

Being discrete - Star formation on (sub-)galactic scales

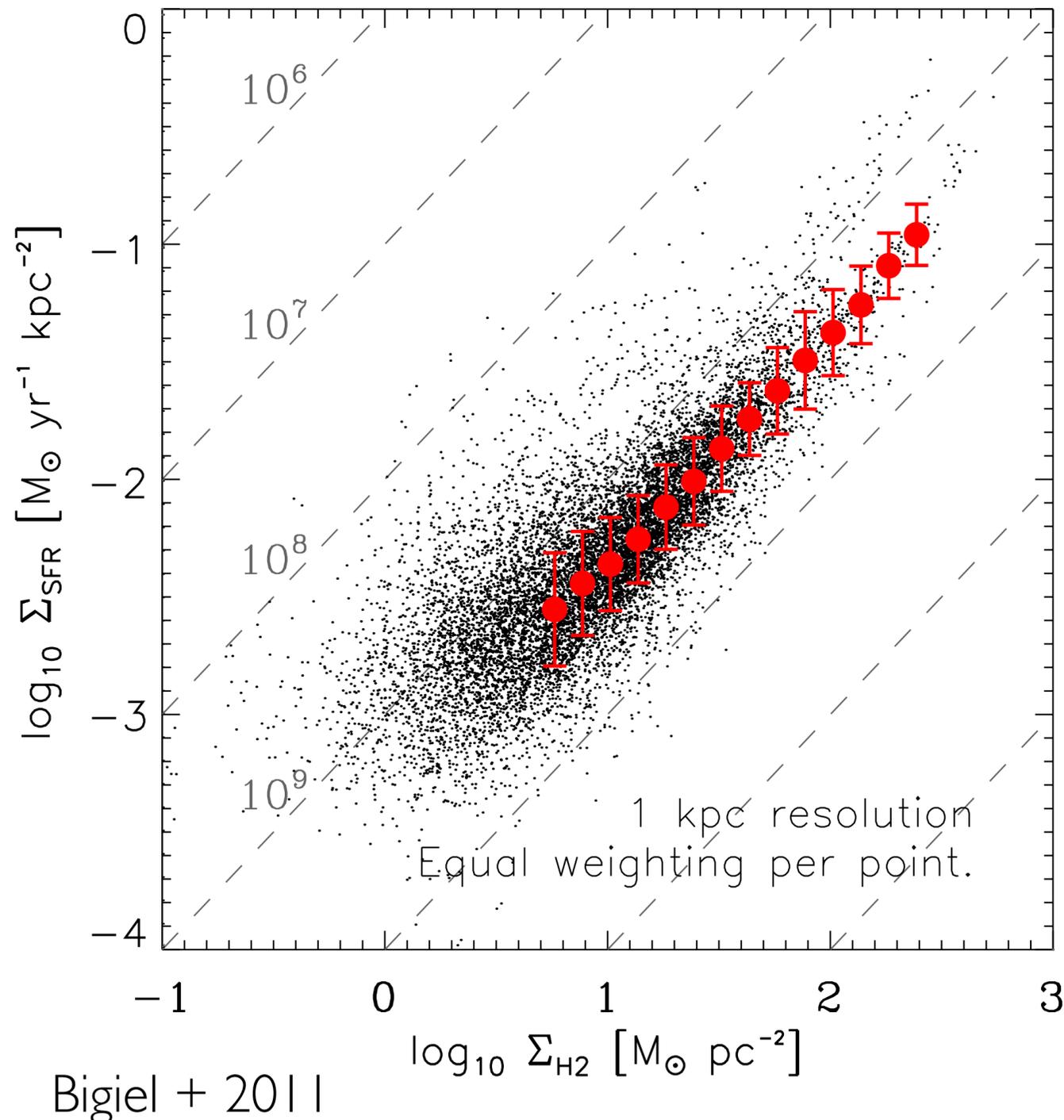
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Outline

1. Running on gas: the Σ_{H_2} - Σ_{SFR} relation
2. H_2 , CO and SF in simulations
3. The H_2 - HI transition and mid-plane pressure
4. Slope and scatter of the Σ_{H_2} - Σ_{SFR} relation



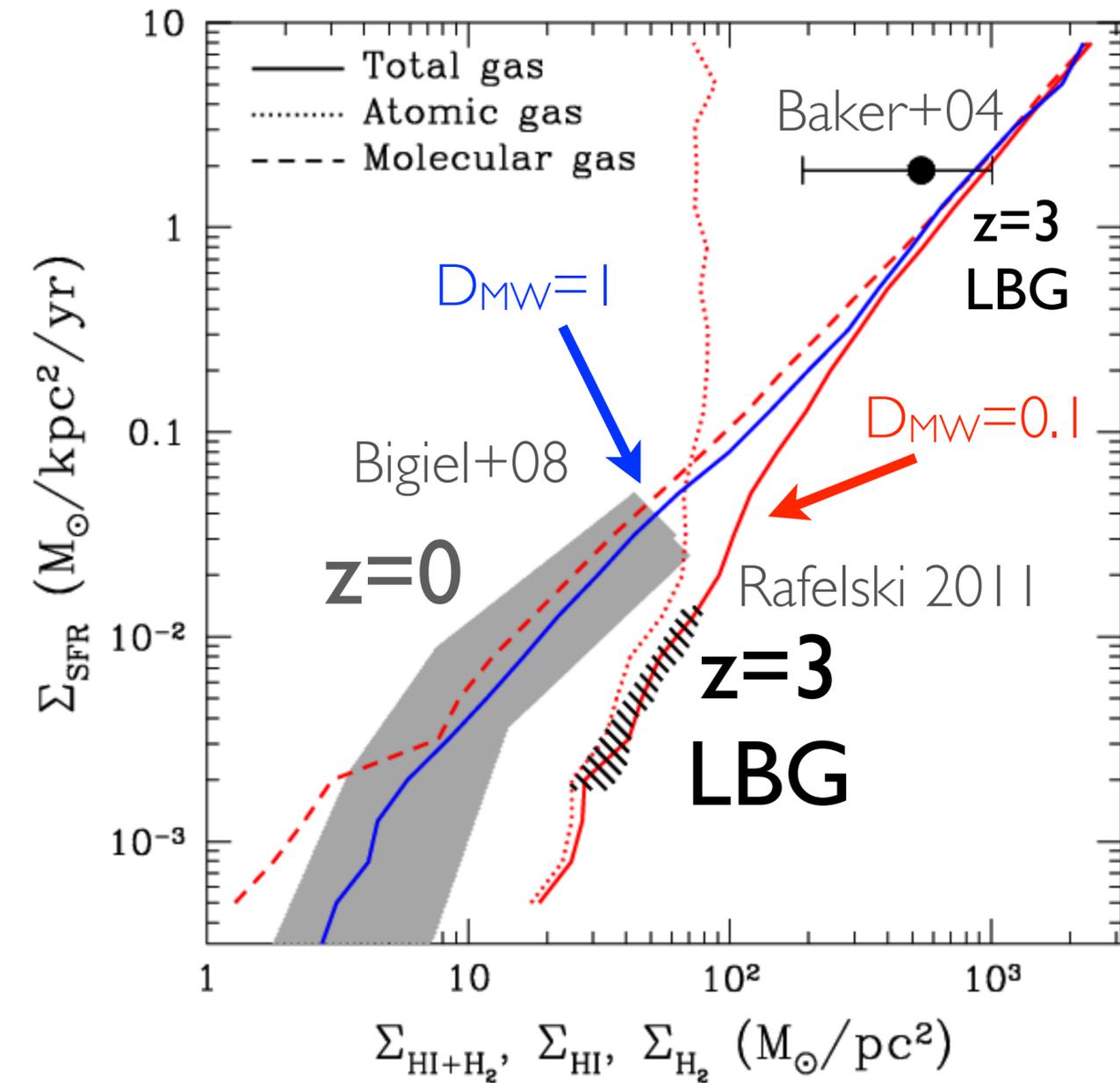
The Σ_{H_2} - Σ_{SFR} relation



- observations: tight relation on kpc scales
- relation extends into HI dominated region
- relation \sim linear \rightarrow constant gas depletion time ($\tau_{\text{dep}} \sim 1\text{-}2$ Gyr)
- slightly super-linear (1.1-1.2) when measured over 4 orders of magnitude (e.g. Genzel+10)

