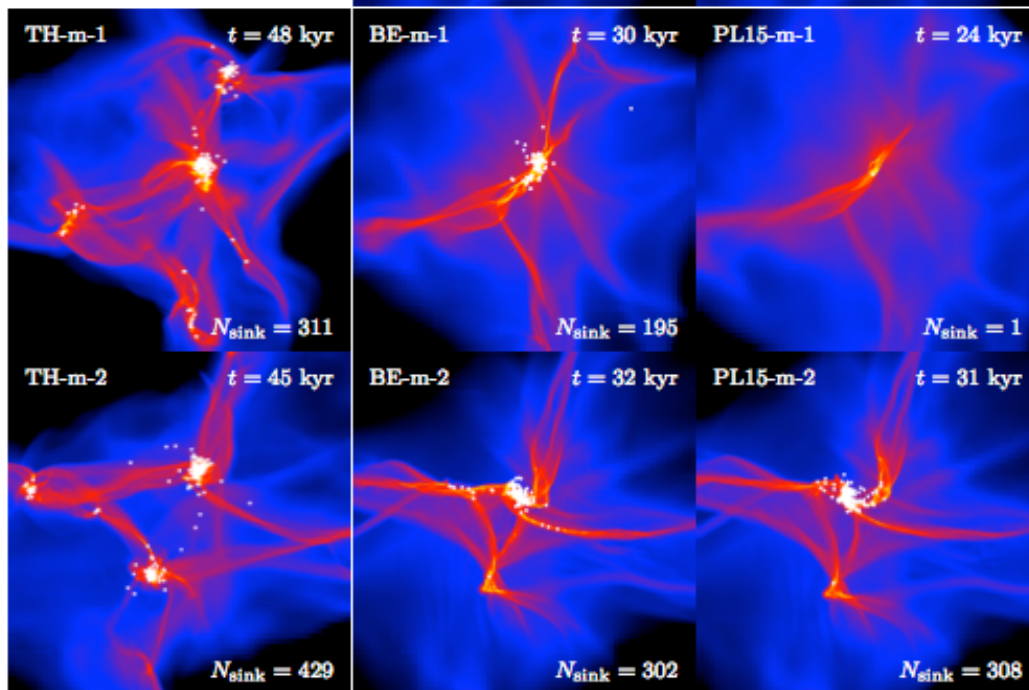
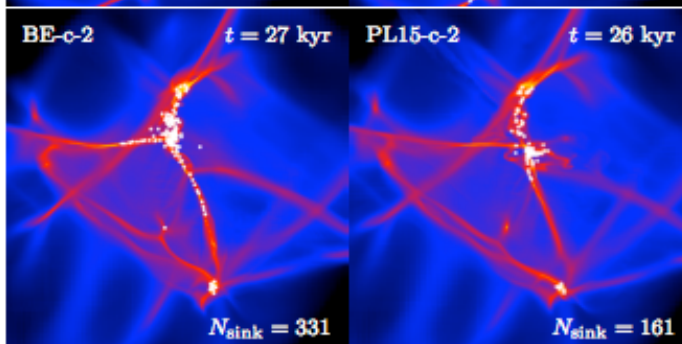
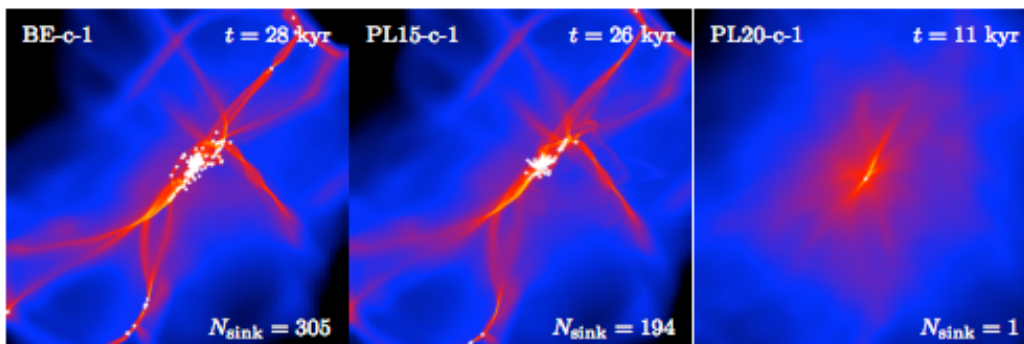
- does the density profile matter?
 -
 -
 -
- in comparison to
 - turbulence ...
 - radiative feedback ...
 - magnetic fields ...
 - thermodynamics ...



different density profiles

- address question in simple numerical experiment
- perform extensive parameter study
 - different profiles (top hat, BE, $r^{-3/2}$, r^{-3})
 - different turbulence fields
 - ▶ different realizations
 - ▶ different Mach numbers
 - ▶ solenoidal turbulence
 - ▶ dilatational turbulence
 - ▶ both modes
 - no net rotation, no B-fields (at the moment)



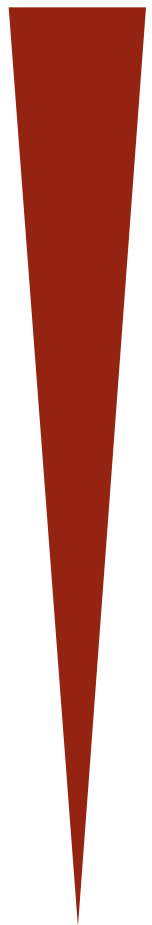


Run	t_{sim} [kyr]	$t_{\text{sim}}/t_{\text{ff}}^{\text{core}}$	$t_{\text{sim}}/t_{\text{ff}}$	N_{sinks}	$\langle M \rangle [M_{\odot}]$	M_{max}
TH-m-1	48.01	0.96	0.96	311	0.0634	0.86
TH-m-2	45.46	0.91	0.91	429	0.0461	0.74
BE-c-1	27.52	1.19	0.55	305	0.0595	0.94
BE-c-2	27.49	1.19	0.55	331	0.0571	0.97
BE-m-1	30.05	1.30	0.60	195	0.0873	1.42
BE-m-2	31.94	1.39	0.64	302	0.0616	0.54
BE-s-1	30.93	1.34	0.62	234	0.0775	1.14
BE-s-2	35.86	1.55	0.72	325	0.0587	0.51
PL15-c-1	25.67	1.54	0.51	194	0.0992	8.89
PL15-c-2	25.82	1.55	0.52	161	0.1244	12.3
PL15-m-1	23.77	1.42	0.48	1	20	20.0
PL15-m-2	31.10	1.86	0.62	308	0.0653	6.88
PL15-s-1	24.85	1.49	0.50	1	20	20.0
PL15-s-2	35.96	2.10	0.72	422	0.0478	4.50
PL20-c-1	10.67	0.92	0.21	1	20	20.0

ICs with flat inner density profile on average form more fragments

number of protostars

Run	t_{sim} [kyr]	$t_{\text{sim}}/t_{\text{ff}}^{\text{core}}$	$t_{\text{sim}}/t_{\text{ff}}$	N_{sinks}	$\langle M \rangle [M_{\odot}]$	M_{max}
TH-m-1	48.01	0.96	0.96	311	0.0634	0.86
TH-m-2	45.46	0.91	0.91	429	0.0461	0.74
BE-c-1	27.52	1.19	0.55	305	0.0595	0.94
BE-c-2	27.49	1.19	0.55	331	0.0571	0.97
BE-m-1	30.05	1.30	0.60	195	0.0873	1.42
BE-m-2	31.94	1.39	0.64	302	0.0616	0.54
BE-s-1	30.93	1.34	0.62	234	0.0775	1.14
BE-s-2	35.86	1.55	0.72	325	0.0587	0.51
PL15-c-1	25.67	1.54	0.51	194	0.0992	8.89
PL15-c-2	25.82	1.55	0.52	161	0.1244	12.3
PL15-m-1	23.77	1.42	0.48	1	20	20.0
PL15-m-2	31.10	1.86	0.62	308	0.0653	6.88
PL15-s-1	24.85	1.49	0.50	1	20	20.0
PL15-s-2	35.96	2.10	0.72	422	0.0478	4.50
PL20-c-1	10.67	0.92	0.21	1	20	20.0



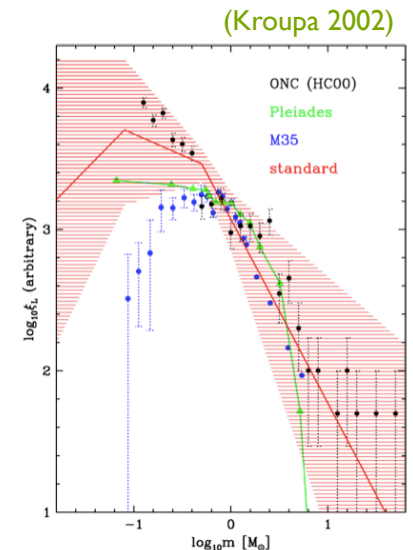
ICs with flat inner density profile on average form more fragments

however, the real situation is very complex: details of the initial turbulent field matter

number of protostars

stellar mass function

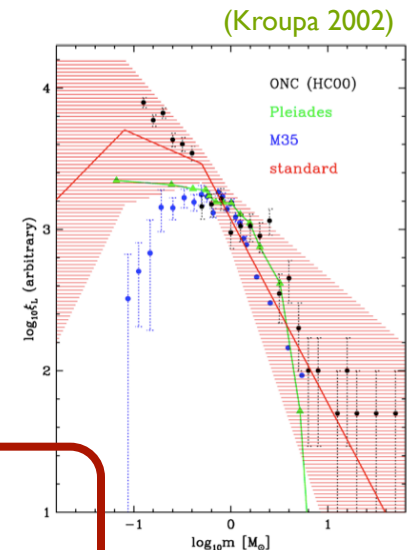
- distribution of stellar masses depends on
 - turbulent initial conditions
 - > mass spectrum of prestellar cloud cores
 - collapse and interaction of prestellar cores
 - > accretion and N -body effects
 - thermodynamic properties of gas
 - > balance between heating and cooling
 - > EOS (determines which cores go into collapse)
 - (proto) stellar feedback terminates star formation
 - ionizing radiation, bipolar outflows, winds, SN, etc.



stellar mass function

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application to early star formation



thermodynamics & fragmentation

degree of fragmentation depends on *EOS!*

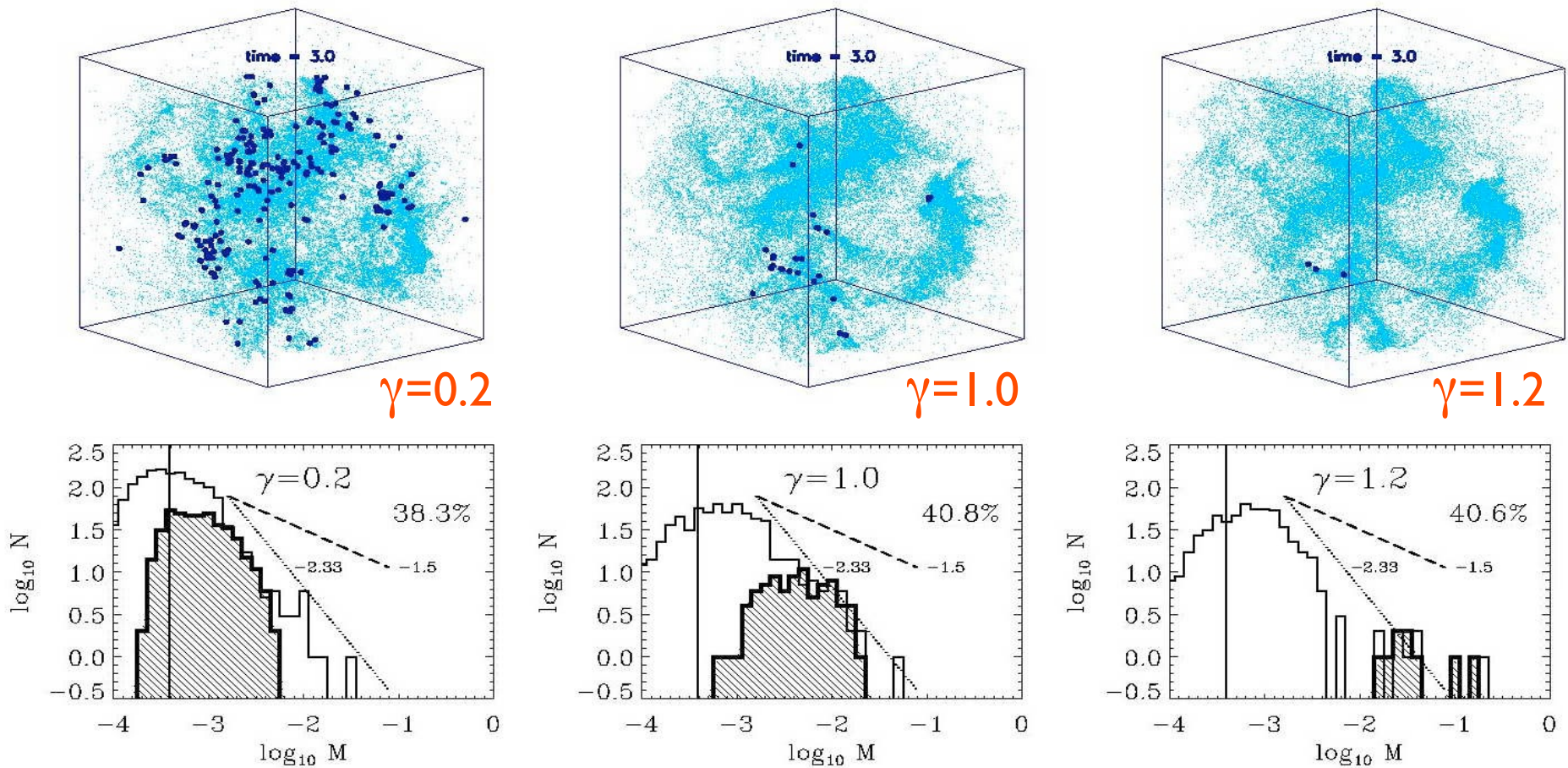
polytropic EOS: $p \propto \rho^\gamma$

$\gamma < 1$: dense cluster of low-mass stars

$\gamma > 1$: isolated high-mass stars

(see Li et al. 2003; also Kawachi & Hanawa 1998, Larson 2003)

dependency on EOS

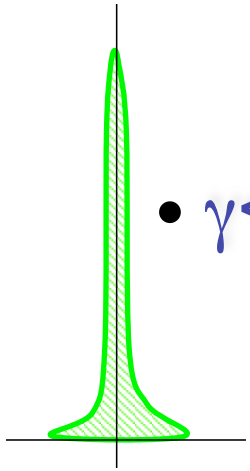


for $\gamma < 1$ fragmentation is enhanced \rightarrow *cluster of low-mass stars*
for $\gamma > 1$ it is suppressed \rightarrow *isolated massive stars*

