Numerics on a grid

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Outline

Nonparametric galaxy cluster reconstruction

- The physical problem
- Reconstruction algorithm

2 Regular Grids

- Finite differences
- Solving differential equations

3 Adaptive grids

- The problem
- The implementation

Summary

- Reconstruction of a galaxy cluster out of observed quantitites by using gravitational lensing
- Weak lensing gives constraints on a big field size but with bad resolution
- Strong lensing is limited to the cluster core with well observable effects
- Combine both effects to use all avaiable knowledge and to reduce their individual weaknesses
- To reconstruct the cluster you describe it by a adequate quantity \implies lensing potential ψ

and you use the thin lens approximation which projects the cluster on a 2D plane $% \left({{\left[{{{\rm{D}}_{\rm{T}}} \right]}_{\rm{T}}} \right)$

The reconstruction method

- The reconstruction method is based on a least χ^2 minimisation with respect to the lensing potential
- We define a penalty function which contains weak and strong lensing constraints

$$\chi^2(\psi) = \chi^2_w(\psi) + \chi^2_s(\psi)$$

and minimize it with respect to ψ

• Finally we have to connect the lensing potential to observable lensing quantities which is done by convergence and shear

$$\kappa = \frac{1}{2} (\partial_{1,1} + \partial_{2,2}) \psi, \quad \gamma_1 = \frac{1}{2} (\partial_{1,1} - \partial_{2,2}) \psi, \quad \gamma_2 = \partial_{1,2} \psi$$
$$\hat{\epsilon} = \frac{\gamma}{1-\kappa} \equiv g$$
$$\Longrightarrow \chi_w^2 = \frac{|\epsilon - \hat{\epsilon}(\psi)|^2}{\sigma_w^2} = \frac{|\epsilon - \frac{\gamma(\psi)}{1-\kappa(\psi)}|^2}{\sigma_w^2}, \quad \chi_s^2 = \frac{(\det \mathcal{A}(\psi))^2}{\sigma_s^2} = \frac{[(1-\kappa(\psi))^2 - |\gamma(\psi)|^2]^2}{\sigma_s^2}$$

How the grid comes into play

- We now cover the lens plane by a grid of a given resolution, this resolution is normally limited by the quality of the weak lensing data
- Instead of continous quantities we now calculate only with values on every discrete grid point
- The values of convergence and shear which are connected to the potential by second derivatives are approximated by finite differences



	f_{i-2}	f_{i-1}	fi	f_{i+1}	f_{i+2}		fe-3	f_{t-2}	f_{t-1}	ħ	fi+1	f_{i+2}	f_{i+3}
$2(\Delta x)\frac{\partial f}{\partial x}$		- 1	0	1		$12(\Delta x)\frac{\partial f}{\partial x}$		1	-8	0	8	-1	
$(\Delta x)^2 \frac{\partial^2 f}{\partial x^2}$		1	- 2	1		$12(\Delta x)^2 \frac{\partial^2 f}{\partial x^2}$		-1	16	-30	16	-1	
$2(\Delta x)^2 \frac{\partial^3 f}{\partial x^3}$	-1	2	0	-2	1	$8(\Delta x)^3 \frac{\partial^2 f}{\partial x^5}$	1	-8	13	0	-13	8	- 1
$(\Delta x)^4 \frac{\partial^4 f}{\partial x^4}$	1	-4	6	-4	1	$6(\Delta x)^4 \frac{\partial^4 f}{\partial x^4}$	-1	12	-39	56	-39	12	- 1



Implementation of finite differencing

- Values of a certain quantity on the grid are now written as a 1D array $\psi_i, \epsilon_i, \kappa_i, \gamma_i$
- This allows us to write finite differences, means derivatives, as matrix multiplications:

$$\kappa_i = \mathcal{K}_{ij}\psi_j$$

$$\gamma_i^1 = \mathcal{G}_{ij}^1\psi_j$$

$$\gamma_i^2 = \mathcal{G}_{ij}^2\psi_j$$

• As χ^2 -function we now define the sum over all grid points

$$\chi^2 = \sum_{i=1}^{n} \frac{|\epsilon_i - \hat{\epsilon}_i(\psi_j)|^2}{\sigma_{wi}^2} + \sum_{i=1}^{n^*} \frac{(\det \mathcal{A})_i^2}{\sigma_{si}^2}$$