In this talk: I will assume $n=1$ throughout

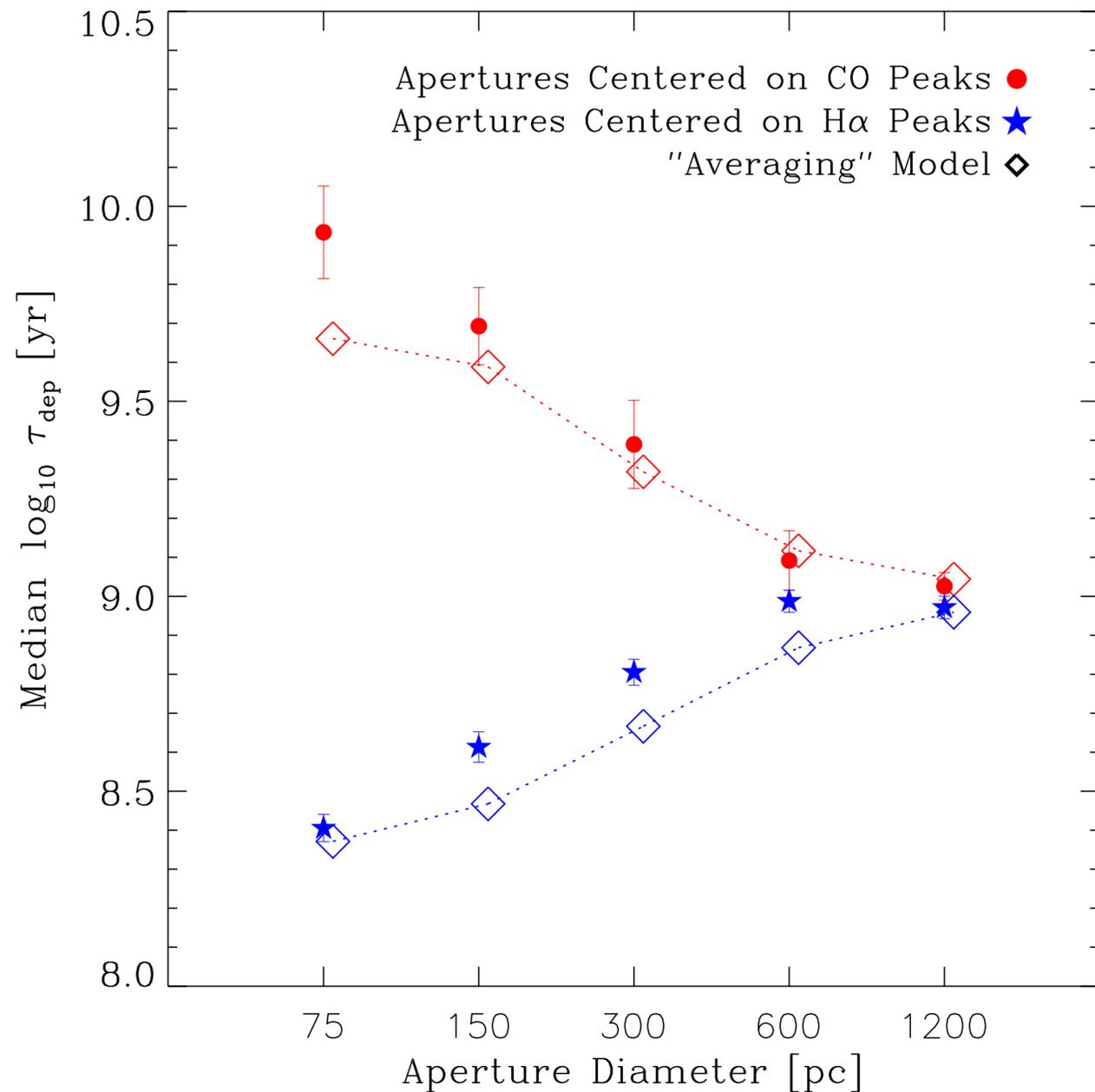
The Kennicutt-Schmidt ($\Sigma_{\text{tot}} - \Sigma_{\text{SFR}}$) relation



Gnedin & Kravtsov 2010

- Traditional K-S relation arises naturally from the $\text{HI} \leftrightarrow \text{H}_2$ transition (around Σ_c) for given Σ_{tot}
 - $\Sigma_{\text{tot}} \ll \Sigma_c \rightarrow$ gas predominantly $\text{HI} \rightarrow$ little SF indep. of Σ_{tot}
 - $\Sigma_{\text{tot}} \sim \Sigma_c \rightarrow$ H_2 fraction increases \rightarrow steep increase of SF with Σ_{tot}
 - $\Sigma_{\text{tot}} \gg \Sigma_c \rightarrow$ gas predominantly $\text{H}_2 \rightarrow \text{SF} \sim \Sigma_{\text{tot}}$
- $\text{HI} \leftrightarrow \text{H}_2$ transition
 - larger $\Sigma_{\text{tot}} \rightarrow$ increased shielding
 - larger $\Sigma_{\text{tot}} \rightarrow$ higher formation rates
 - Σ_c changes with dust-to-gas ratio

The Scatter of the Σ_{H_2} - Σ_{SFR} relation



Schruba + 2010

- scatter relatively small (~ 0.1 - 0.2 dex) on kpc scales
- but increases rapidly with resolution \rightarrow "break down of SF laws"
- On $<$ few 100 pc scales: clear separation between CO and H α emission
- Different interpretations:
 - Evolution model: CO regions (high τ_{dep}) evolve into H α regions (low τ_{dep})
 - regional variations of τ_{dep}
 - Discreteness of SF (see later)

Questions

Why does the $\Sigma_{\text{H}_2} - \Sigma_{\text{SFR}}$ relation have a slope $\sim 1.1-1.2$ when measured over 4 orders of magnitude ?

What is the origin of the scatter of the $\Sigma_{\text{H}_2} - \Sigma_{\text{SFR}}$ relation & can we understand its change with scale?

What is driving the $\text{H}_2 - \text{HI}$ transition in galaxies?

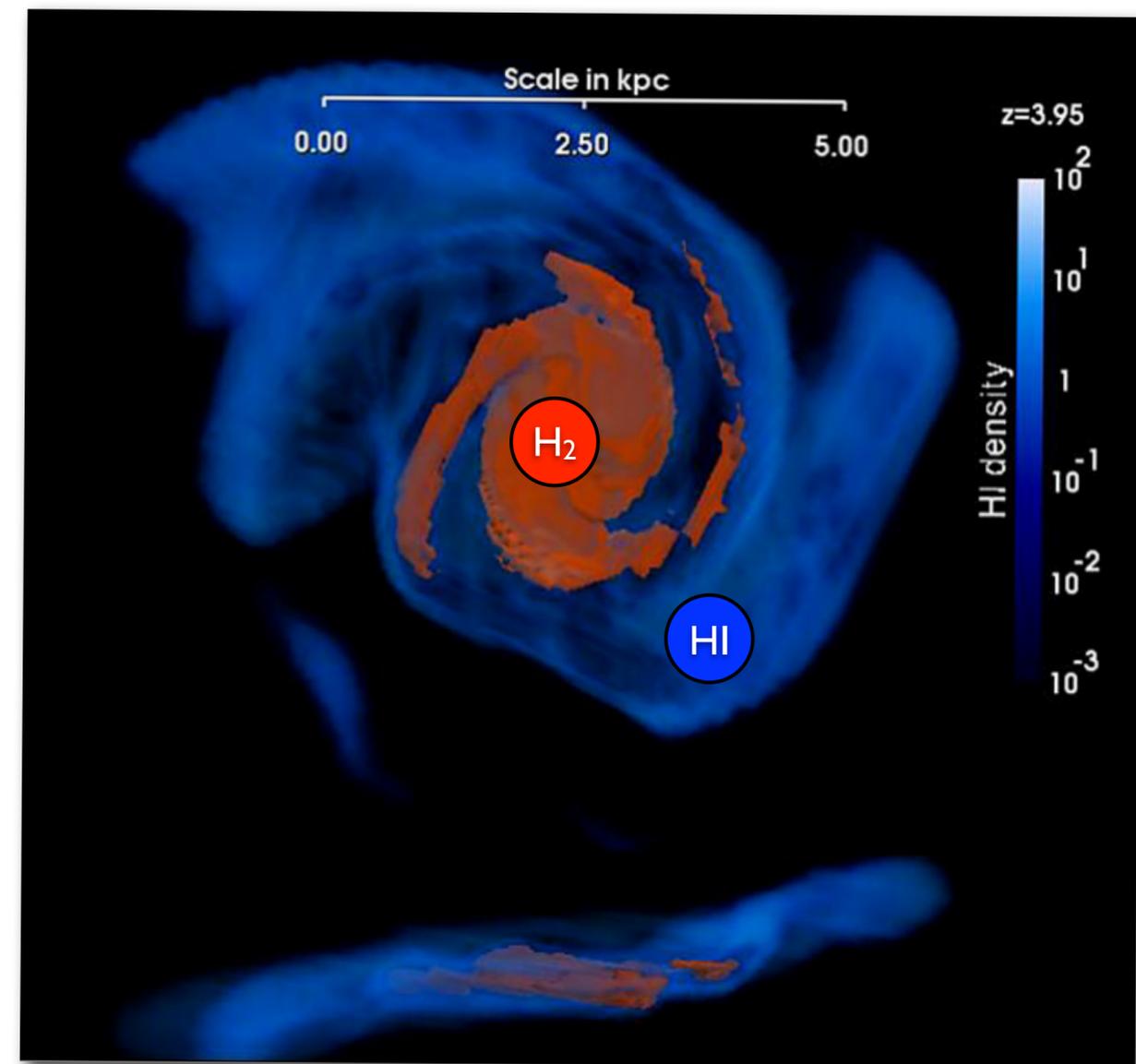
What role does the CO-H_2 conversion factor play?