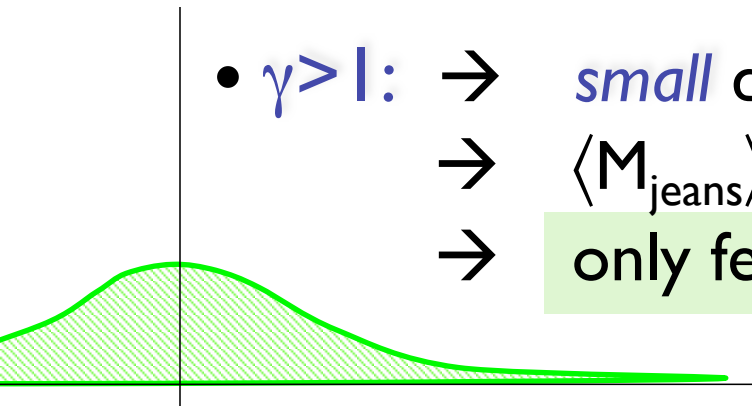
how does that work?

$$(1) \quad p \propto \rho^\gamma \quad \rightarrow \quad \rho \propto p^{1/\gamma}$$

$$(2) \quad M_{\text{jeans}} \propto \gamma^{3/2} \rho^{(3\gamma-4)/2}$$

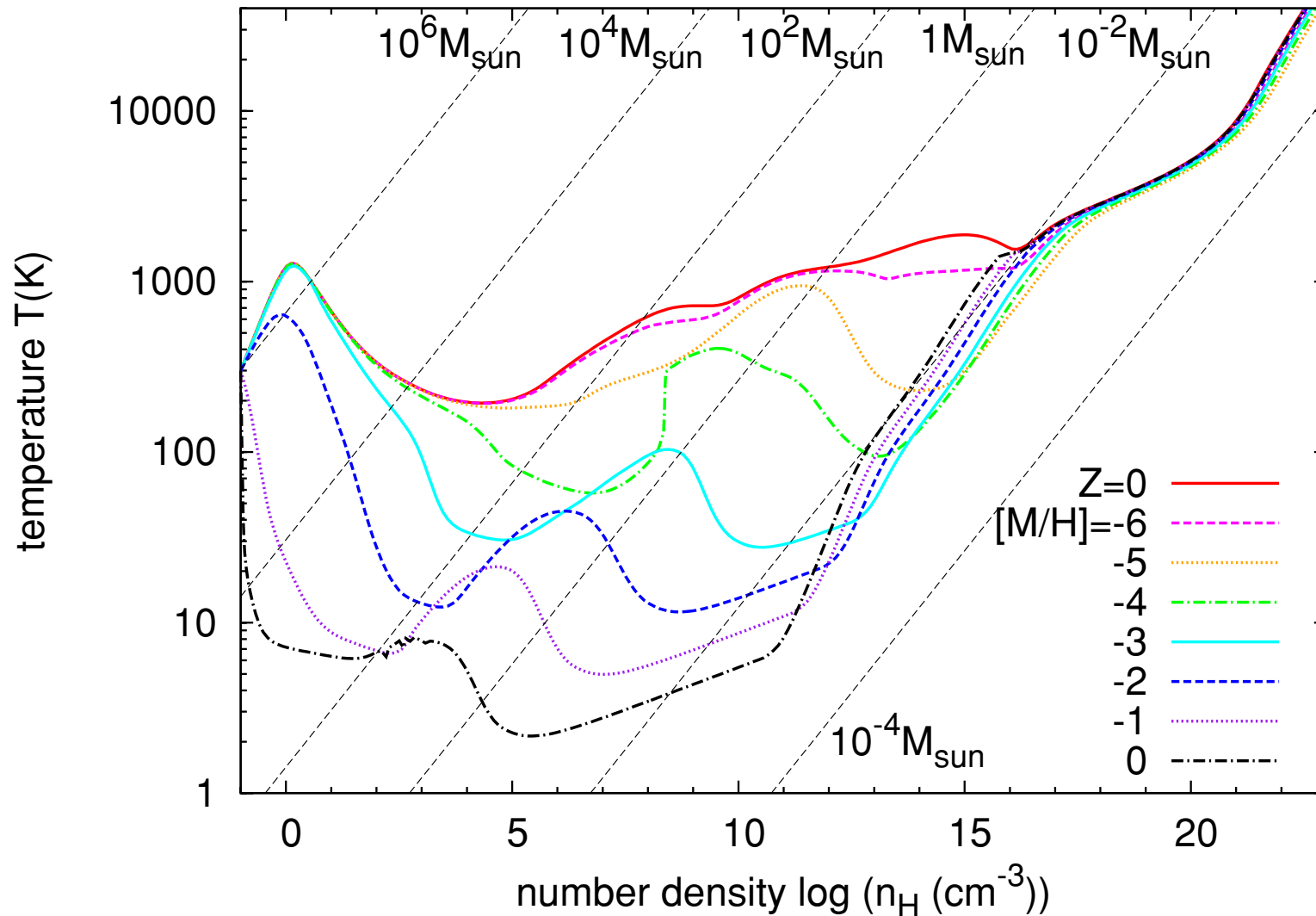


- $\gamma < 1$: \rightarrow *large* density excursion for given pressure
 \rightarrow $\langle M_{\text{jeans}} \rangle$ becomes small
 \rightarrow number of fluctuations with $M > M_{\text{jeans}}$ is large



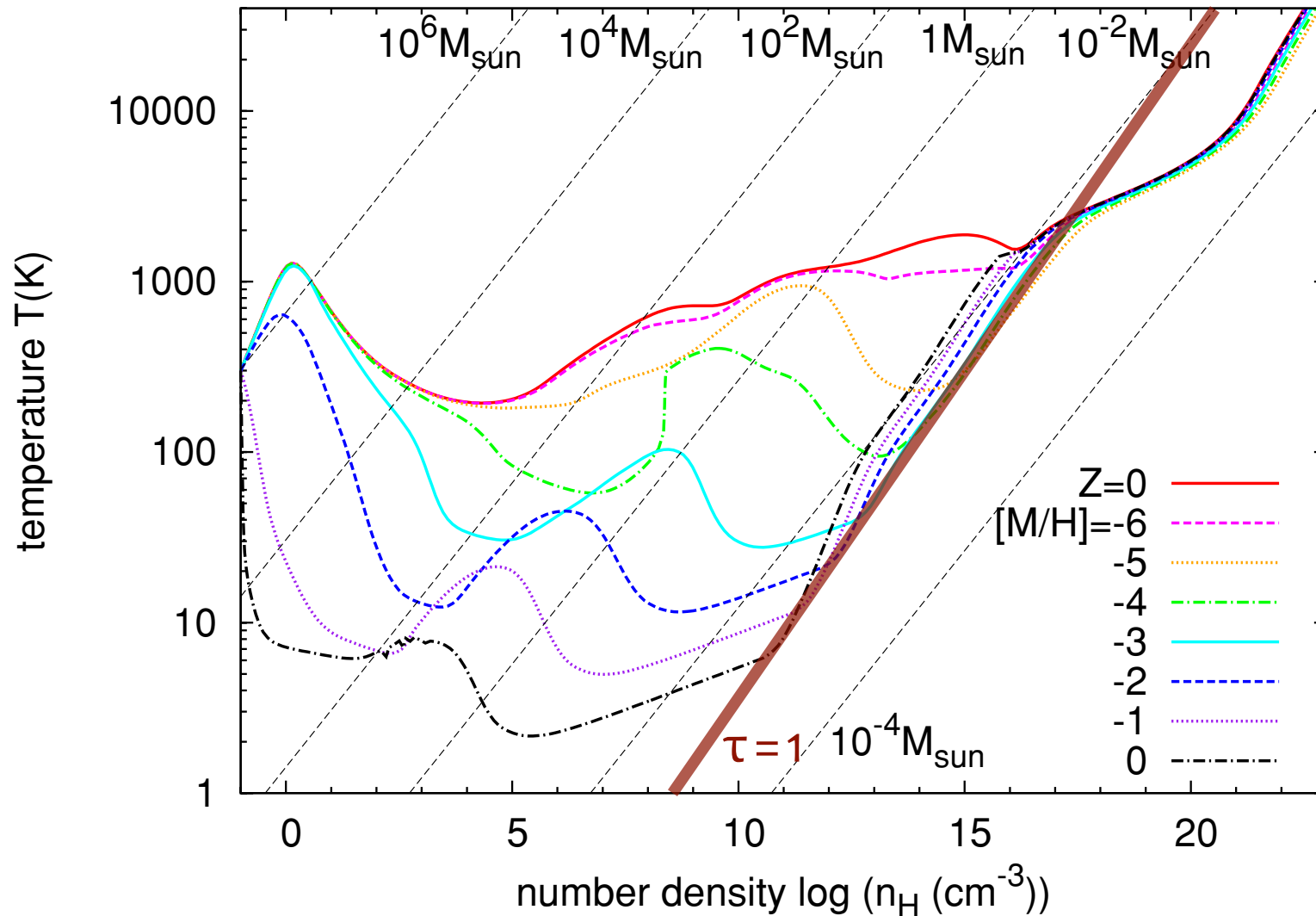
- $\gamma > 1$: \rightarrow *small* density excursion for given pressure
 \rightarrow $\langle M_{\text{jeans}} \rangle$ is large
 \rightarrow only few and massive clumps exceed M_{jeans}

EOS as function of metallicity



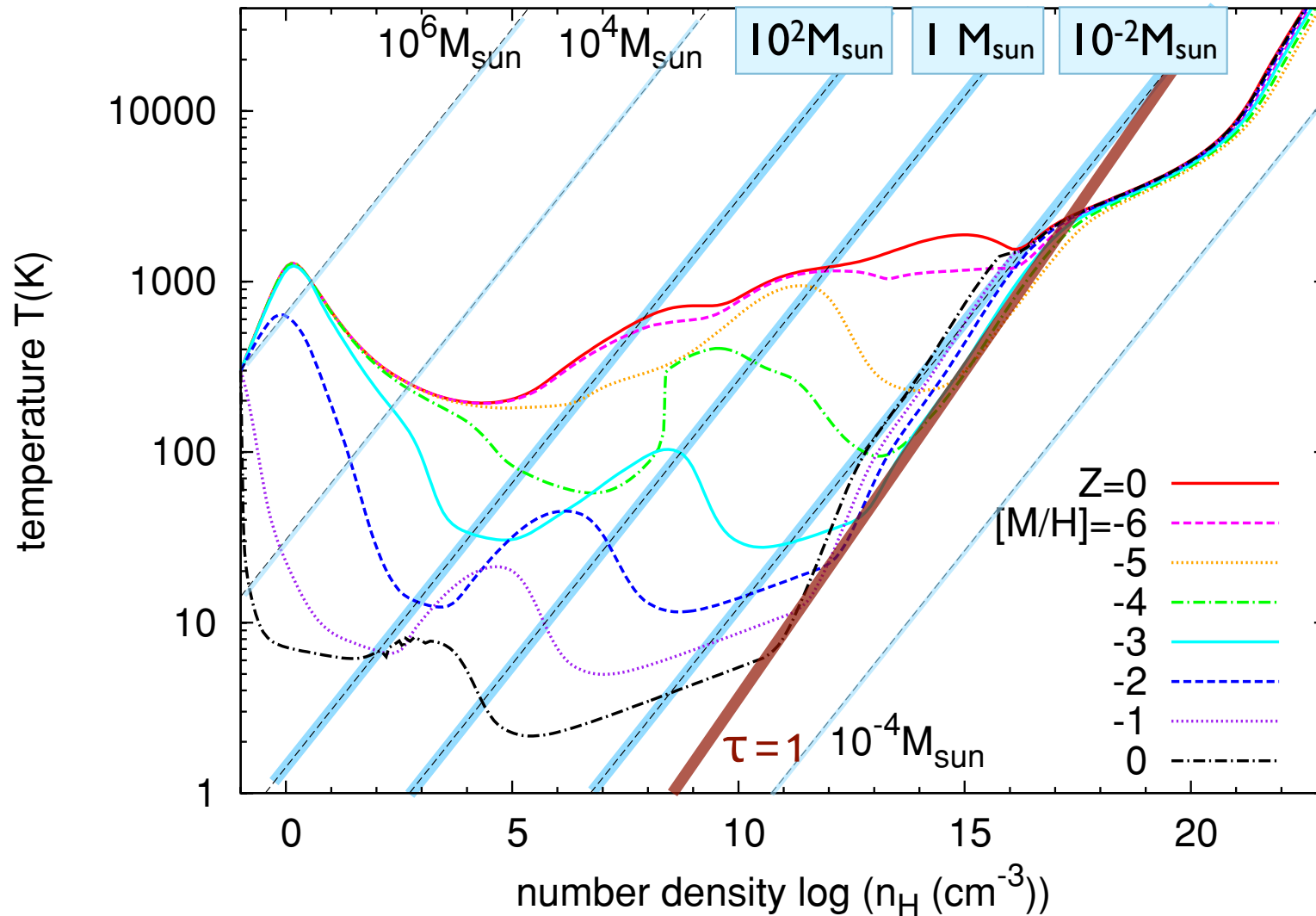
(Omukai et al. 2005, 2010)

EOS as function of metallicity



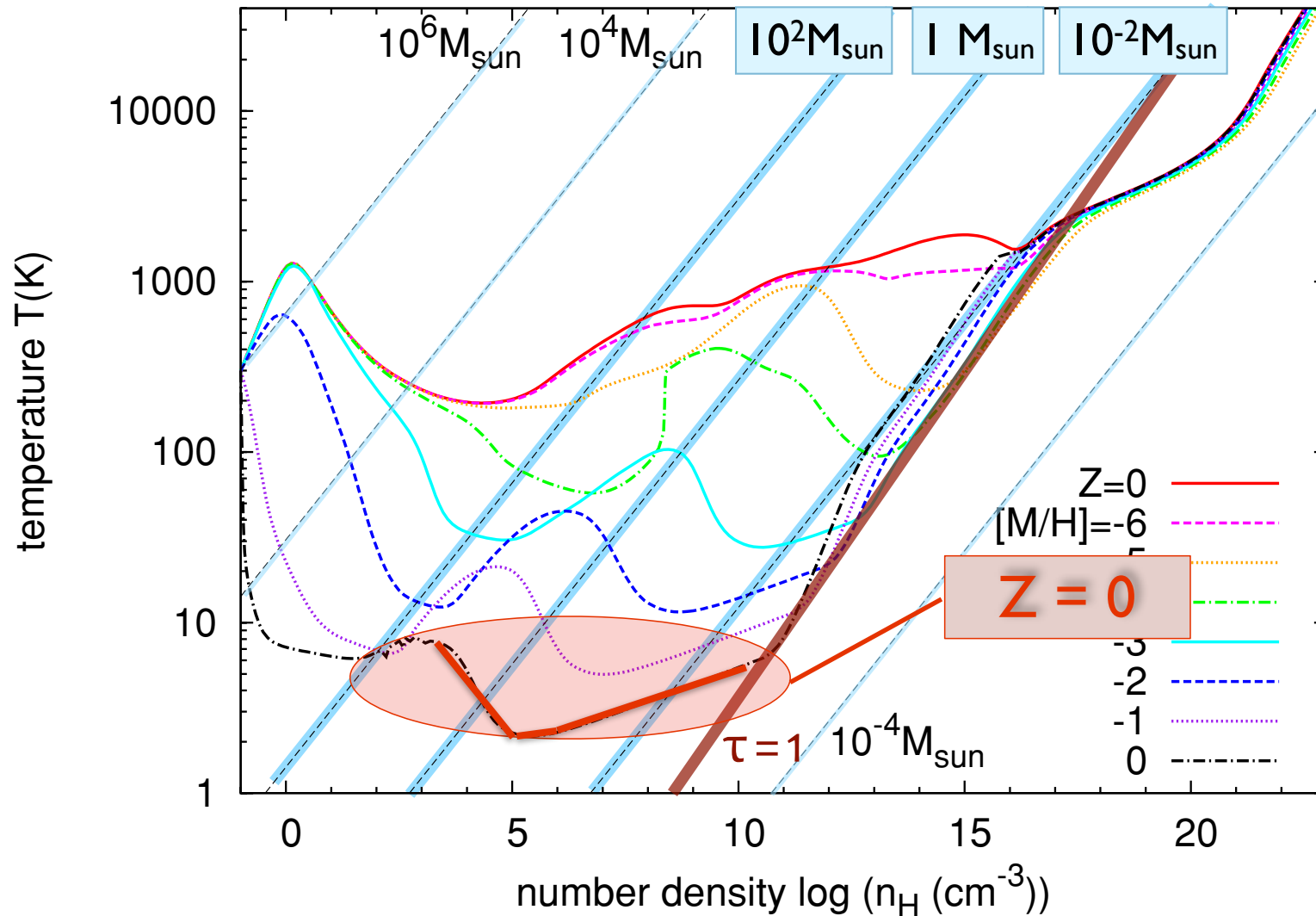
(Omukai et al. 2005, 2010)