Finite Differences Matrix



Figure: Finite differences Matrix for convergence



Figure: A more detailed look at the matrix

Solving a differential equation on the grid

- What we now want to solve is: $\frac{\partial \chi^2(\psi_k)}{\partial \psi_i} = 0$
- This can now be linearized: $\frac{\partial \chi_w^2(\psi_k)}{\partial \psi_i} = 0$

$$-2\sum_{i=1}^{n}\frac{1}{\sigma_{w_{i}}^{2}}\left[\left(\epsilon_{i}^{1}-\gamma_{i}^{1}\right)\frac{\partial\gamma_{i}^{1}}{\partial\psi_{j}}+\left(\epsilon_{i}^{2}-\gamma_{i}^{2}\right)\frac{\partial\gamma_{i}^{2}}{\partial\psi_{j}}\right]=0$$

$$-2\sum_{i=1}^{n}\frac{1}{\sigma_{wi}^{2}}\left[\left(\epsilon_{i}^{1}-\gamma_{i}^{1}\right)\frac{\partial(\mathcal{G}_{ik}^{1}\psi_{k})}{\partial\psi_{j}}+\left(\epsilon_{i}^{2}-\gamma_{i}^{2}\right)\frac{\partial(\mathcal{G}_{ik}^{2}\psi_{k})}{\partial\psi_{j}}\right]=0$$

$$\sum_{i=1}^{n} \frac{2}{\sigma_{wi}^2} \left[\left(\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 \right) \psi_k - \left(\epsilon_i^1 \mathcal{G}_{ij}^1 + \epsilon_i^2 \mathcal{G}_{ij}^2 \right) \right] = 0$$

 $\mathcal{B}_{jk}\psi_k=\mathcal{V}_j$

Image: Image:

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 $\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} [\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 + \epsilon_i^1 (\mathcal{G}_{ij}^1 \mathcal{K}_{ik} + \mathcal{G}_{ij}^2 \mathcal{K}_{ik})$

 $+\epsilon_i^2(\mathcal{G}_{ij}^1\mathcal{K}_{ik}+\mathcal{G}_{ij}^2\mathcal{K}_{ik})+|\epsilon_i|^2\mathcal{K}_{ij}\mathcal{K}_{ik}+\eta\mathcal{K}_{ij}\mathcal{K}_{ik}]\psi_k$

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 $\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} [\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 + \epsilon_i^1 (\mathcal{G}_{ij}^1 \mathcal{K}_{ik} + \mathcal{G}_{ij}^2 \mathcal{K}_{ik})$

 $+\epsilon_i^2(\mathcal{G}_{ij}^1\mathcal{K}_{ik}+\mathcal{G}_{ij}^2\mathcal{K}_{ik})+|\epsilon_i|^2\mathcal{K}_{ij}\mathcal{K}_{ik}+\eta\mathcal{K}_{ij}\mathcal{K}_{ik}]\psi_k$

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$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} [\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 + \epsilon_i^1 (\mathcal{G}_{ij}^1 \mathcal{K}_{ik} + \mathcal{G}_{ij}^2 \mathcal{K}_{ik})]$$

 $+\epsilon_i^2(\mathcal{G}_{ij}^1\mathcal{K}_{ik}+\mathcal{G}_{ij}^2\mathcal{K}_{ik})+|\epsilon_i|^2\mathcal{K}_{ij}\mathcal{K}_{ik}+\eta\mathcal{K}_{ij}\mathcal{K}_{ik}]\psi_k$

$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} (\epsilon_i^1 \mathcal{G}_{ij}^1 + \epsilon_i^2 \mathcal{G}_{ij}^2 + |\epsilon_i|^2 \mathcal{K}_{ij} + \eta \kappa_i^{(n_2-1)} \mathcal{K}_{ij})$$