N-body + AMR hydro code ART*

- non-equilibrium cooling & ionization
- non LTE chemical network
- radiative transfer in the LW bands (OTVET)
- metal enrichment, SN feedback
- H₂ formation & destruction, including dust-shielding and self-shielding
- subgrid modeling of stochastic SF based on H₂
- CO, FUV, H α emission in post-processing

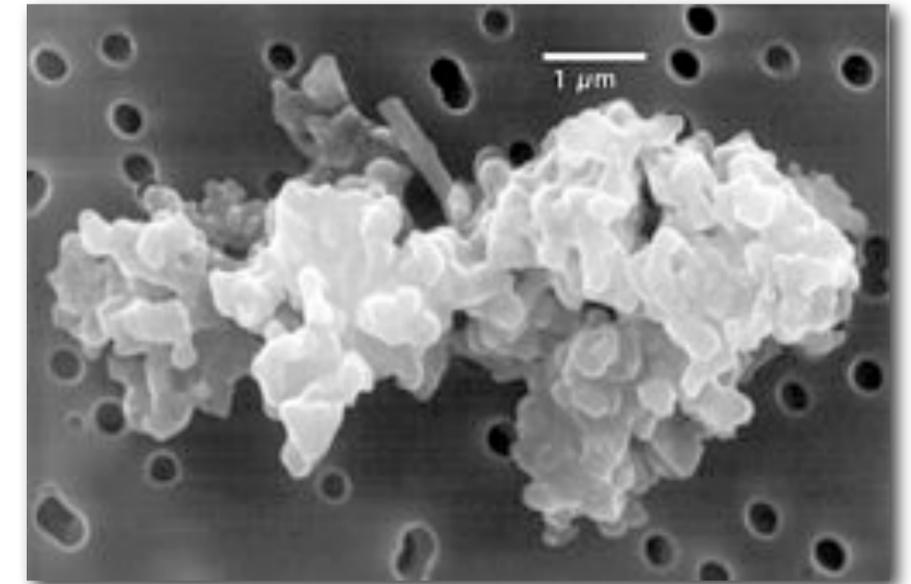
- cosmological, zoom-in simulations
- $\Delta x \sim 60$ pc, $m_{\text{DM}} \sim 1.3 \times 10^6 M_{\odot}$

*Adaptive Refinement Tree (Kravtsov+97,02)



Modeling: H₂

- H₂ density is an advected field with source terms
 - formation catalyzed on dust grains
 - dissociation by UV radiation in the LW bands
- Shielding is important: $\Gamma_{\text{LW}} \Rightarrow \Gamma_{\text{LW}} S_{\text{H}_2} S_{\text{D}}$



- Dust shielding
- H₂ self-shielding

$$S_{\text{D}} = e^{-D_{\text{MW}} \sigma_0 n_{\text{H}} L_{\text{Sob}}}$$

$$S_{\text{H}_2} = \begin{cases} 1, & \text{for } N_{\text{H}_2} < 10^{14} \text{ cm}^{-2}, \\ (N_{\text{H}_2}/10^{14} \text{ cm}^{-2})^{-3/4}, & \text{for } N_{\text{H}_2} > 10^{14} \text{ cm}^{-2}, \end{cases}$$

$$N_{\text{H}_2} \approx n_{\text{H}_2} L_c$$

CO

starting point:

- MHD, driven turbulence ISM simulations (Glover, Mac Low, 2010)
- CO abundance as function of Z & N

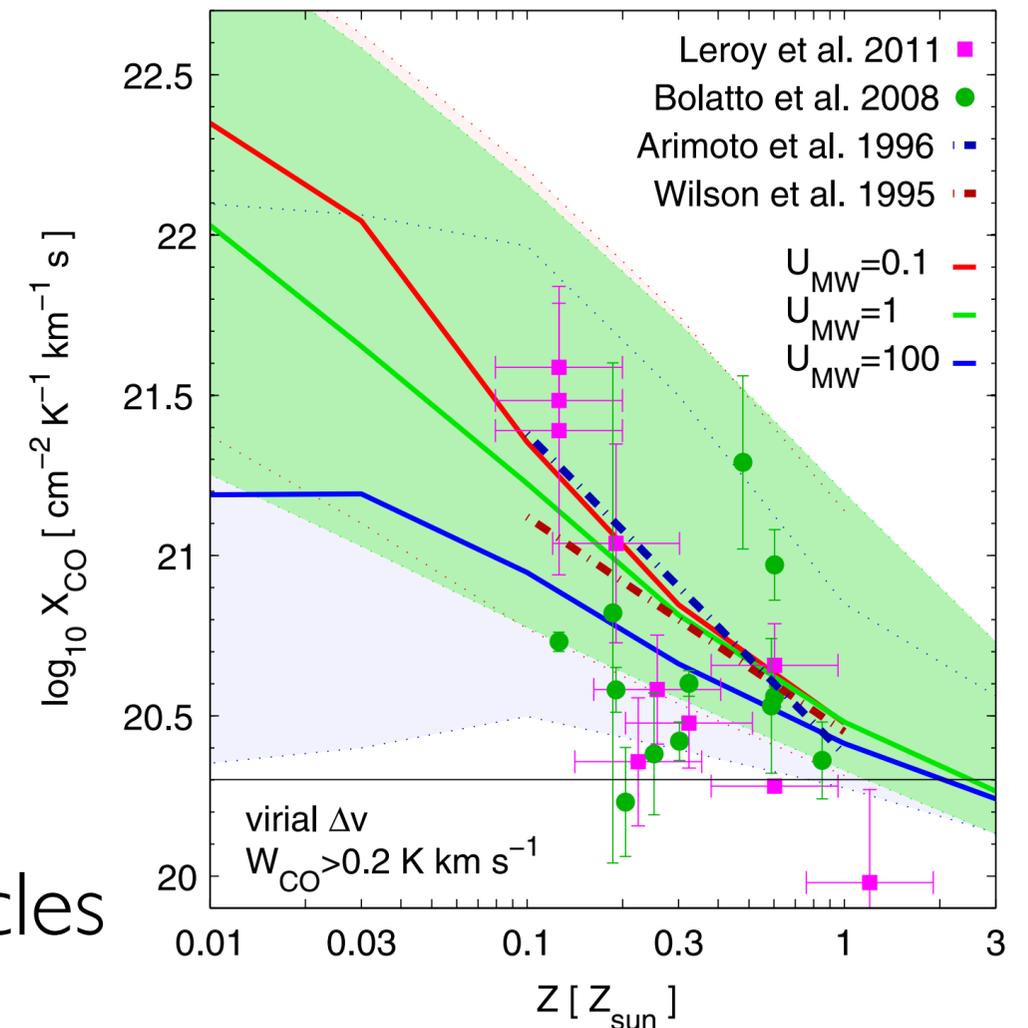
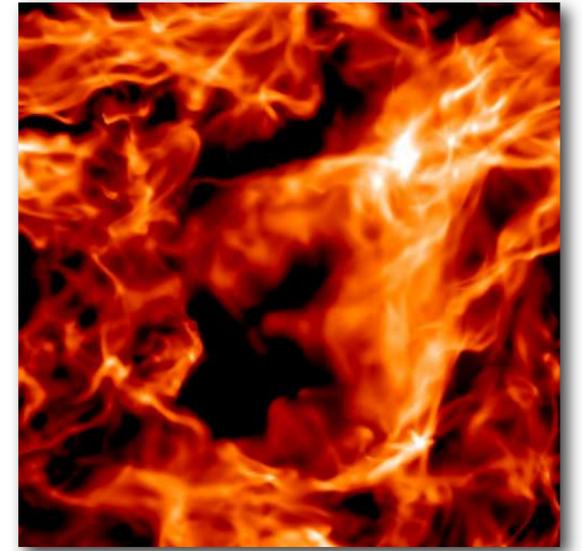
added:

- modeled of dependence on UV (photodiss. eq.)
- virial scaling of CO line width with gas surface density
- compute CO emission from each cell in the escape probability formalism

see Feldmann et al. 2012, APJ, 747, 2 *arxiv:1112.1732*

SFR tracers ($H\alpha$, FUV)