EOS as function of metallicity



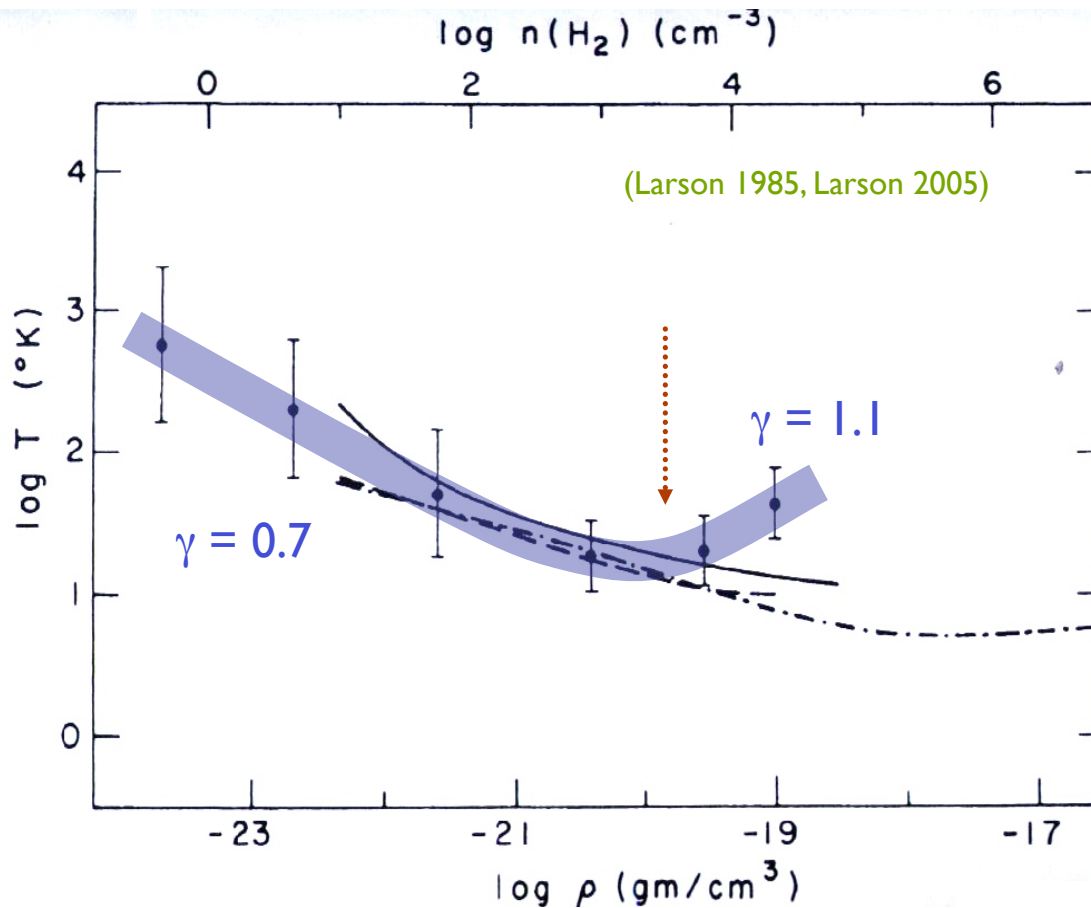
(Omukai et al. 2005, 2010)

EOS as function of metallicity

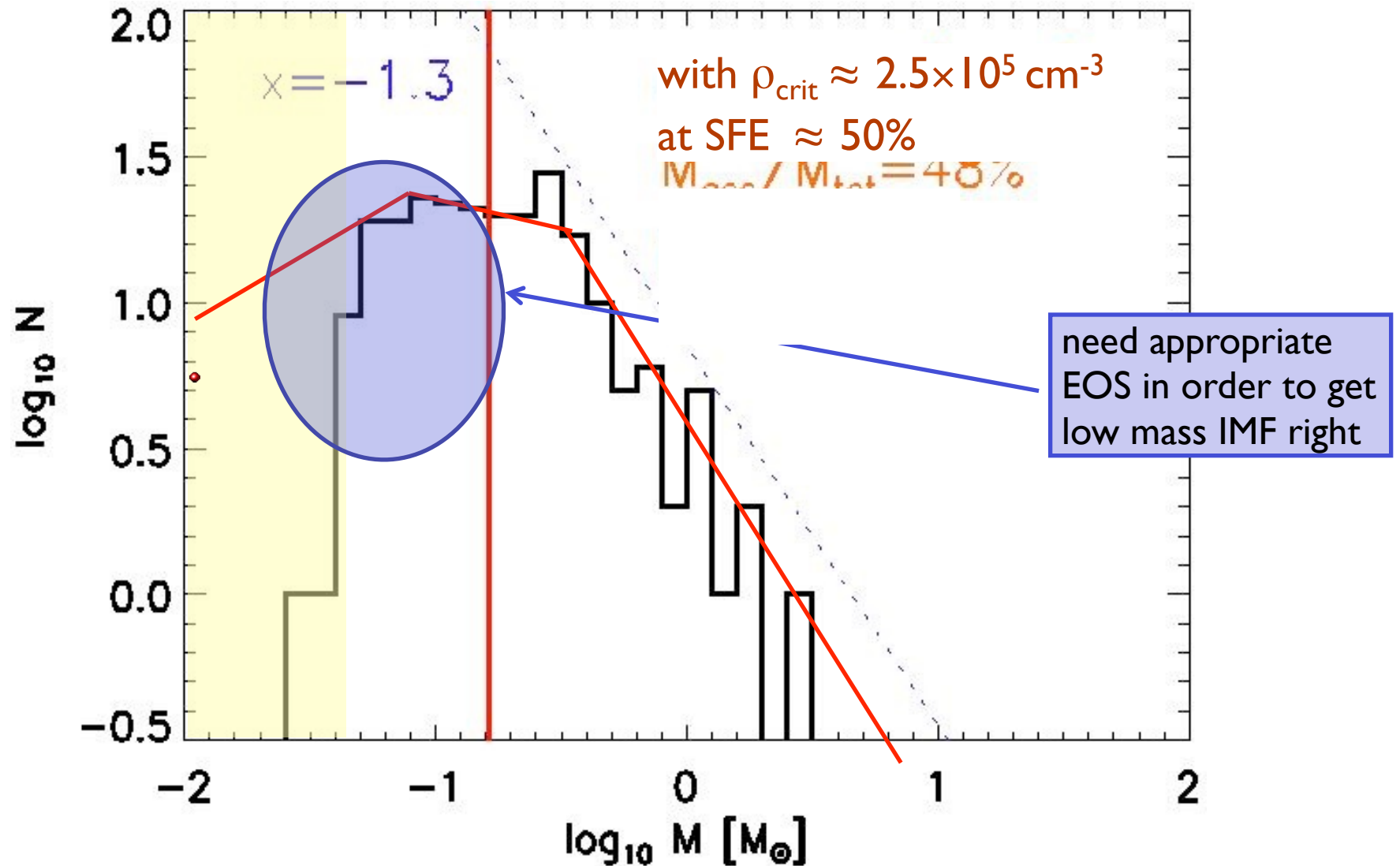


(Omukai et al. 2005, 2010)