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$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} [\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 + \epsilon_i^1 (\mathcal{G}_{ij}^1 \mathcal{K}_{ik} + \mathcal{G}_{ij}^2 \mathcal{K}_{ik})]$$

 $+\epsilon_i^2(\mathcal{G}_{ij}^1\mathcal{K}_{ik}+\mathcal{G}_{ij}^2\mathcal{K}_{ik})+|\epsilon_i|^2\mathcal{K}_{ij}\mathcal{K}_{ik}+\eta\mathcal{K}_{ij}\mathcal{K}_{ik}]\psi_k$

$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} (\epsilon_i^1 \mathcal{G}_{ij}^1 + \epsilon_i^2 \mathcal{G}_{ij}^2 + |\epsilon_i|^2 \mathcal{K}_{ij} + \eta \kappa_i^{(n_2-1)} \mathcal{K}_{ij})$$

$$-\sum_{i=1}^{n^*}rac{-4det\mathcal{A}_i}{\sigma_{si}^2}((1-\kappa_i)\mathcal{K}_{ij}-\gamma_i^1\mathcal{G}_{ij}^1-\gamma_i^2\mathcal{G}_{ij}^1)$$

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$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} [\mathcal{G}_{ij}^1 \mathcal{G}_{ik}^1 + \mathcal{G}_{ij}^2 \mathcal{G}_{ik}^2 + \epsilon_i^1 (\mathcal{G}_{ij}^1 \mathcal{K}_{ik} + \mathcal{G}_{ij}^2 \mathcal{K}_{ik})]$$

$$+\epsilon_i^2(\mathcal{G}_{ij}^1\mathcal{K}_{ik}+\mathcal{G}_{ij}^2\mathcal{K}_{ik})+|\epsilon_i|^2\mathcal{K}_{ij}\mathcal{K}_{ik}+\eta\mathcal{K}_{ij}\mathcal{K}_{ik}]\psi_k$$

$$\sum_{i=1}^{n} \frac{1}{(1-\kappa_i)^2 \sigma_{wi}^2} \left(\epsilon_i^1 \mathcal{G}_{ij}^1 + \epsilon_i^2 \mathcal{G}_{ij}^2 + |\epsilon_i|^2 \mathcal{K}_{ij} + \eta \kappa_i^{(n_2-1)} \mathcal{K}_{ij} \right)$$

$$-\sum_{i=1}\frac{-4det\mathcal{A}_i}{\sigma_{si}^2}((1-\kappa_i)\mathcal{K}_{ij}-\gamma_i^1\mathcal{G}_{ij}^1-\gamma_i^2\mathcal{G}_{ij}^1)$$

The problem now is that this system is not linear in ψ anymore. This is solved by an iterative process.

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Solving the linear system

- The last task is solving the linear system to obtain the solution for ψ
- In my code this is done by standard GSL routines (LU decomposition)
- In the end it is a little bit more complicated: iterative process, increasing grid resolution
- Especially you have to think about fast routines to build up the coefficient matrix



Figure: The complete coefficient matrix

Adaptive grids: The problem

- During the cluster reconstruction you want to focus on the innermost cluster center and cut out the core
- If you do so you are not allowed to work on regular rectangular grids anymore, they become arbitrarily shaped now and things become more complicated
- Again all points of the grid are stored in a 1D array
- Problems: Positions of points for finite differencing, you have to be more careful with the borders





Adaptive grids: The implementation

- First, a proper border is constructed to avoid the finite differencing from taking points which don't exist on that grid
- Every point of the grid needs properties on what finite differencing scheme could be used
- To take into account the more complex shape of the grid, additional information for every point needs to be stored
- These informations are: Position in original grid, type of the point, position of points up and down in the grid



Figure: Typical type map of an adaptive grid

Again some matrices



Figure: Adaptive finite differences matrix



Figure: LU decomposition

- Use of a grid with finite resolution to describe a lensing plane
- Use of finite differencing schemes to approximate derivatives
- Linearising of differential equations on our grid
- Use of adaptive grids to focus on more complicated structures