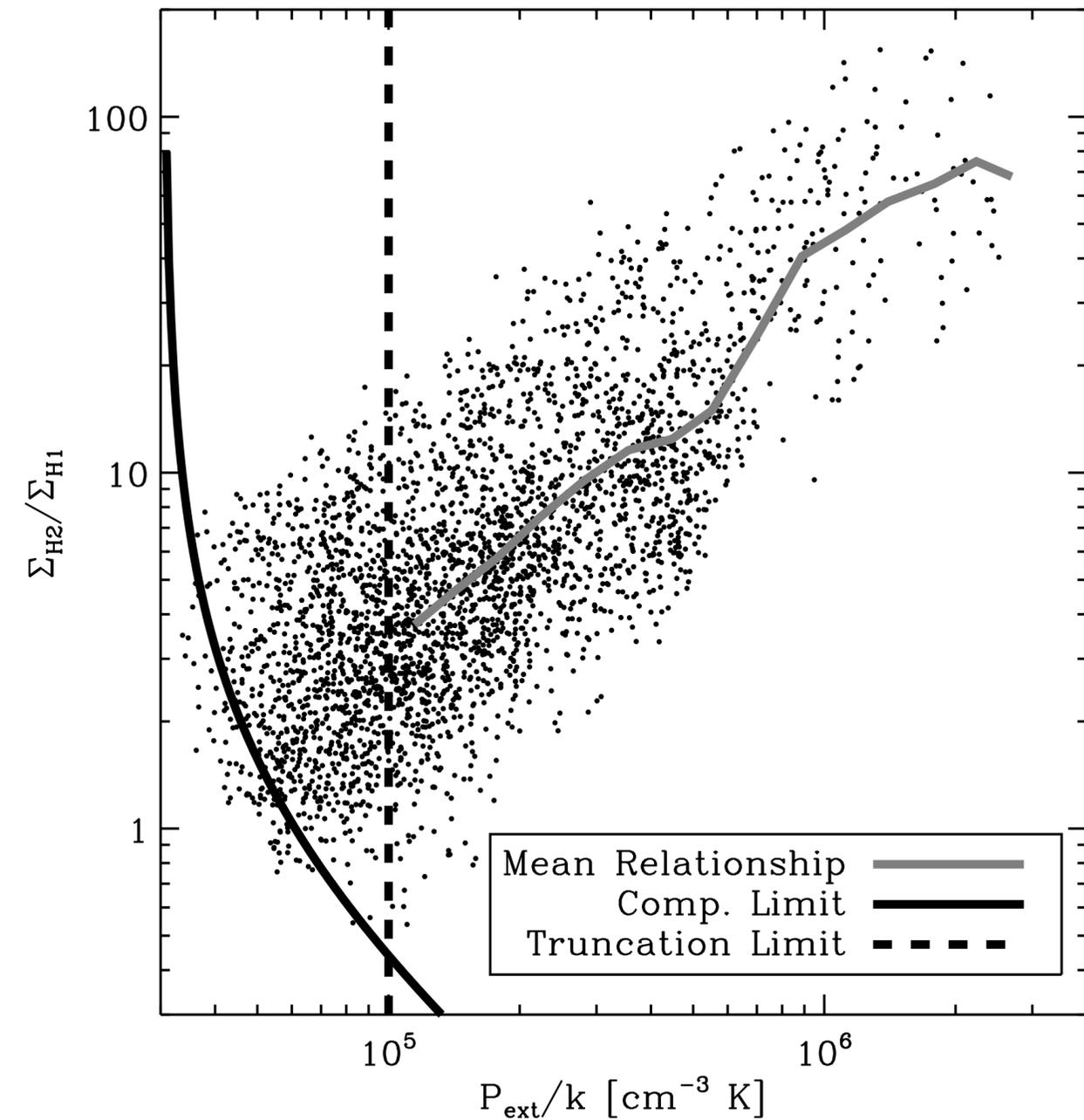
- use Starburst-99 + ages and metallicities of star particles



The H₂ to HI transition

The H₂ to HI ratio scales with midplane pressure

$$P_{\text{ext}} = (2G)^{0.5} \Sigma_g v_g \left[\left(\frac{\Sigma_*}{h_*} \right)^{0.5} + \left(\frac{\pi}{4} \frac{\Sigma_g}{h_g} \right)^{0.5} \right]$$



Blitz & Rosolowski 2006

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often $R_{\text{mol}} > 1 \Rightarrow$ H₂ dominated regime

Elmegreen '93 formula: $R_{\text{mol}} \sim P^{2.2}/j$ derived for $f_{\text{H}_2} \approx 0.1$!

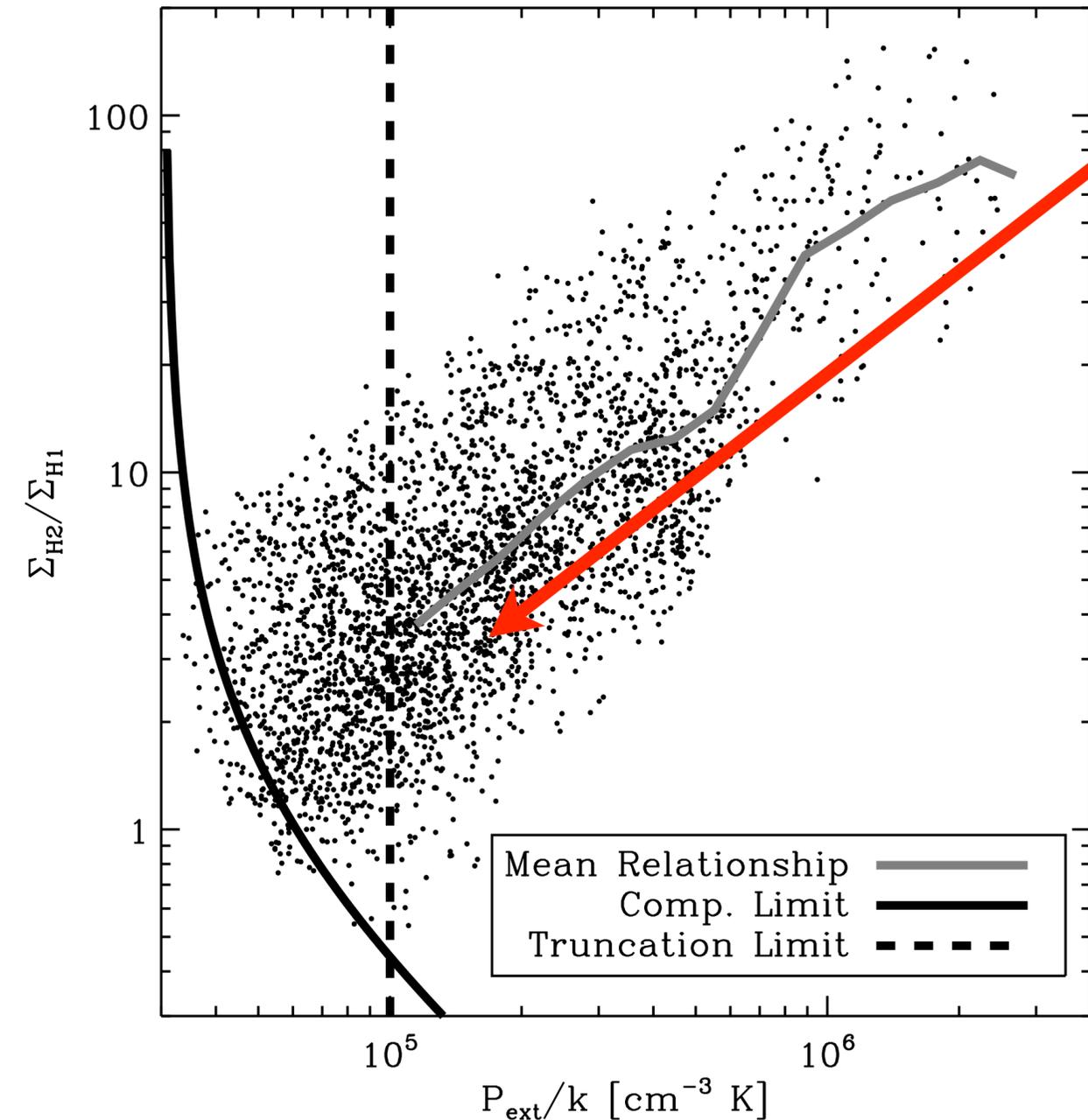
no reason to expect it to describe the asymptotic ($f_{\text{H}_2} \sim 1$) slope