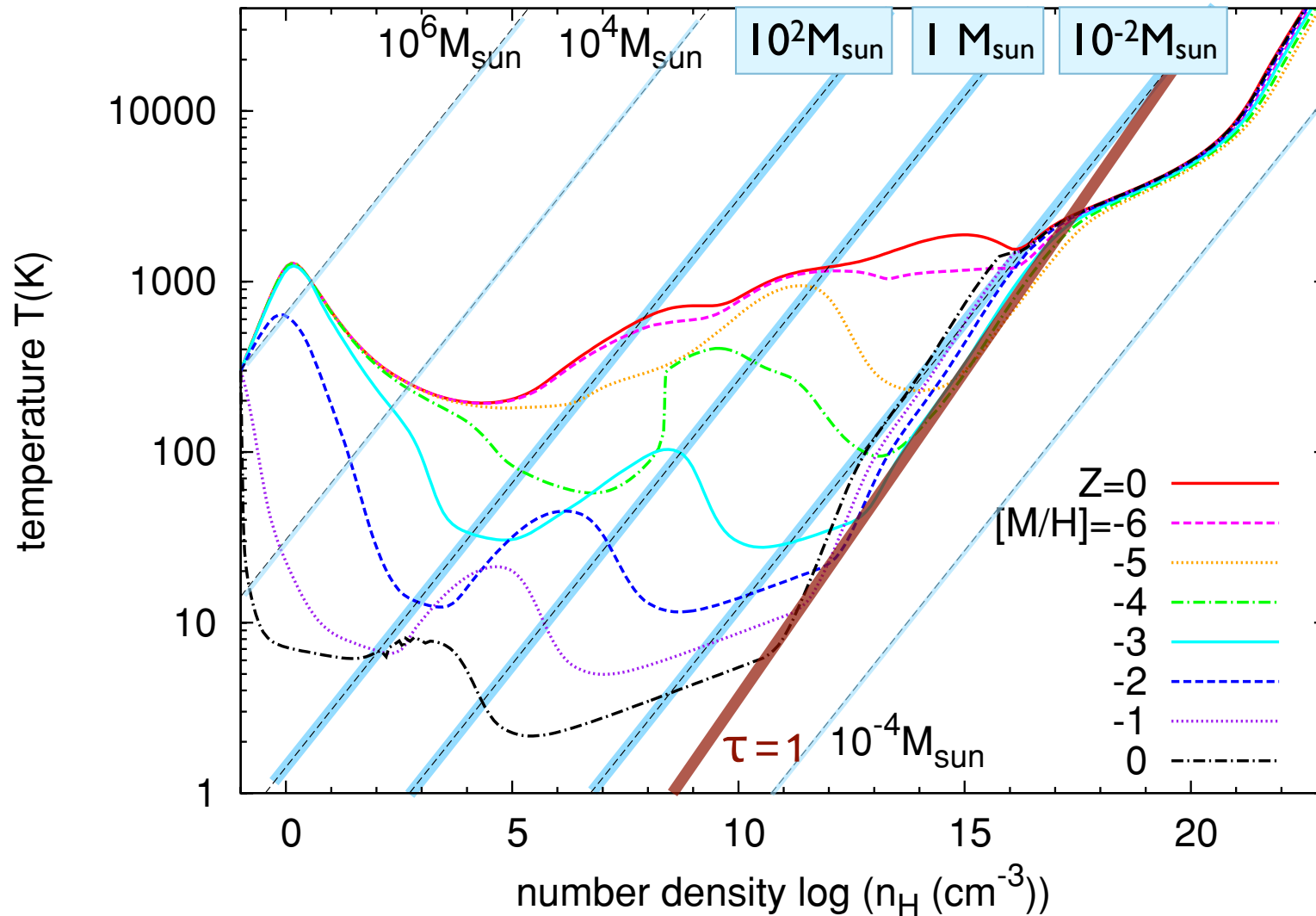
present-day star formation



IMF in nearby molecular clouds

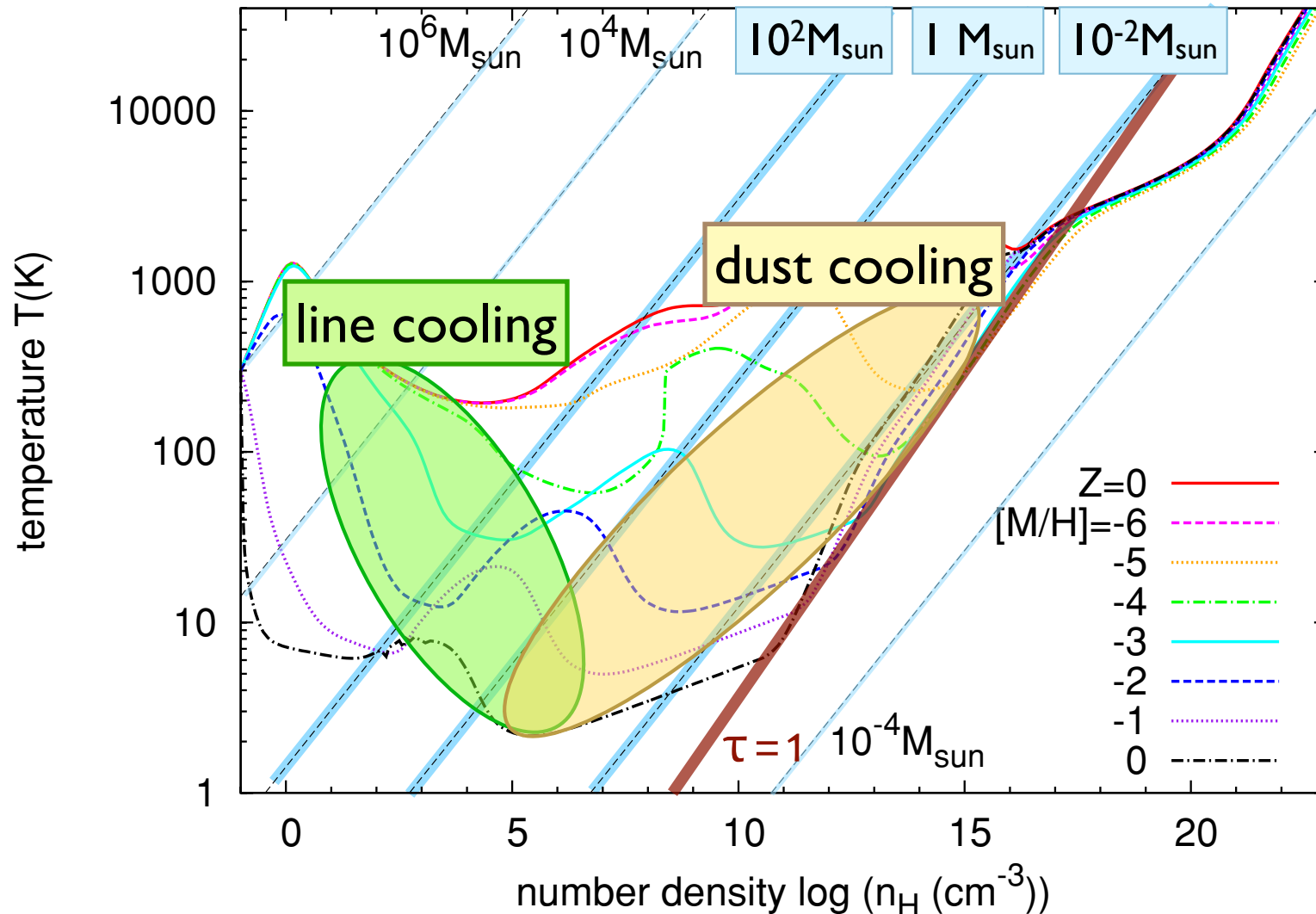


EOS as function of metallicity



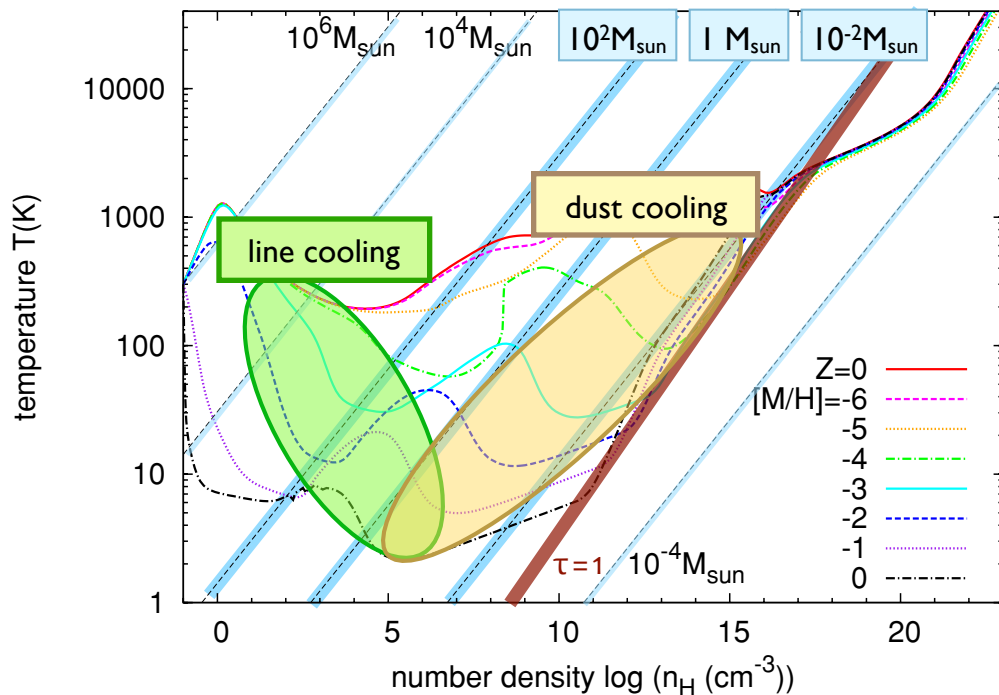
(Omukai et al. 2005, 2010)

EOS as function of metallicity



(Omukai et al. 2005, 2010)

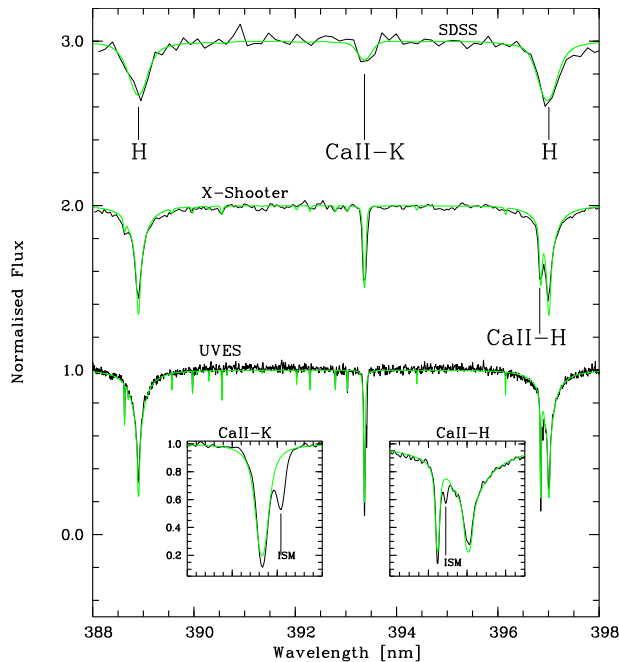
transition: Pop III to Pop II.5



two competing models:

- cooling due to atomic fine-structure lines ($Z > 10^{-3.5} Z_{\text{sun}}$)
- cooling due to coupling between gas and dust ($Z > 10^{-5...-6} Z_{\text{sun}}$)
- which one explains origin of extremely metal-poor stars?
NB: lines would only make very massive stars, with $M > \text{few} \times 10 M_{\text{sun}}$.

transition: Pop III to Pop II.5



SDSS J1029151+172927

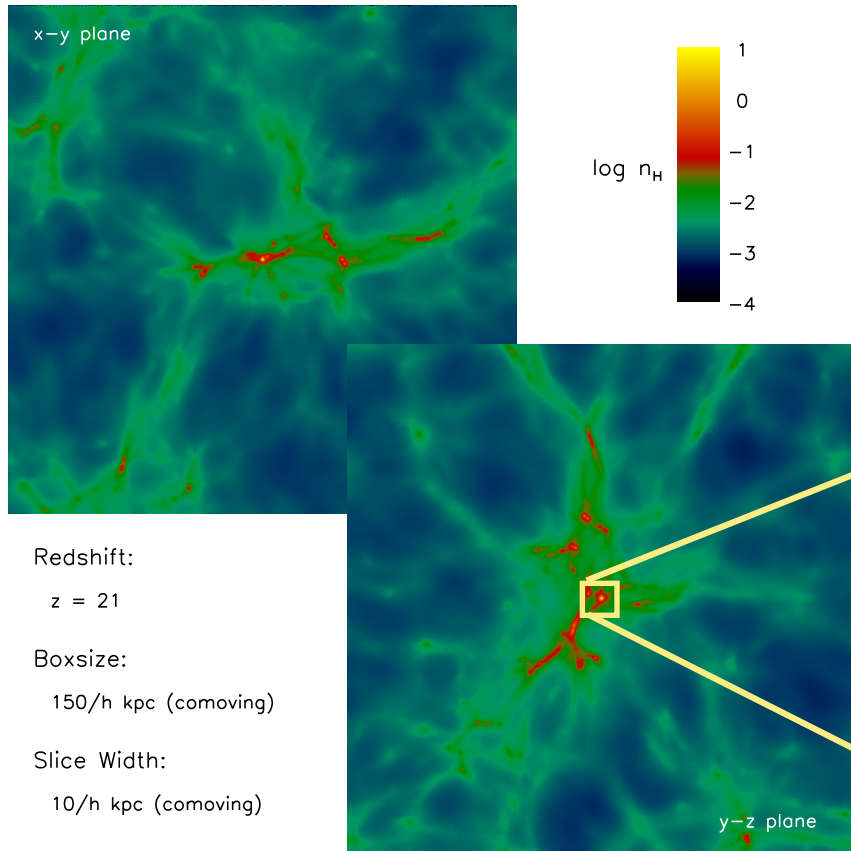
- is first ultra metal-poor star with $Z \sim 10^{-4.5} Z_{\text{sun}}$ for all metals seen (Fe, C, N, etc.)
[see Caffau et al. 2011]
- this is in regime, where metal-lines cannot provide cooling
[e.g. Schneider et al. 2011, 2012, Klessen et al. 2012]

- TOPoS ESO large program to find more of these stars (120h x-shooter, 30h UVES)

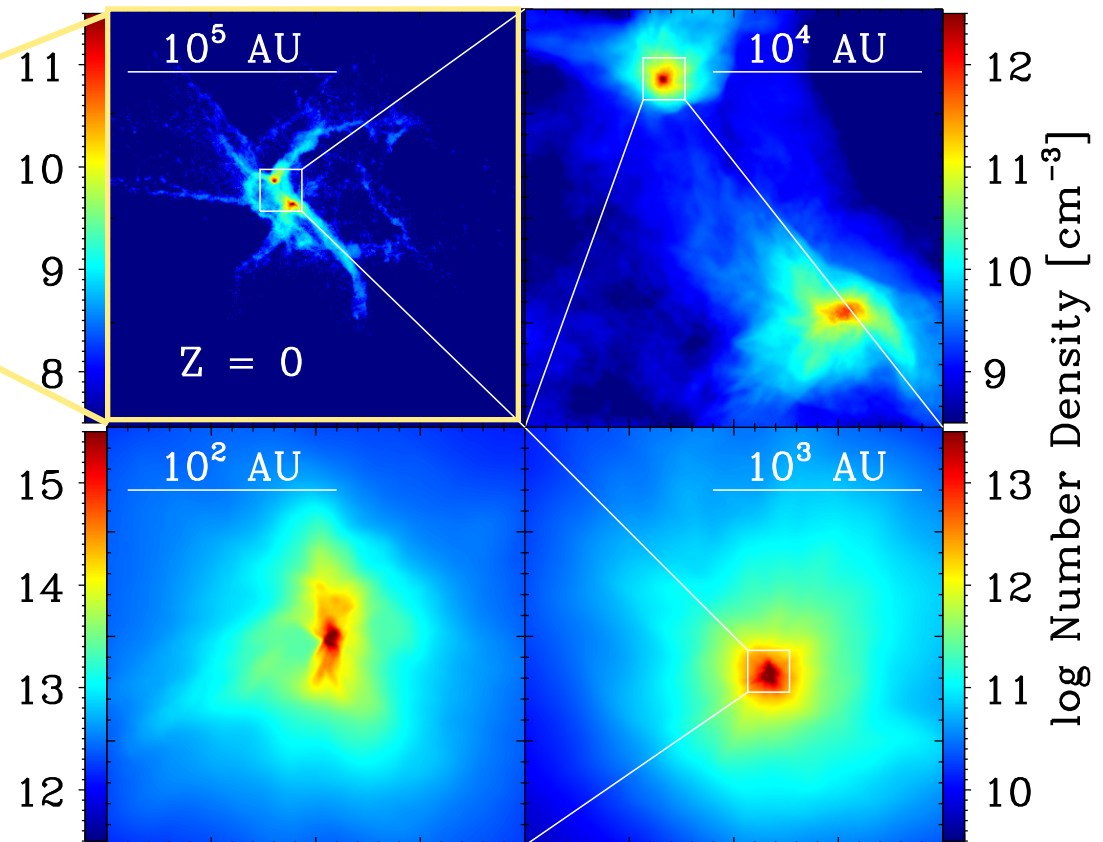
(Caffau et al. 2013, A&A, 560,A71, Bonifacio et al. 2014, in prep)

Element		+3Dcor.	[X/H] _{1D} +NLTE cor.	+ 3D cor + NLTE cor	N lines	S _H	A(X) _⊙
C	≤ -3.8	≤ -4.5			G-band		8.50
N	≤ -4.1	≤ -5.0			NH-band		7.86
Mg I	-4.71 ± 0.11	-4.68 ± 0.11	-4.52 ± 0.11	-4.49 ± 0.12	5	0.1	7.54
Si I	-4.27	-4.30	-3.93	-3.96	1	0.1	7.52
Ca I	-4.72	-4.82	-4.44	-4.54	1	0.1	6.33
Ca II	-4.81 ± 0.11	-4.93 ± 0.03	-5.02 ± 0.02	-5.15 ± 0.09	3	0.1	6.33
Ti II	-4.75 ± 0.18	-4.83 ± 0.16	-4.76 ± 0.18	-4.84 ± 0.16	6	1.0	4.90
Fe I	-4.73 ± 0.13	-5.02 ± 0.10	-4.60 ± 0.13	-4.89 ± 0.10	43	1.0	7.52
Ni I	-4.55 ± 0.14	-4.90 ± 0.11			10		6.23
Sr II	≤ -5.10	≤ -5.25	≤ -4.94	≤ -5.09	1	0.01	2.92