simpler:
for $f_{\text{H}_2} \sim 1$

$$\Sigma_* \propto \Sigma_g^{1+\alpha} \text{ or } \Sigma_* \ll \Sigma_g$$

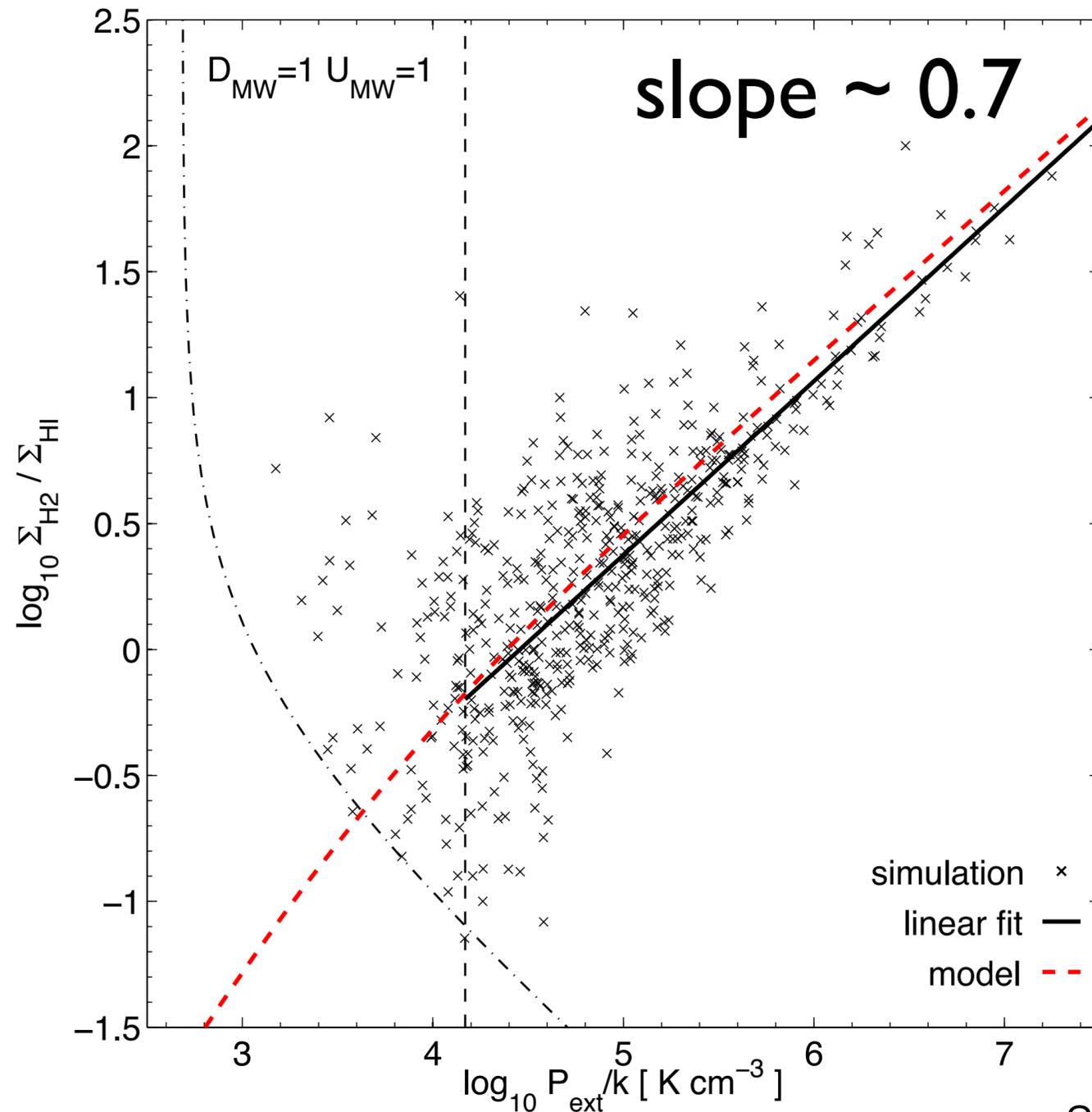
$$P_{\text{ext}} \propto \Sigma_g^{1.5+0.5\alpha}$$

$$R_{\text{mol}} \propto \Sigma_g \Rightarrow R_{\text{mol}} \propto P_{\text{ext}}^{\frac{2}{3+\alpha}}$$



Blitz & Rosolowski 2006

The H₂ to HI transition



built simple model based on fitting formula by Gnedin & Kratsov 2011

Model: $\Sigma_{\text{HI}+\text{H}_2}$ given

$$\Sigma_{\text{H}_2} = \frac{\Sigma_{\text{HI}+\text{H}_2}}{(1 + \Sigma_{\text{c}}/\Sigma_{\text{HI}+\text{H}_2})^2}$$

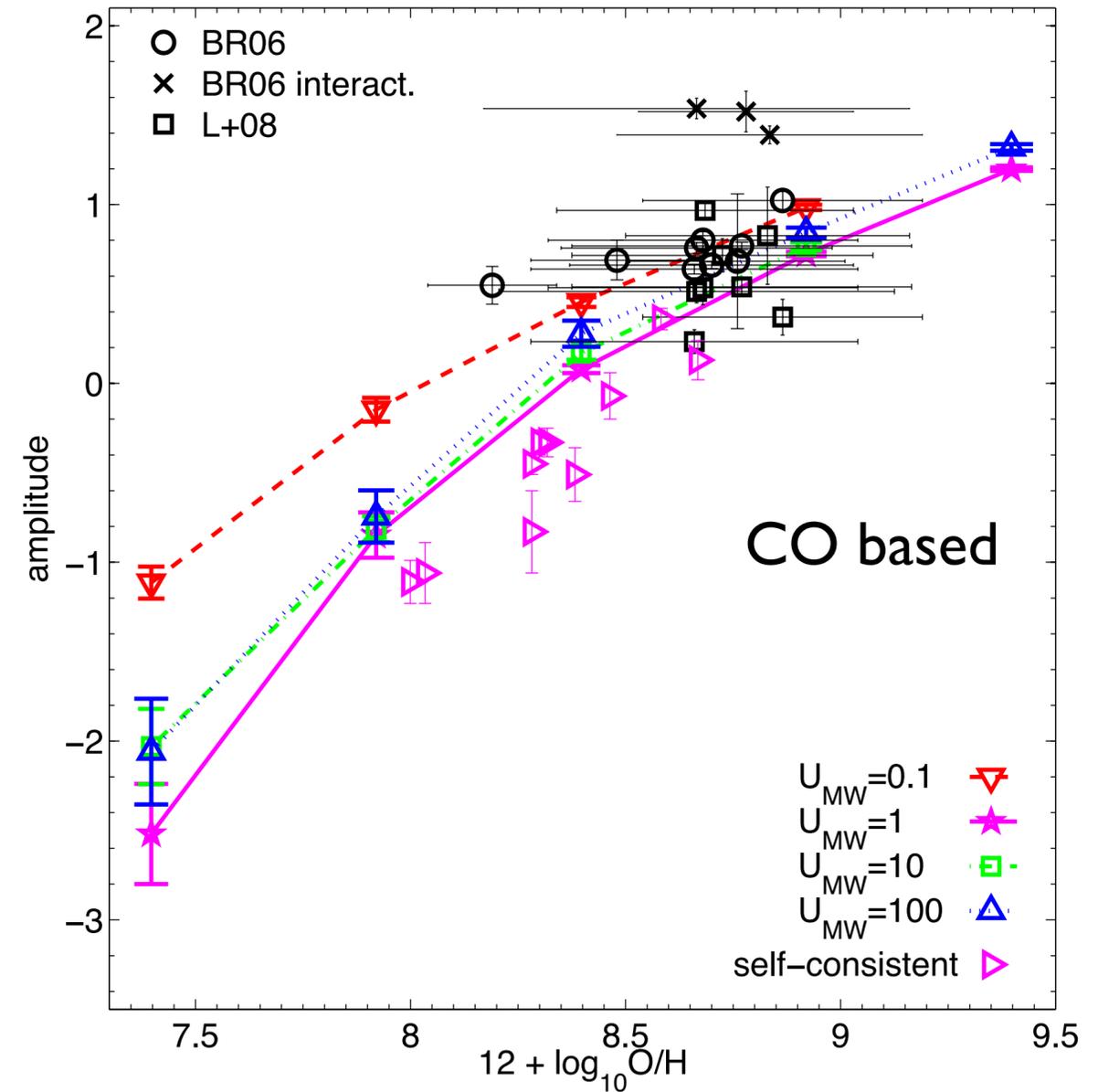
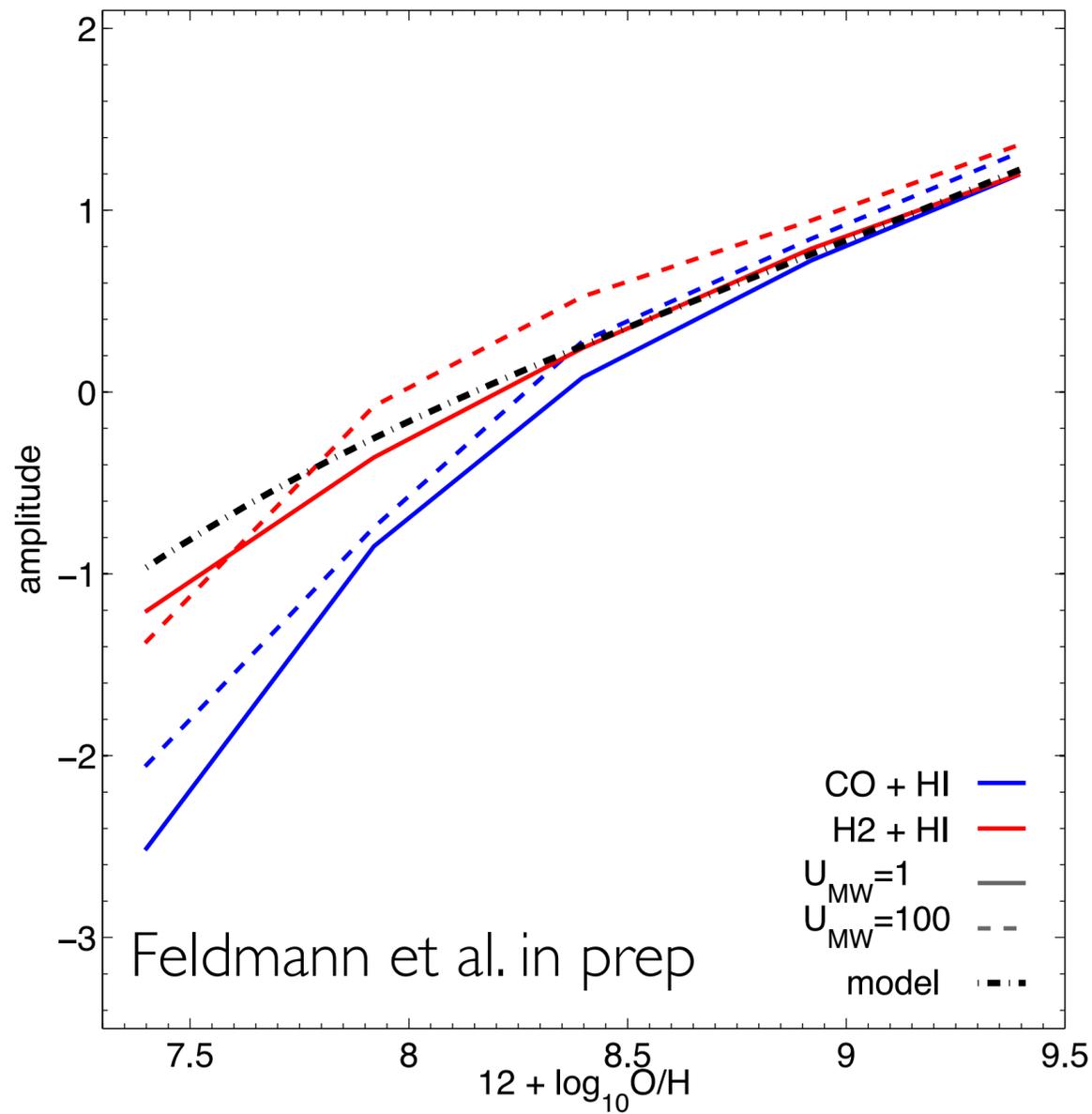
$$\Sigma_{\text{HI}} = \Sigma_{\text{HI}+\text{H}_2} - \Sigma_{\text{H}_2}$$

$$\Sigma_{\text{c}} = \Sigma_{\text{c}}(D_{\text{MW}}, U_{\text{MW}})$$

Feldmann et al. in prep

find slope ~ 0.7 (vs. B&R'06 ~ 0.9 , W&B'02 ~ 0.8),
consistent amplitude

The H₂ to HI transition



- steep decline of amplitude with Z
- well explained by model \Rightarrow i.e. by $\Sigma_c(Z)$
- more severe if H₂ based on CO (X_{CO} effect)

- no strong trend with UV field
- not in conflict with observations (but not confirmed yet either)



Modeling: Star formation

Extremely crude model!

Ansatz: SF = Poisson Process with rate $\Delta t_{\text{SF}}^{-1}$

on avg. 1 SF event per Δt_{SF}
per volume element!

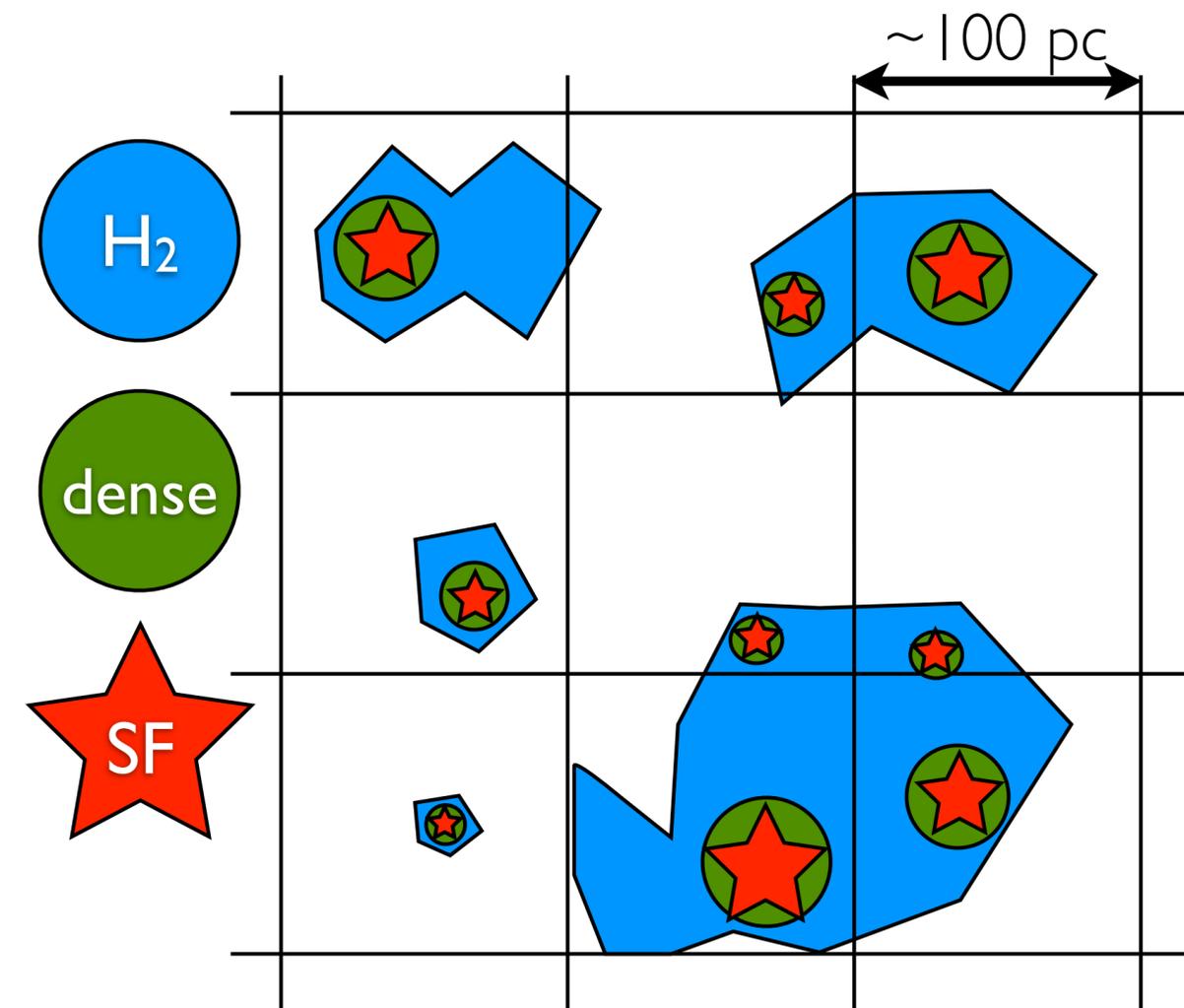
$\Delta t_{\text{SF}} \sim 10\text{-}20 \text{ Myr}$

$$\langle \text{SFR} \rangle = \frac{M_{\text{H}_2}}{\tau_{\text{dep}}} \quad (\text{ensemble avg. SFR})$$

- $N_{\Delta t}$
- draw Poisson distr. random variable with mean and variance $\Delta t / \Delta t_{\text{SF}}$
 - number of “SF events” in time Δt

$$\text{SFR}_{\Delta t} = \frac{N_{\Delta t}}{\langle N_{\Delta t} \rangle} \langle \text{SFR} \rangle \quad (\text{observed SFR})$$

Feldmann + 2012 (arXiv:1204.3910)



Consequences of the stochastic SF model

- observed depletion time given a SF tracer with lifetime Δt

$$M_{\text{H}_2}/\text{SFR}_{\Delta t} = \tau_{\text{dep}} \frac{\langle N_{\Delta t} \rangle}{N_{\Delta t}} = \tau_{\text{dep}} \frac{\Delta t / \Delta t_{\text{SF}}}{N_{\Delta t}}$$