modeling the formation of the first/second stars

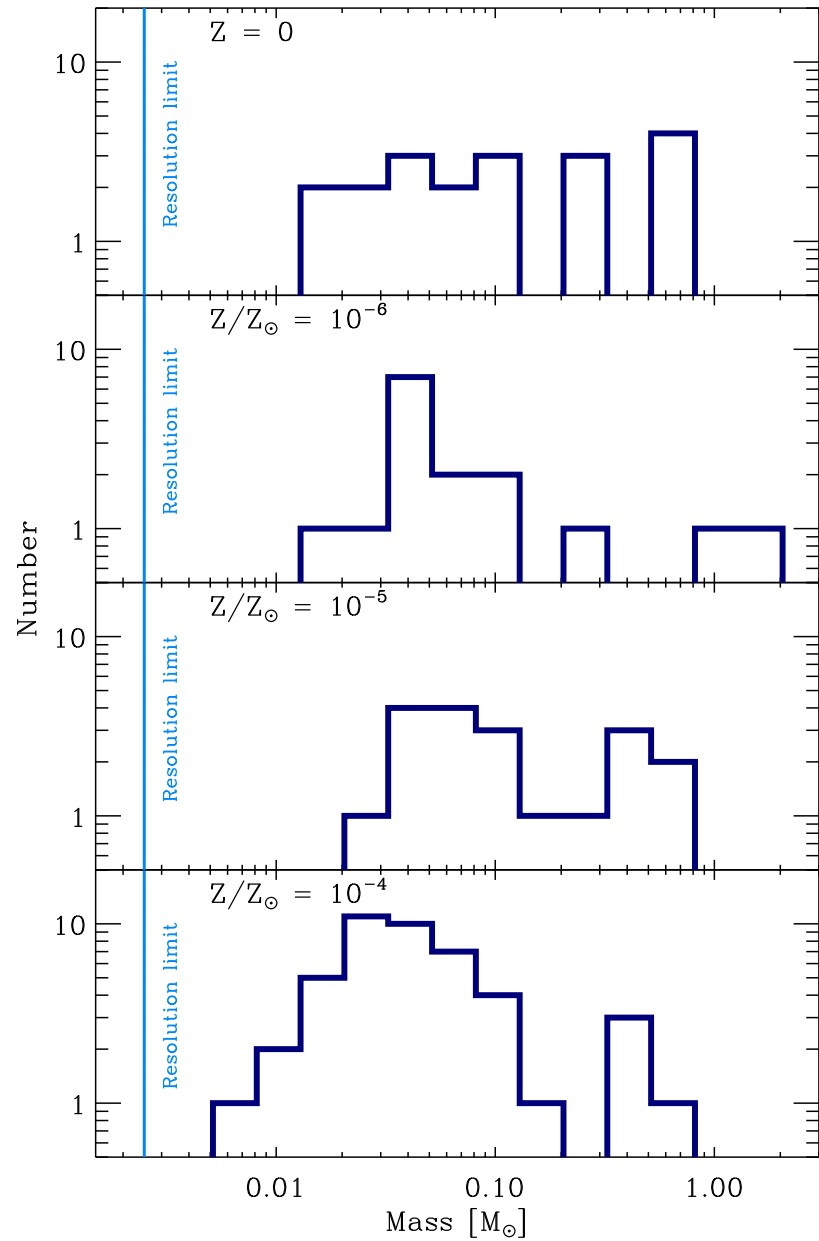
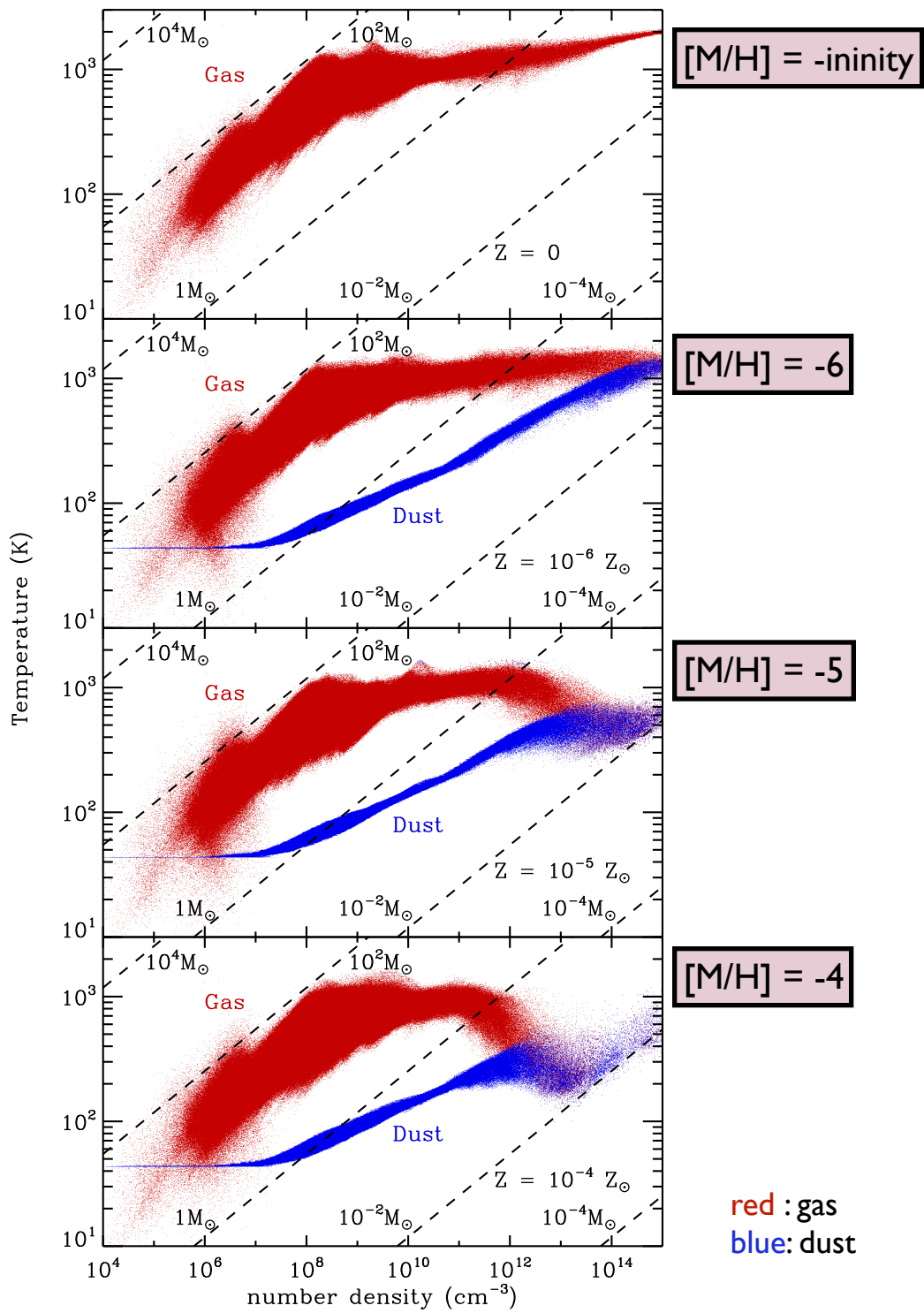


successive zoom-in calculation from cosmological initial conditions (using SPH and new grid-code AREPO)



(Greif et al., 2007, ApJ, 670, 1)

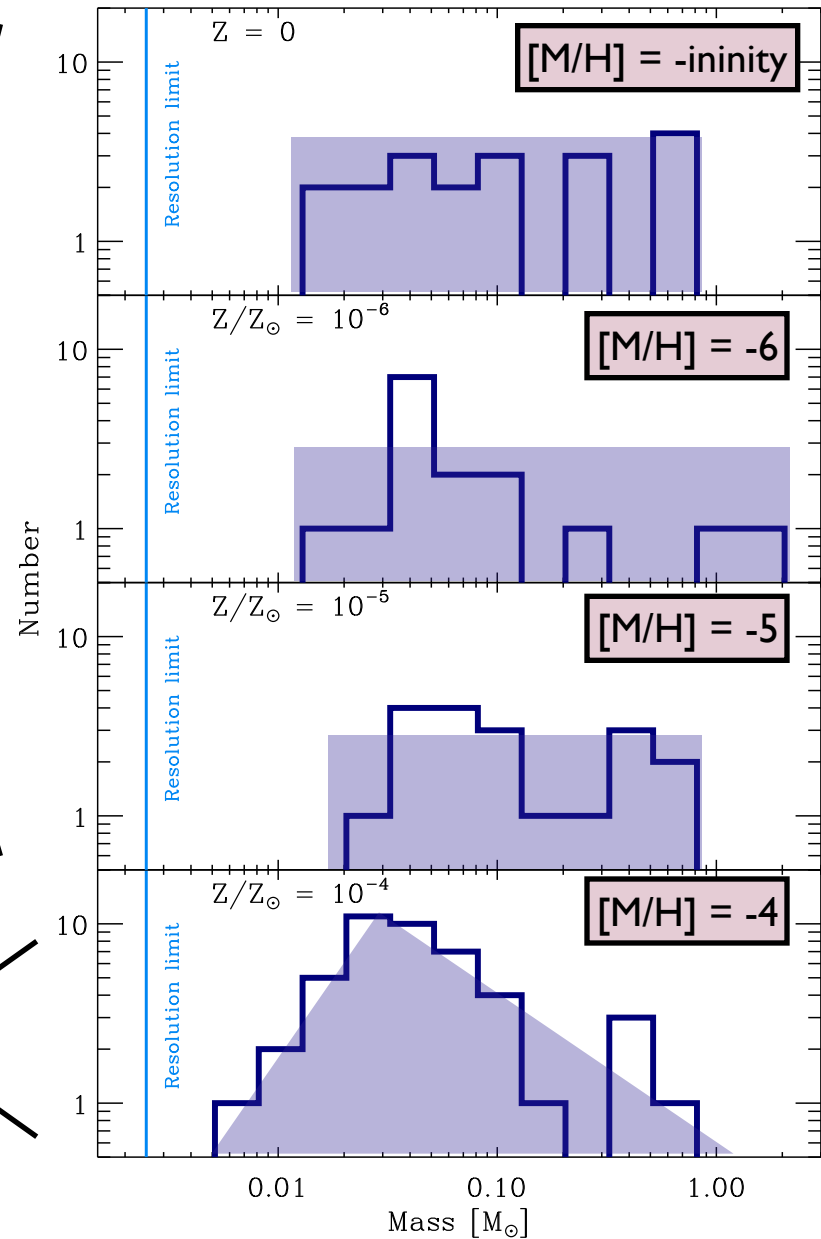
(Greif et al. 2011, ApJ, 737, 75, Greif et al. 2012, MNRAS, 424, 399, Dopcke et al. 2013, ApJ, 776, 103)



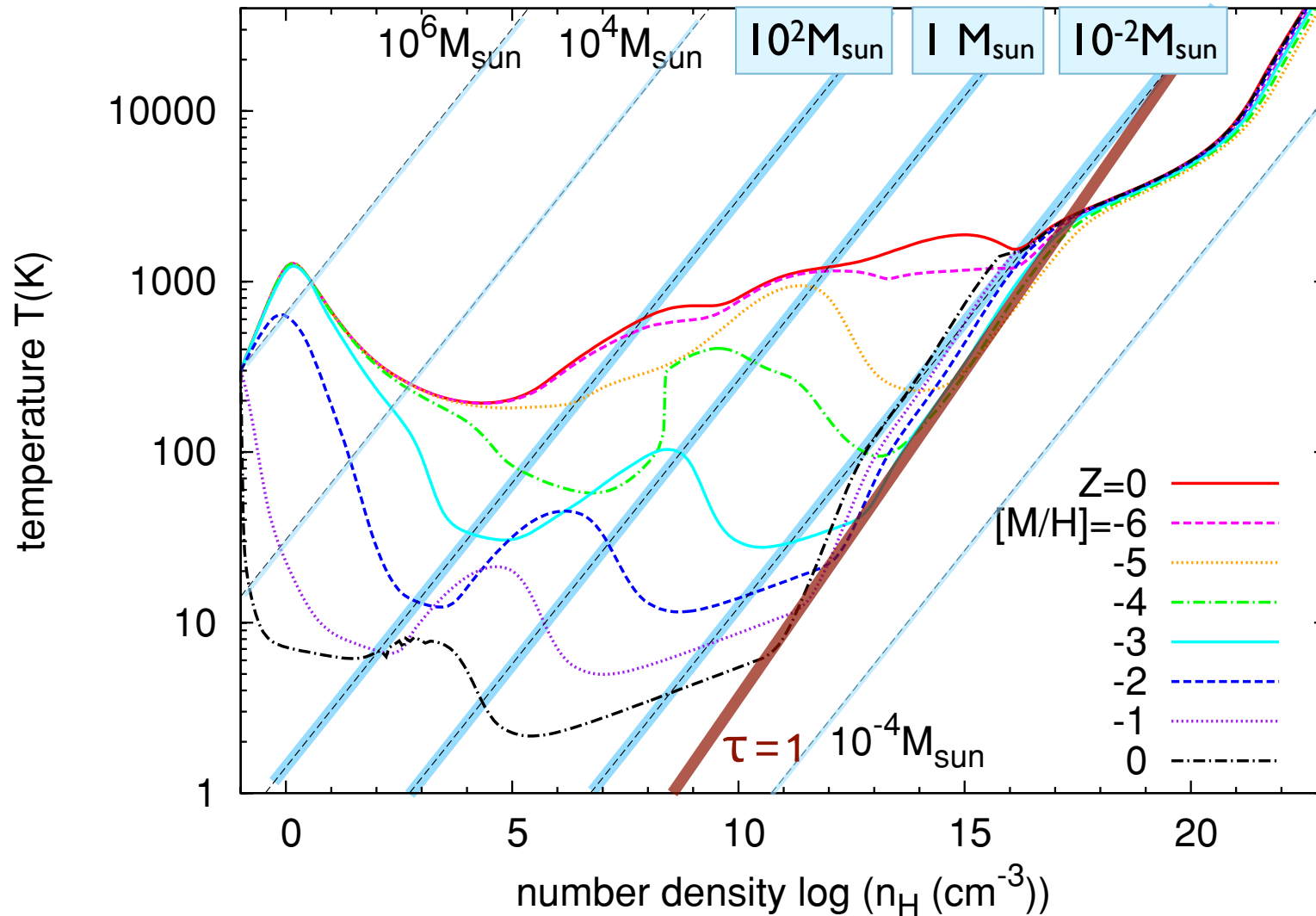
hints for differences
in mass spectrum

disk fragmentation mode

gravoturbulent fragmentation mode

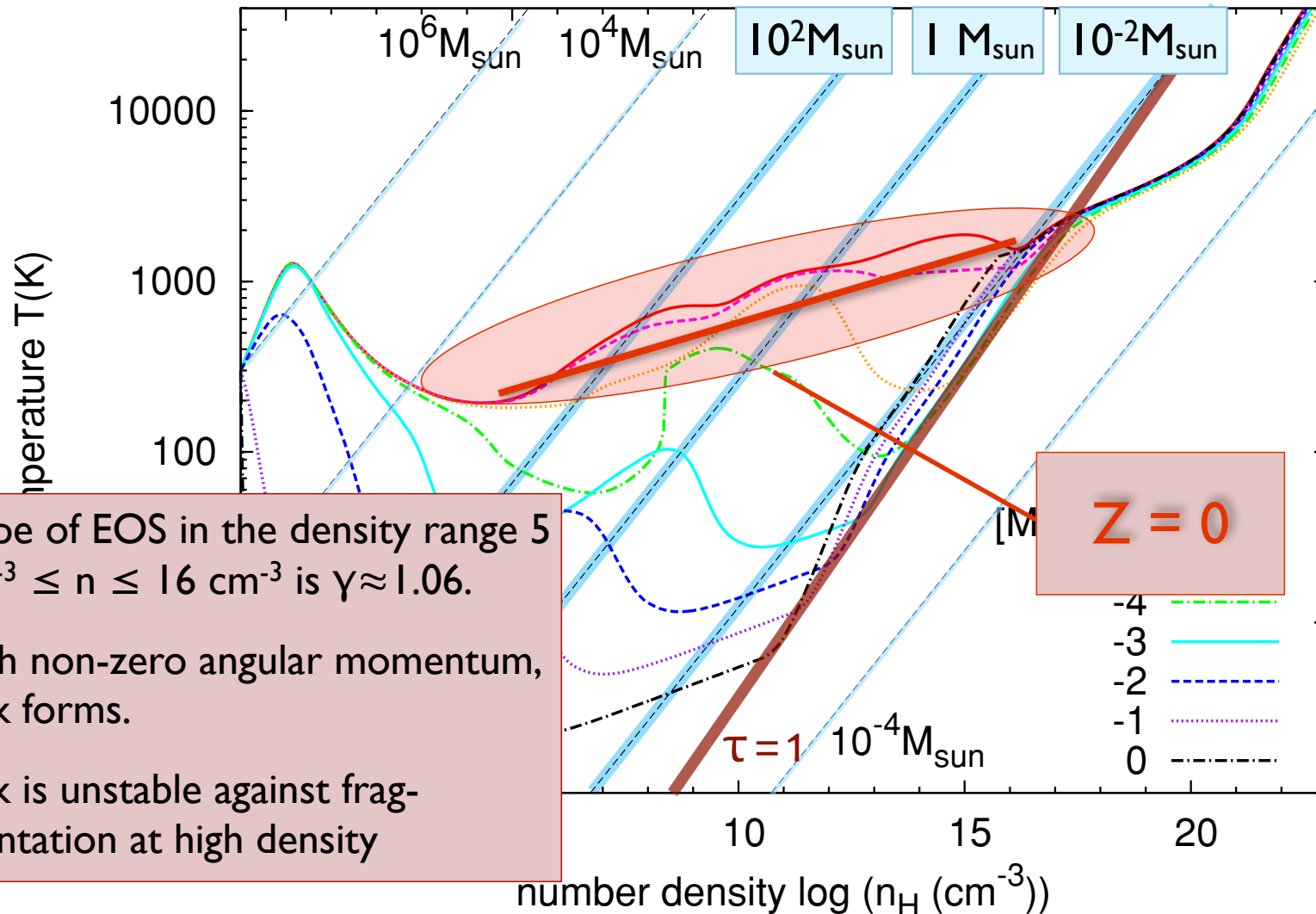


EOS as function of metallicity



(Omukai et al. 2005, 2010)