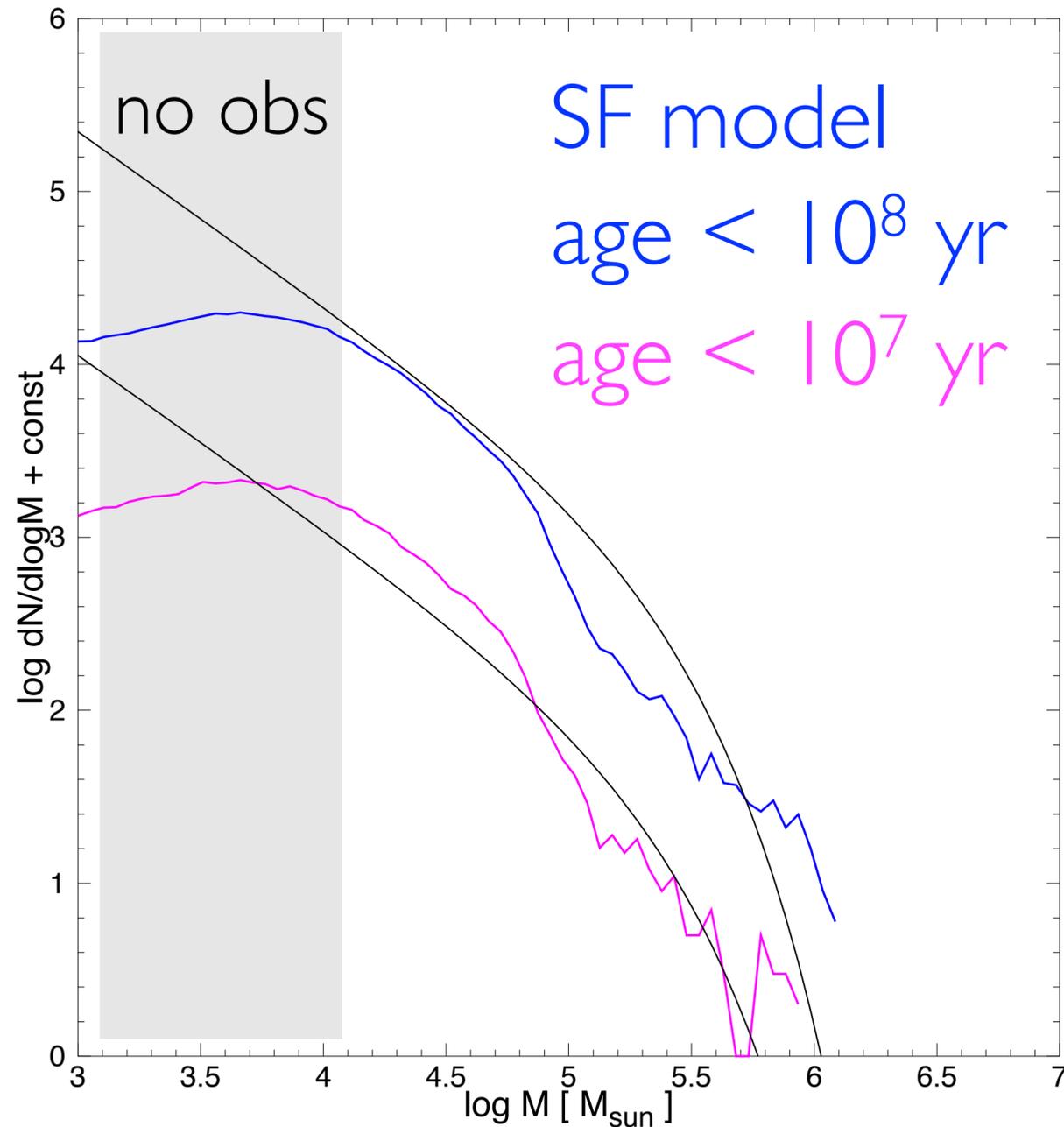
⇒ observed depletion time can be much shorter than the global average & depends on tracer

- introduces scatter in the $\Sigma_{\text{H}_2} - \Sigma_{\text{SFR}}$ relation via
 - Poisson (discreteness) noise in the number of SF events
 - fluctuations in the avg. SFR due to H_2 variations
 - their covariance

Modeling: Star formation

MF of young massive clusters

($M^* \sim 2.5 \times 10^5 M_\odot$, see Portegies Zwart et al. 2010)



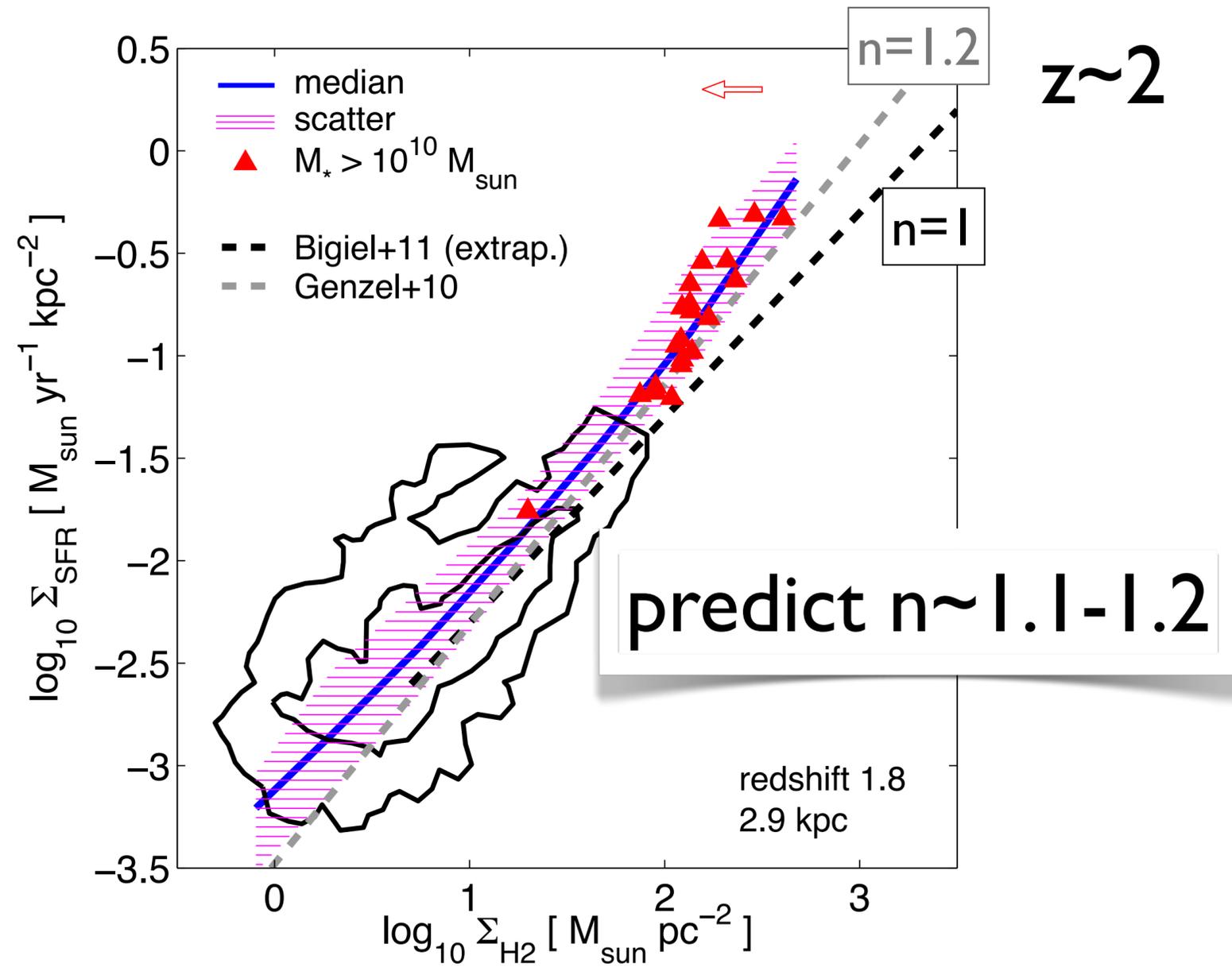
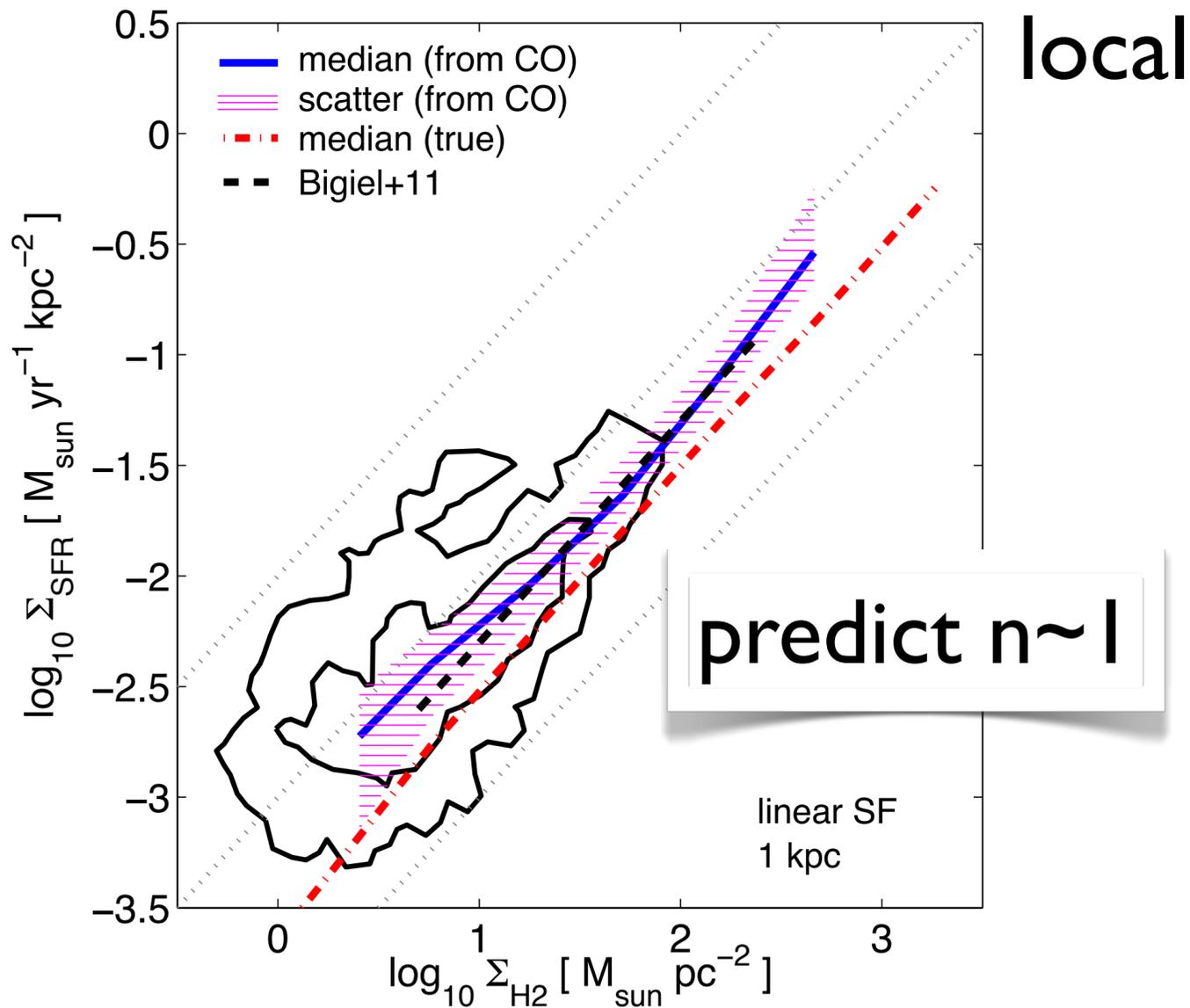
Arguments for $\Delta t_{\text{SF}} \sim 10\text{-}20 \text{ Myr}$
on $\sim 100 \text{ pc}$ scales

- lower limit: crossing time $L/\sigma \sim 10 \text{ Myr}$ (at $l=100 \text{ pc}$)
- upper limit: scatter in the $\Sigma_{\text{H}_2} - \Sigma_{\text{SFR}}$ relation
- reproduces reasonable cluster mass function
- consistent with observed relation between SFR & maximum embedded star cluster mass (Weidner+04)

Δt_{SF} fundamentally different from t_{ff}

- time between SF events
- duration of SF event
- decreases with scale
- increases with scale

The Slope of the Σ_{H_2} - Σ_{SFR} relation (reloaded)

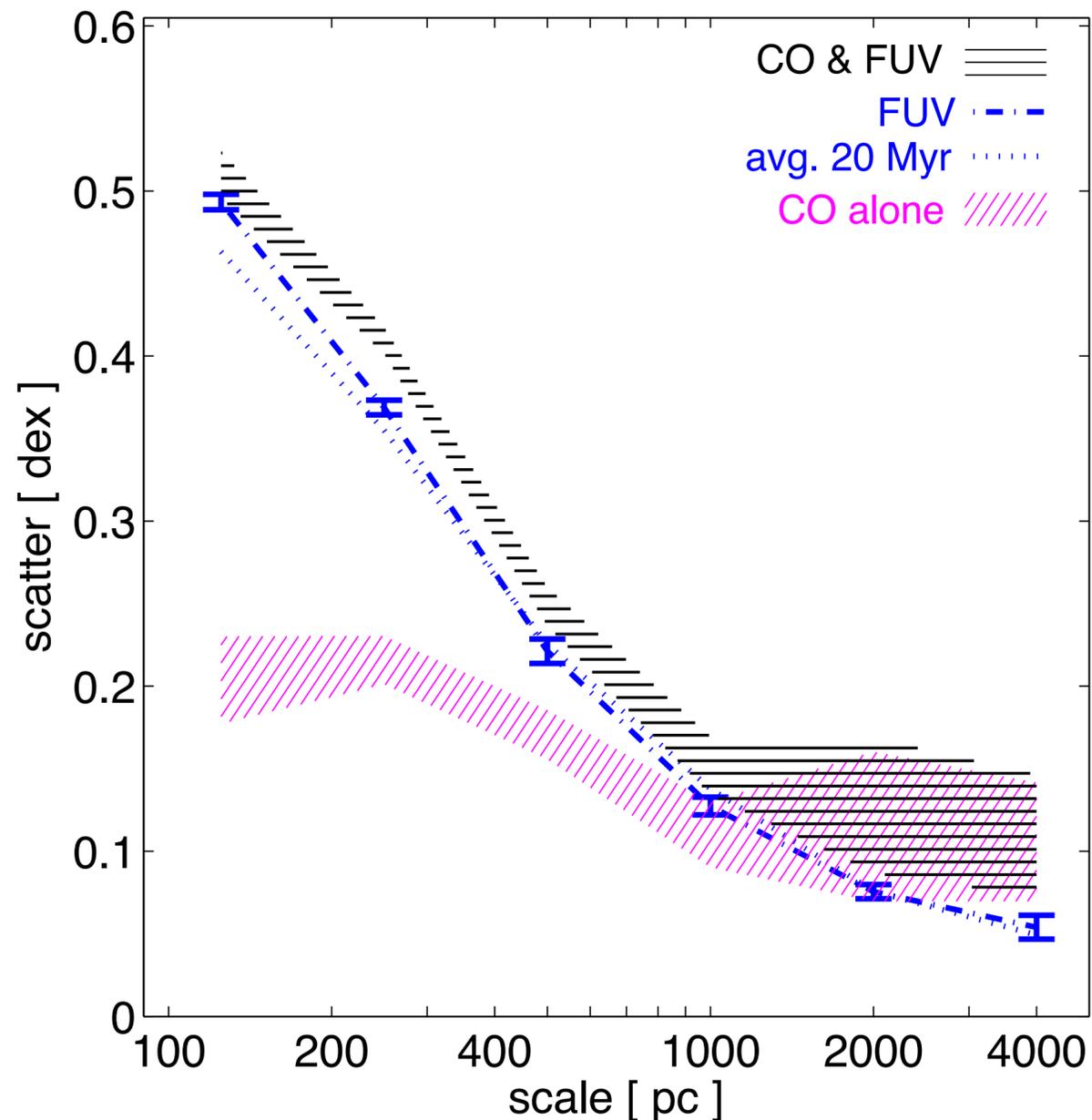


Slight super-linearity in observed molecular KS relation explained by:

- Increase of X_{CO} with column density (primarily)
- Increase of X_{CO} with decreasing metallicity of higher z galaxies

Feldmann + 2012
(arXiv:1204.3910)

The Scatter of the Σ_{H_2} - Σ_{SFR} relation (reloaded)



Scatter

- strong increase with decreasing scale
- above ~ 1 kpc: X_{CO} variations dominate
- below ~ 1 kpc: discreteness noise of SF
- decreases with increasing tracer lifetime
- increases with increasing Δt_{SF}
- scaling $\sim l^{-0.5}$: 2D configuration + finite width of density pdf

Feldmann + 2011 (arXiv:1010.1539)

Feldmann + 2012 (arXiv:1204.3910)

Conclusions

1. Σ_{H_2} - Σ_{SFR} relation clearly important test-bed for SF modeling
2. CO-H₂ conversion factor can lead to increase in slope of the Σ_{H_2} - Σ_{SFR} relation at large surface densities
3. Scatter in the Σ_{H_2} - Σ_{SFR} relation
 - potentially caused by discreteness of SF
 - testable predictions for the scatter and τ_{dep} as function of SF tracer lifetime
4. Scatter roughly $\sim \text{scale}^{-0.5}$; explained by 2D set-up of the gas disk & width of the density pdf
5. Strong Z dependence of the amplitude of the pressure - H₂/HI relation in self-consistent H₂ models

Thank you!