EOS as function of metallicity



(Omukai et al. 2005, 2010)

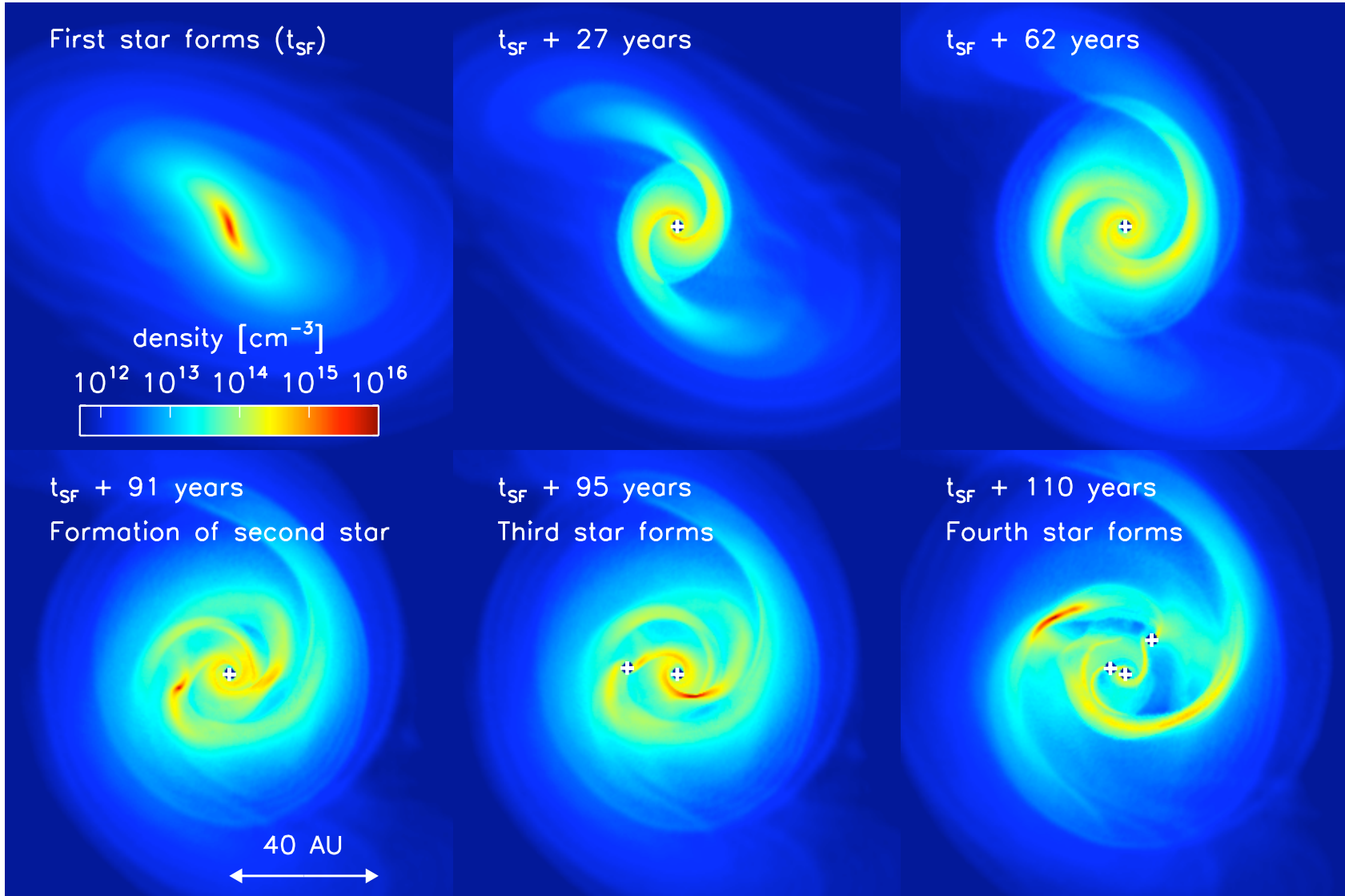
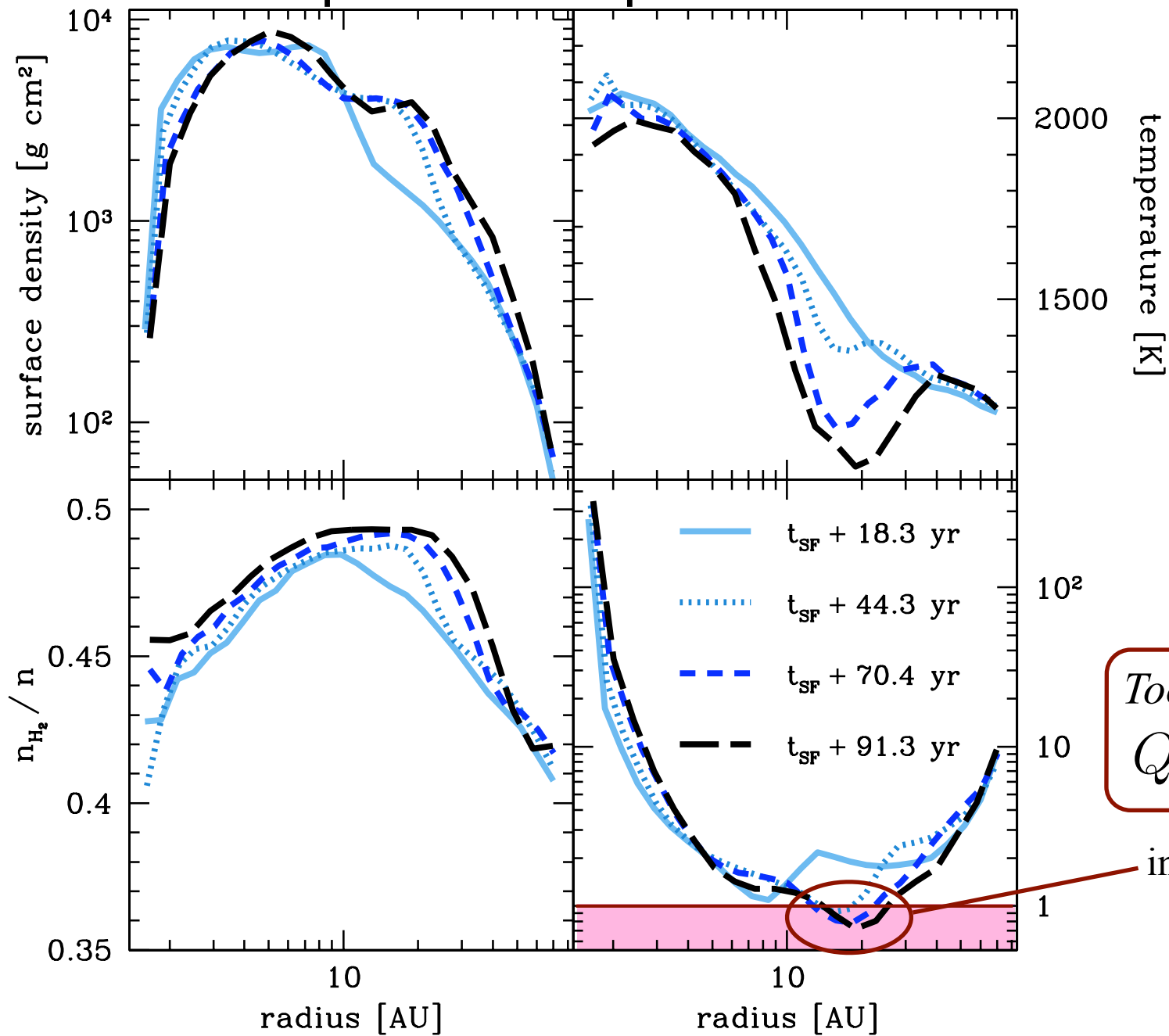


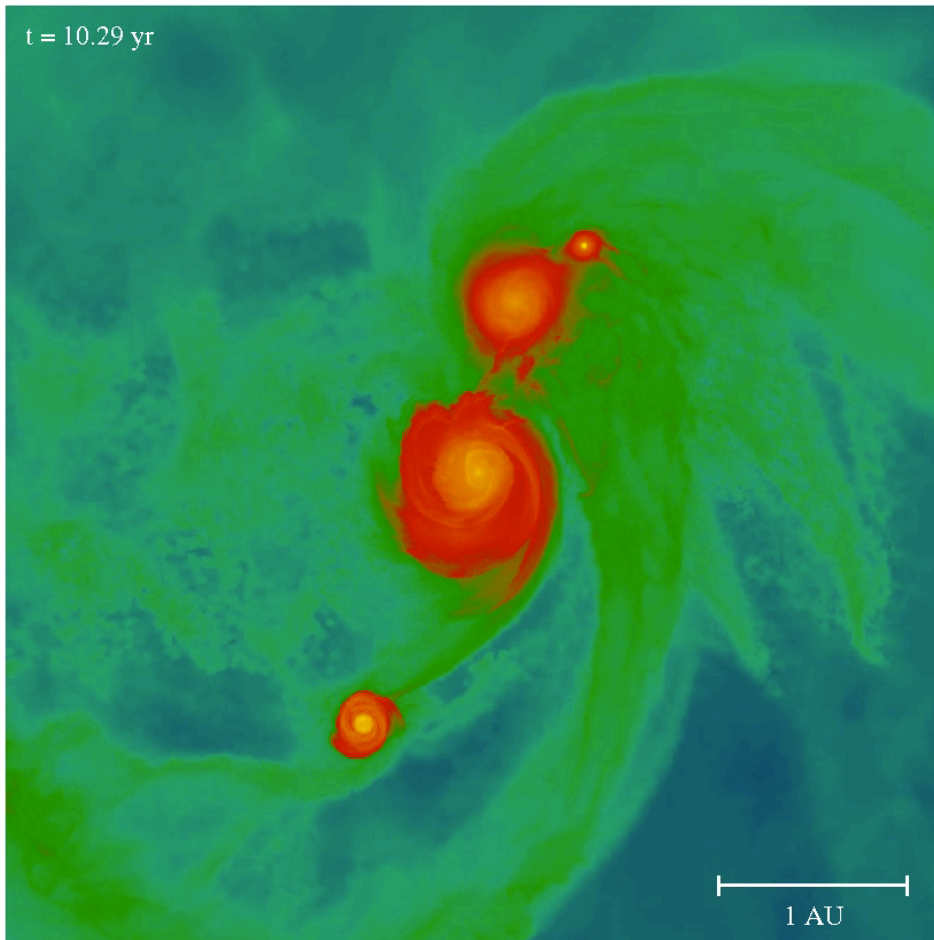
Figure 1: Density evolution in a 120 AU region around the first protostar, showing the build-up of the protostellar disk and its eventual fragmentation. We also see ‘wakes’ in the low-density regions, produced by the previous passage of the spiral arms.

important disk parameters

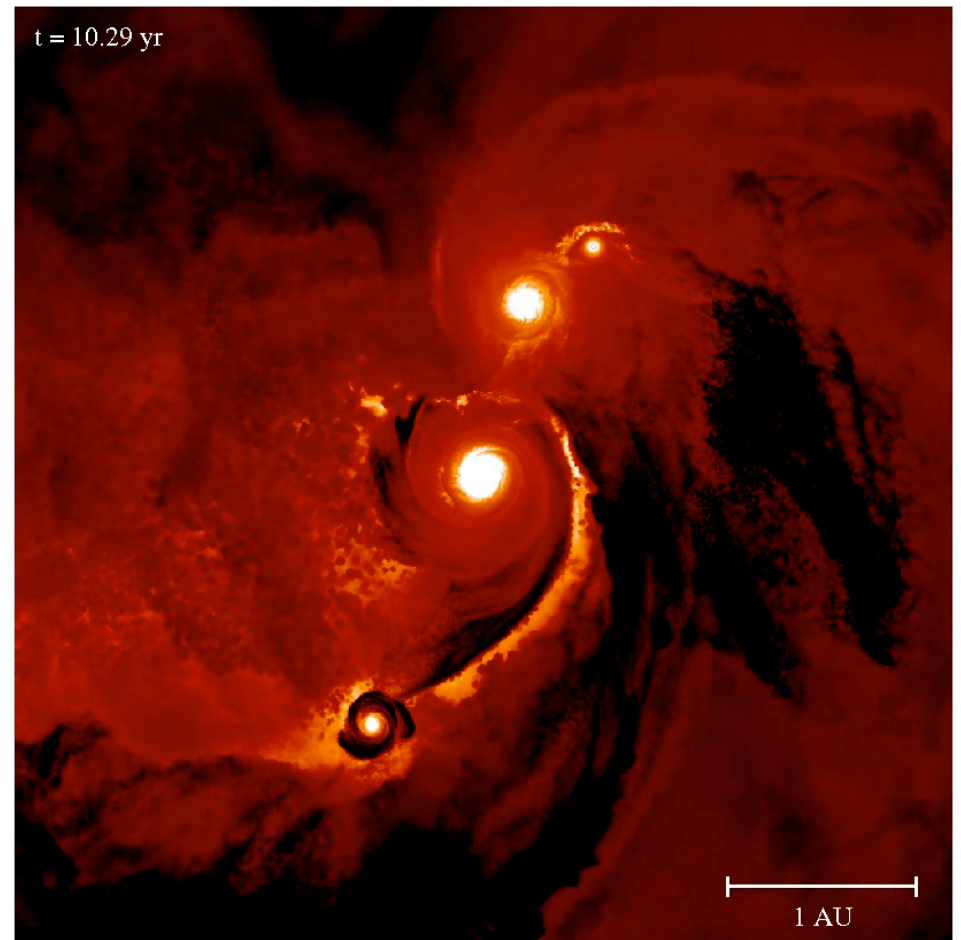


Most recent calculations:

*fully sink-less simulations, following the disk build-up over ~ 10 years
(resolving the protostars - first cores - down to 10^5 km $\sim 0.01 R_{\odot}$)*

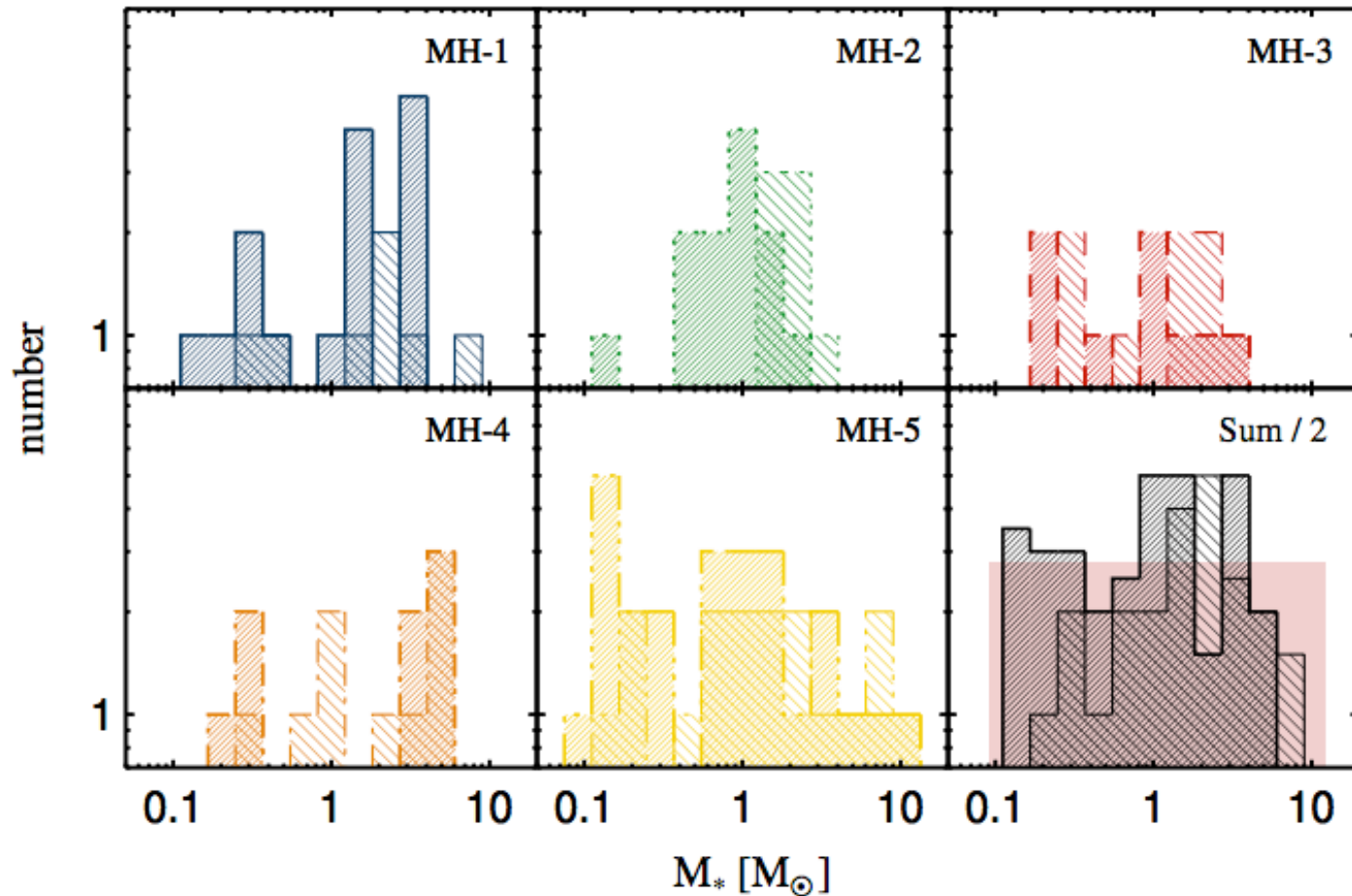


density



temperature

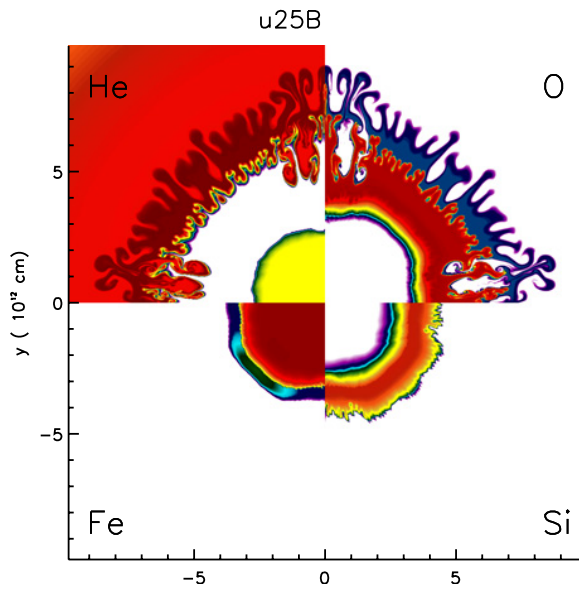
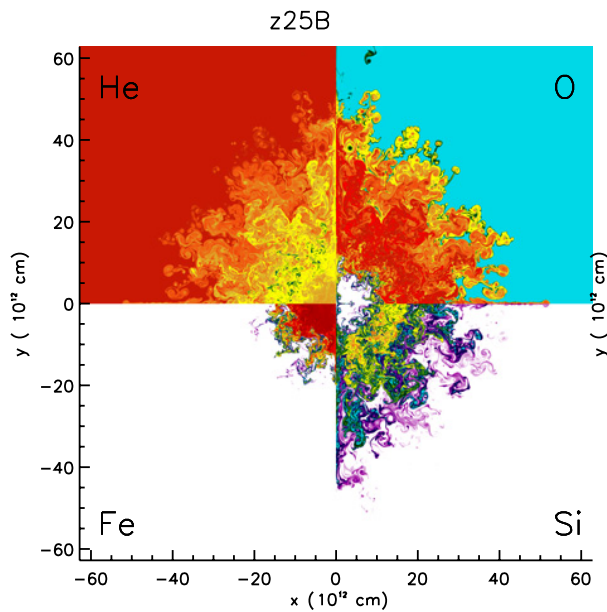
expected mass spectrum



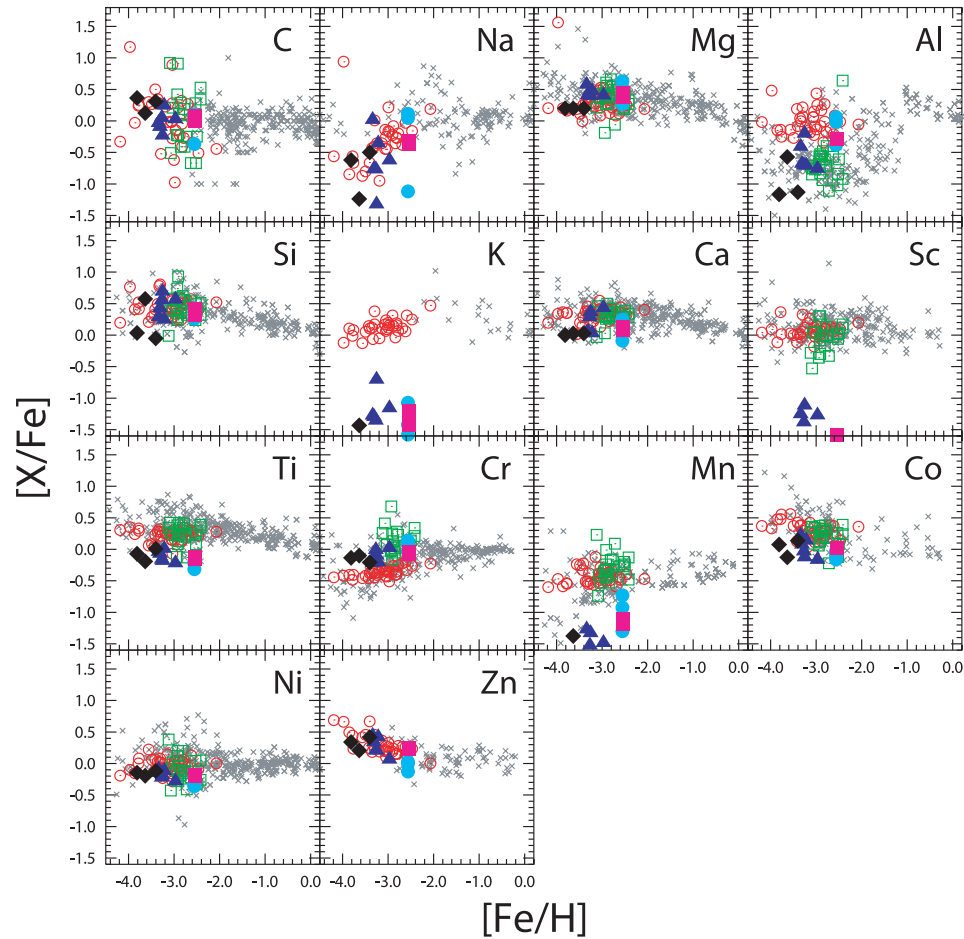
we see “flat”
mass spectrum

expected mass spectrum

- *expected IMF is flat* and covers a wide range of masses
- implications
 - because slope > -2 , most *mass is in massive objects* as predicted by most previous calculations
 - most high-mass Pop III stars should be in *binary systems* --> source of *high-redshift gamma-ray bursts*
 - because of ejection, some *low-mass objects* ($< 0.8 M_{\odot}$) might have *survived* until today and could potentially be found in the Milky Way
- consistent with abundance patterns found in second generation stars



(Joggerst et al. 2009, 2010)



(Tominaga et al. 2007)

The metallicities of extremely metal-poor stars in the halo are consistent with the yields of core-collapse supernovae, i.e. progenitor stars with 20 - 40 M_{\odot}

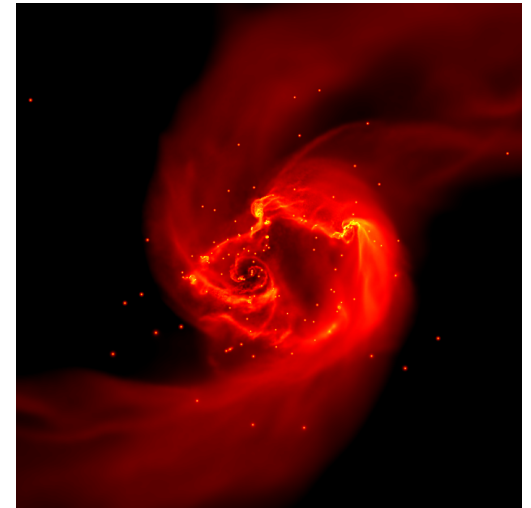
(e.g. Tominaga et al. 2007, Izutani et al. 2009, Joggerst et al. 2009, 2010)

primordial star formation

- just like in present-day SF, we expect
 - *turbulence*
 - *thermodynamics (i.e. heating vs. cooling)*
 - *feedback*
 - *magnetic fields*

to influence first star formation.

- masses of first stars still *uncertain*, but we expect a *wide mass range* with *typical masses* of several *10s* of M_{\odot}
- disks unstable: first stars in *binaries* or *part of small clusters*
- current frontier: include *feedback* and *magnetic fields* and possibly *dark matter annihilation...*



primordial star formation

- from present-day star formation theory we know, that
 - magnetic fields: Peters et al. 2011, Seifried et al. 2012, Hennebelle et al. 2011
 - accretion heating: Peters et al. 2010, Krumholz et al. 2009, Kuipers et al. 2011can influence the fragmentation behavior.
- in the context of Pop III
 - radiation: Hosokawa et al. 2012, Stacy et al. 2012a
 - magnetic fields: Turk et al. 2012, but see also Bovino et al. 2013
Schleicher et al. 2010, Sur et al. 2010, Federrath et al. 2011, Schober et al. 2012ab, 2013
- all these will reduce degree of fragmentation
(but not by much, see Rowan Smith et al. 2011, 2012, at least for accretion heating)
- DM annihilation might become important for disk dynamics and fragmentation (Ripamonti et al. 2011, Stacy et al. 2012b, Rowan Smith et al. 2012)

stellar archeology

- *if* genuine Pop III stars with $M < 0.8 M_{\odot}$ have been formed, they should be still be around !
- could be seen in current (and future) surveys of searching for extremely metal-poor stars
- QUESTION:
can we constrain the *low-mass end* of the *primordial IMF*?

stellar archeology

- can we constrain the *low-mass end* of the *primordial IMF*?

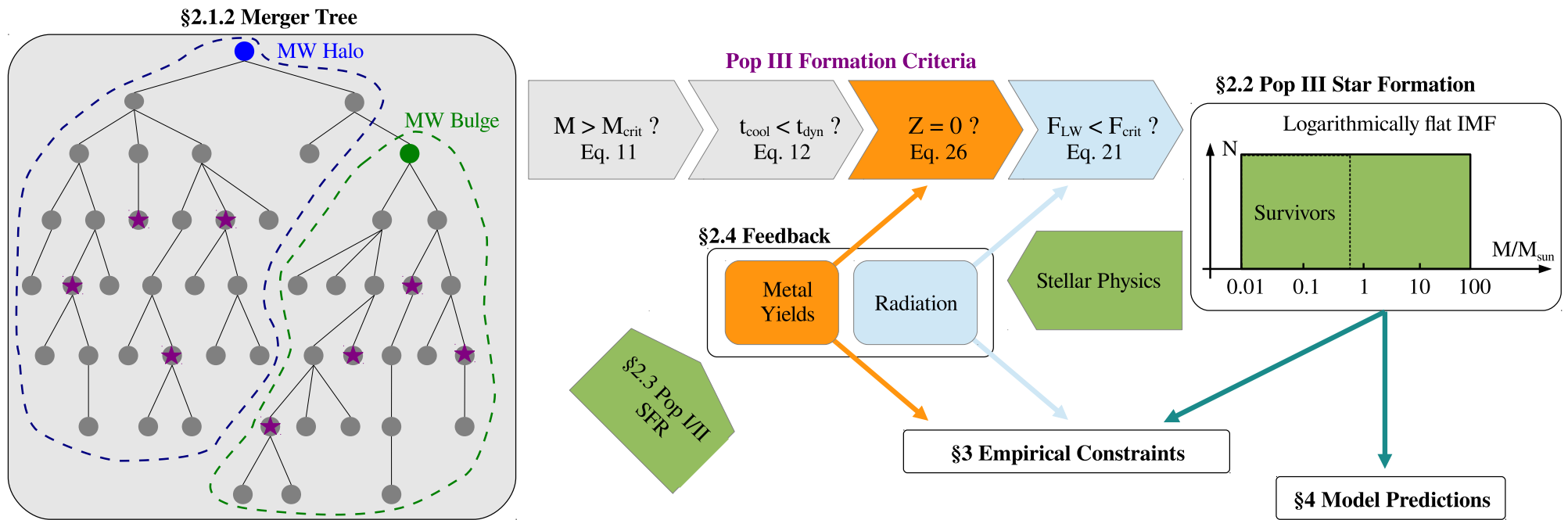
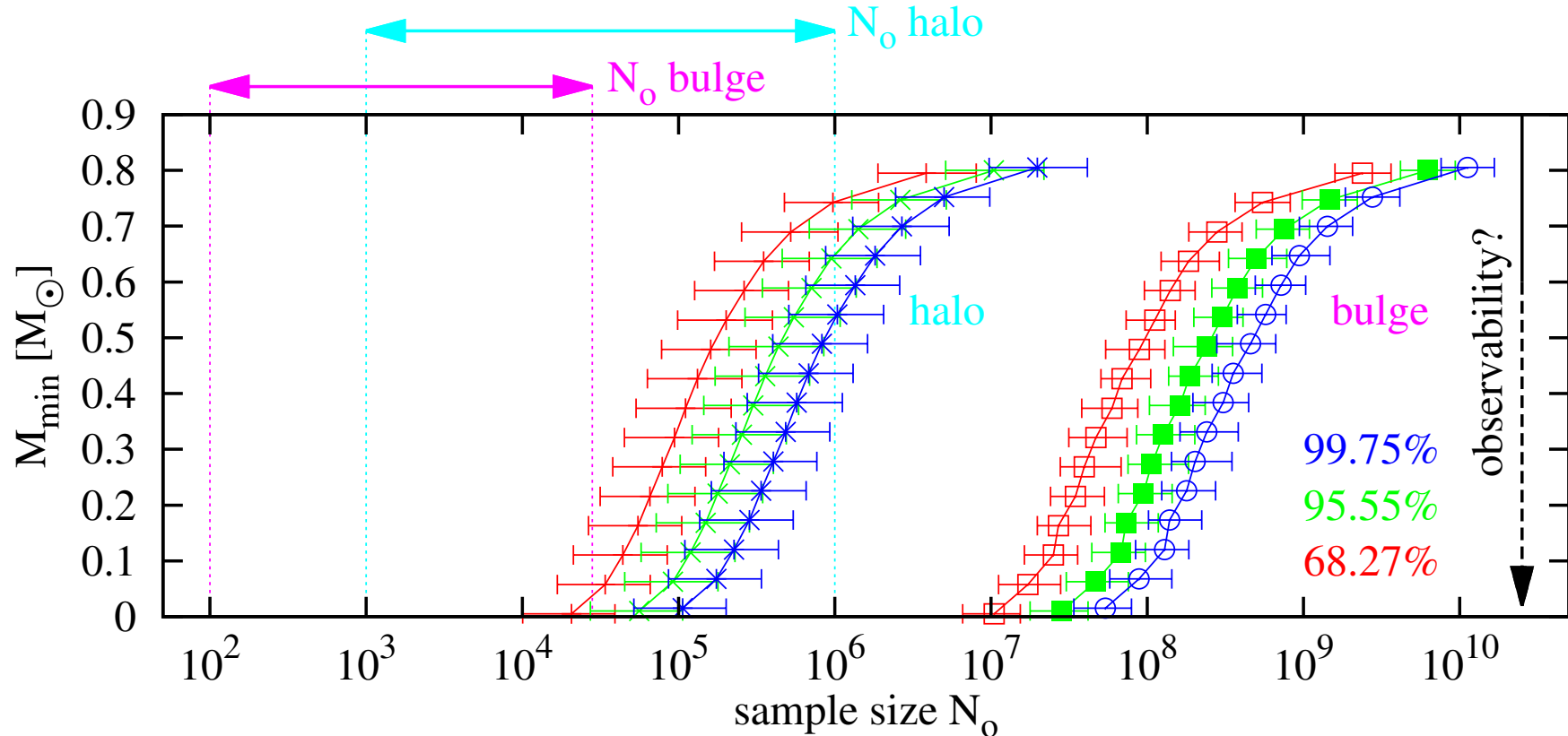


Figure 1. Roadmap, illustrating our model, with references to the relevant sections and equations. Based on the merger tree, we check which haloes are able to form Pop III stars. These checks include the critical mass, the absence of dynamical heating due to mergers, no pollution by metals and the strength of the LW background. We assign an individual number of Pop III stars to each successful halo and determine the influence on their environment. The contribution of Pop I/II star formation is modelled based on the analytical cosmic star formation history. By comparing to existing observations, we can calibrate our model parameters. Finally, we derive a prediction for the number of Pop III survivors in the Milky Way and determine constraints on the primordial IMF.

stellar archeology

- can we constrain the *low-mass end* of the *primordial IMF*?



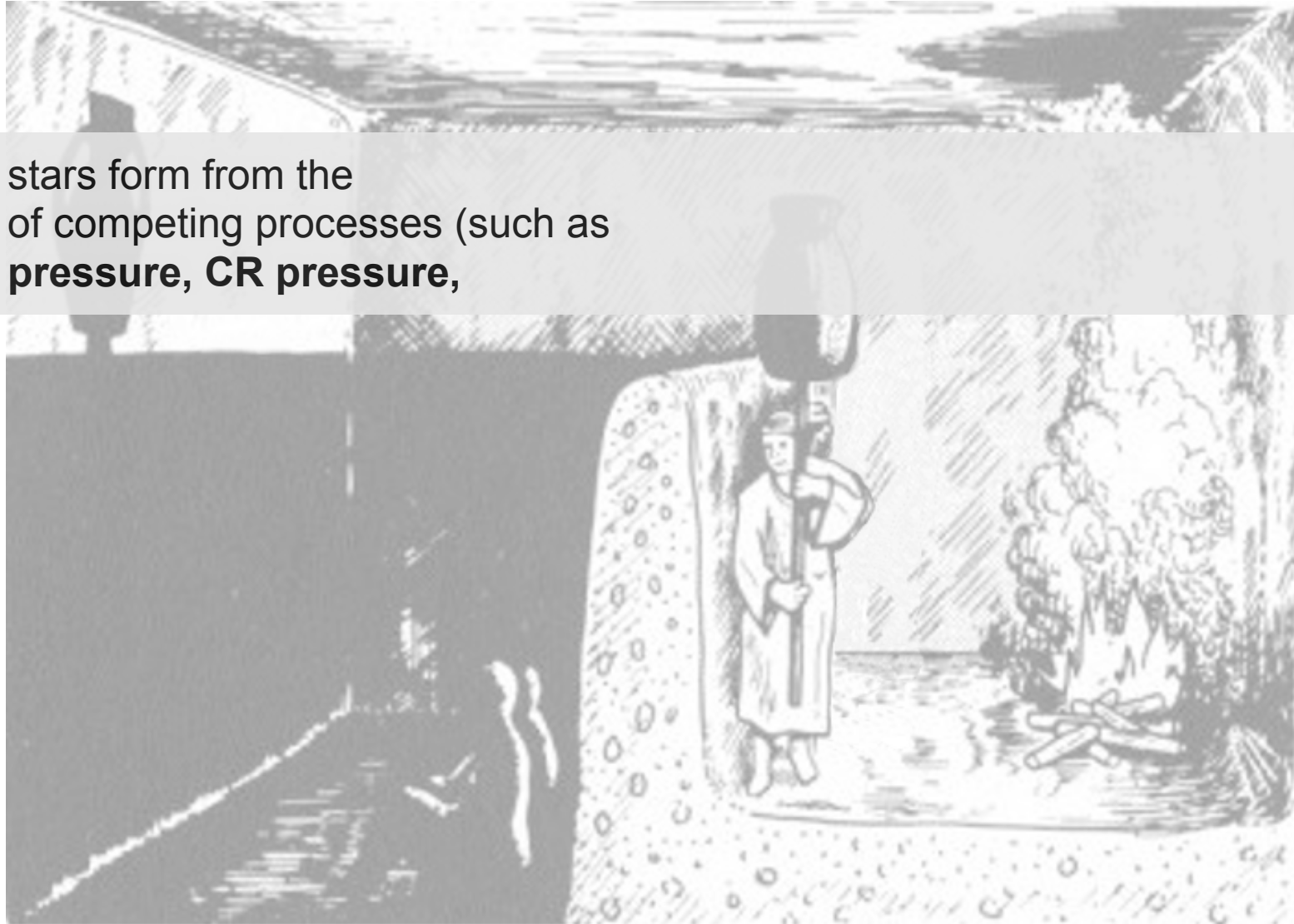
Star formation is intrinsically a multi-scale and multi-physics problem. Many different processes need to be considered simultaneously.



* The Republic
(514a-520a)

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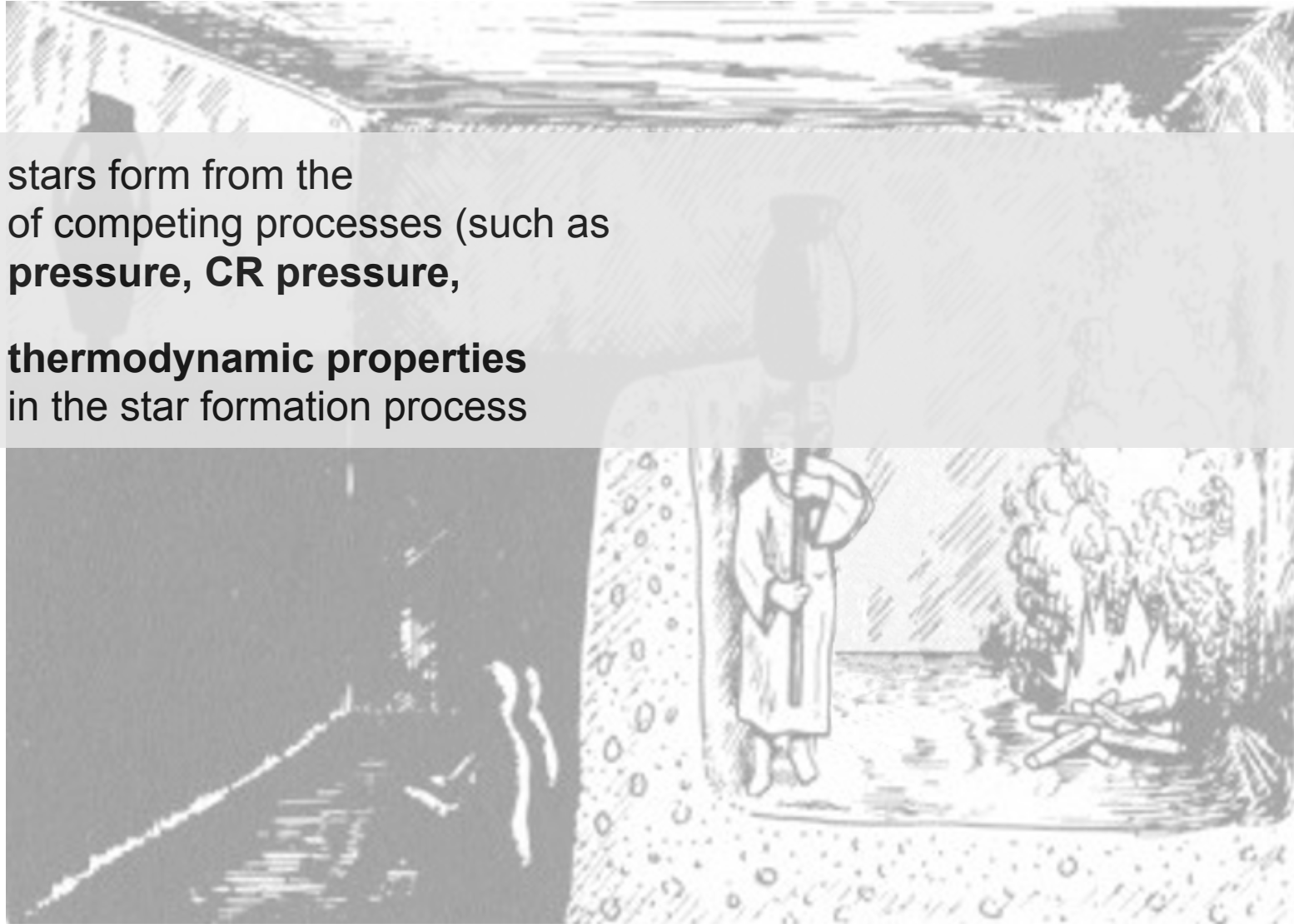
- stars form from the of competing processes (such as pressure, CR pressure,



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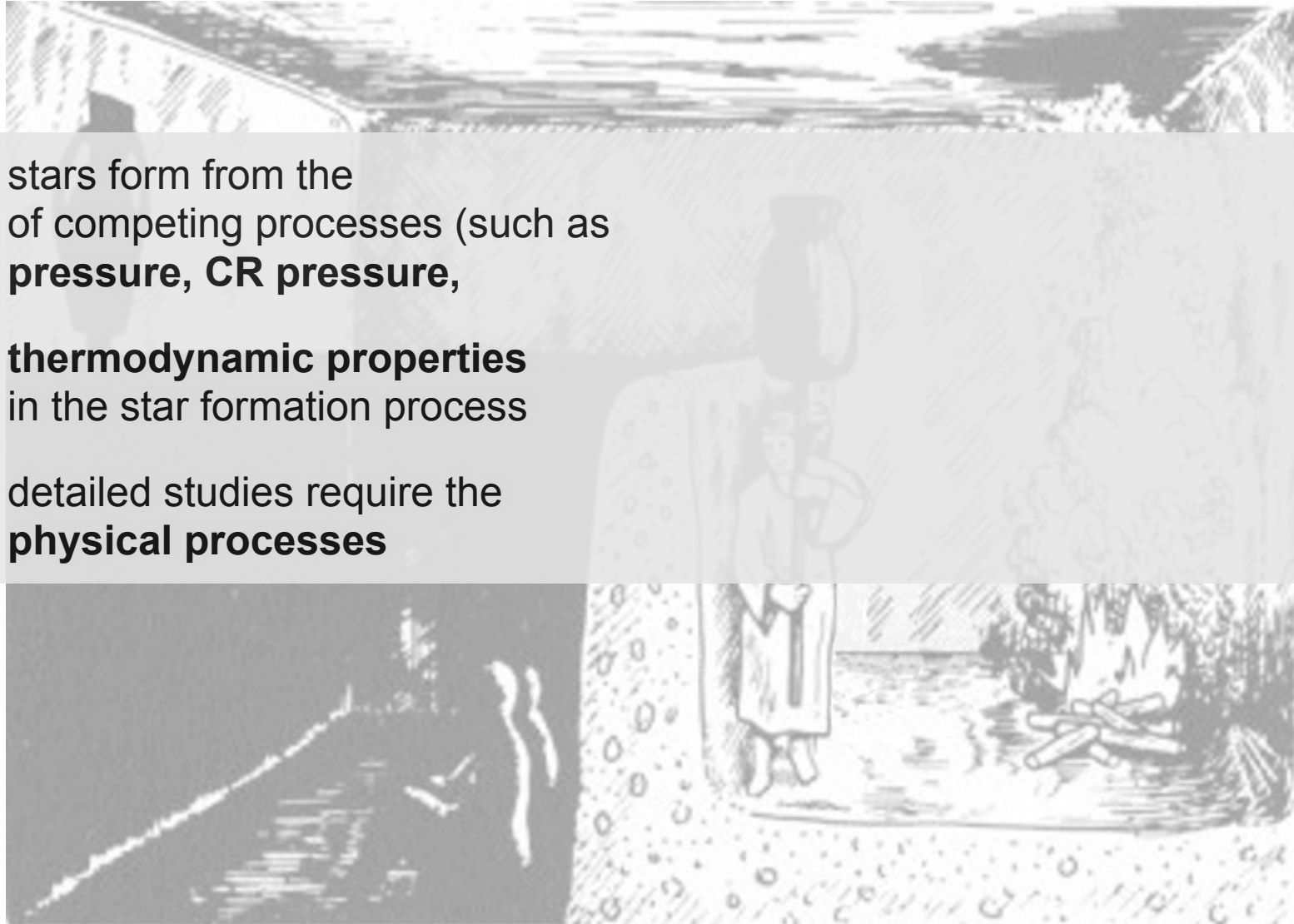
- stars form from the of competing processes (such as **pressure, CR pressure,**
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- detailed studies require the **physical processes**



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- stars form from the of competing processes (such as **pressure, CR pressure,**
- **thermodynamic properties** in the star formation process
- detailed studies require the **physical processes**
- **primordial star formation**
star formation

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